

Nanoassembly technique of carbon nanotubes for hybrid circuit-QED ^{EP}

Cite as: Appl. Phys. Lett. **117**, 114001 (2020); <https://doi.org/10.1063/5.0021838>

Submitted: 15 July 2020 . Accepted: 06 September 2020 . Published Online: 17 September 2020

T. Cubaynes, L. C. Contamin, M. C. Dartailh, M. M. Desjardins, A. Cottet, M. R. Delbecq, and T. Kontos ^{ID}

COLLECTIONS

Paper published as part of the special topic on [Hybrid Quantum Devices](#)

^{EP} This paper was selected as an Editor's Pick



View Online



Export Citation



CrossMark

ARTICLES YOU MAY BE INTERESTED IN

[Black phosphorus field effect transistors stable in harsh conditions via surface engineering](#)
Applied Physics Letters **117**, 111602 (2020); <https://doi.org/10.1063/5.0021335>

[Strong interface-induced spin-charge conversion in YIG/Cr heterostructures](#)
Applied Physics Letters **117**, 112402 (2020); <https://doi.org/10.1063/5.0017745>

[Phononic bandgap and phonon anomalies in HfN and HfN/ScN metal/semiconductor superlattices measured with inelastic x-ray scattering](#)
Applied Physics Letters **117**, 111901 (2020); <https://doi.org/10.1063/5.0020935>



Your Qubits. Measured.

Meet the next generation of quantum analyzers

- Readout for up to 64 qubits
- Operation at up to 8.5 GHz, mixer-calibration-free
- Signal optimization with minimal latency

Find out more



Nanoassembly technique of carbon nanotubes for hybrid circuit-QED

Cite as: Appl. Phys. Lett. **117**, 114001 (2020); doi: [10.1063/5.0021838](https://doi.org/10.1063/5.0021838)

Submitted: 15 July 2020 · Accepted: 6 September 2020 ·

Published Online: 17 September 2020



View Online



Export Citation



CrossMark

T. Cubaynes,^{a)} L. C. Contamin, M. C. Dartailh, M. M. Desjardins, A. Cottet, M. R. Delbecq, and T. Kontos^{b)} 

AFFILIATIONS

Laboratoire de Physique de l'Ecole normale supérieure, ENS, Université PSL, CNRS, Sorbonne Université, Université Paris-Diderot, Sorbonne Paris Cité, 75005 Paris, France

Note: This paper is part of the Special Issue on Hybrid Quantum Devices.

^{a)}Electronic mail: tino.cubaynes@kit.edu

^{b)}Author to whom correspondence should be addressed: takis.kontos@ens.fr

ABSTRACT

A complex quantum dot circuit based on a clean and suspended carbon nanotube embedded in a circuit quantum electrodynamic (cQED) architecture is a very attractive platform to investigate a large spectrum of physics phenomena ranging from qubit physics to nanomechanics. We demonstrate a carbon nanotube transfer process allowing us to integrate clean carbon nanotubes into complex quantum dot circuits inside a cQED platform. This technique is compatible with various contacting materials such as superconductors or ferromagnets. This makes it suitable for hybrid quantum devices. Our results are based on eight different devices demonstrating the robustness of this technique.

Published under license by AIP Publishing. <https://doi.org/10.1063/5.0021838>

One challenge for the investigation of quantum phenomena is the fabrication of systems that are sufficiently decoupled from their environment such that they manifest their quantum nature. This decoupling usually comes at the cost of low tunability of the system parameters. From this perspective, a suspended and clean carbon nanotube embedded in a circuit, hence allowing large tunability of its parameters, and with the “active” part of the system confined to the suspended section of the nanotube, is very attractive.¹

Because it can be suspended, a carbon nanotube can be engineered to be arbitrarily far from any interface, which is known to be one of the main sources of charge noise.² The cleanliness of the nanotube also plays a key role since it ensures a very low amount of charge fluctuators. These two points suggest a long electronic coherence time in such carbon nanotube-based circuits. Besides, there is also the possibility to tailor the phonon spectrum of the suspended nanotube via gate voltages.^{3–5} In the case of a qubit, this means that one can tune the relaxation time due to electron-phonon interaction.⁶ The nanotube can also be connected with different types of metals, making this platform interesting for studying hybrid circuits.⁷

Recently, various carbon nanotube transfer techniques have been developed.^{1,8–13} The common idea is to grow carbon nanotubes on a separate substrate and then transfer one of the nanotubes onto the circuit at the final step of the process. Depending on the growth and

the transfer conditions, several of these works have demonstrated the clean nature of the transferred nanotube,^{1,10,11} up to the observation of a 1D Wigner crystal.¹⁴

This technique has the advantage to provide suspended carbon nanotubes and is now largely used in nanomechanics experiments.^{15–17} The suspended nature of the nanotube is also desirable to limit charge noise in quantum dot circuits. Also, because the fabrication of the circuit and the synthesis of the carbon nanotubes are now completely independent, there is a great flexibility in the circuit fabrication process. More recently, this technique has been adapted to the cQED platform^{11,13} by growing carbon nanotubes on a fork-like chip.^{9,12}

Here, we present a fabrication technique inheriting from these recent advances. The nanotubes are grown on a substrate containing a comb of 48 cantilevers allowing us to do a pre-selection of the transferred nanotube. In addition to producing clean suspended carbon nanotube-based circuits embedded in a cQED platform, we show that this technique is compatible with a superconductor and a ferromagnet as metallic contacts, making it suitable for many types of hybrid quantum devices. Our results are based on eight devices with different geometry and contact materials.

The process includes two distinct samples: the growth chip, on which carbon nanotubes are grown, and the circuit chip. For the

growth chip, we use a commercial cantilever chip with 48 cantilevers and a spacing of $30\text{ }\mu\text{m}$ between cantilevers as depicted in Fig. 1(a). Carbon nanotubes are grown all over the sample, using a standard CVD-growth recipe based on CH_4 feedstock gas and using Mo-Fe nanoparticles as a catalyst.¹⁸ Depending on the catalyst distribution, several carbon nanotubes will be suspended between adjacent cantilevers. The large number of cantilevers is particularly useful to integrate multiple nanotubes on the same circuit chip. Prior to the transfer of a carbon nanotube to the circuit chip, pictures of the tips of all the cantilevers are taken using a Scanning Electron Microscope (SEM) at low acceleration voltage (2 kV). While common belief is that SEM observation induces the deposition of the hydrocarbon layer onto the nanotube,¹⁹ hence should introduce disorder in the electronic spectrum of the nanotube, we have found that it is still possible to obtain clean transport spectra using such a brief observation. One possible explanation for this low contamination of suspended carbon nanotubes by the e-beam exposure is the fact that most of the hydrocarbon is on the surface of the substrate, and thus, suspended nanotubes are much less affected by e-beam induced diffusion of hydrocarbon.²⁰ This first step allows us to pre-select the isolated carbon nanotubes. Images of carbon nanotubes suspended in between adjacent cantilevers are shown in Fig. 1(a). Between the localization of carbon nanotubes and the stamping process, the cantilever chip is stored in a vacuum chamber (pressure: 5×10^{-7} mbar), to minimize

the exposition time of nanotubes to the ambient atmosphere down to approximately 10 min.

On the circuit chip, the niobium microwave cavity and the circuit electrodes are fabricated using e-beam lithography as in previous work.²¹ In addition, two trenches on both sides of the circuit are etched by reactive ion etching using SF_6 gas. This process allows us to etch $10\text{--}15\text{ }\mu\text{m}$ deep trenches, which are used to ensure a proper contacting of the nanotube to the circuit electrodes [see Figs. 1(c) and 1(d)].

The transfer of the carbon nanotube to the circuit chips is realized in a dedicated vacuum chamber with a base pressure of 5×10^{-7} mbar. A schematic of this chamber is presented in Fig. 2(c). The circuit chip is placed on a rotary arm, so that it can be either in an horizontal position for the transfer of the nanotubes or in a vertical position, facing an argon gun, in order to clean the surface of the contact electrodes, as well as remove unwanted nanotubes. The growth chip is fixed with a tilt angle of 45° with respect to the horizontal plane, to ensure a good visibility from the top view, while maintaining a small footprint of the circuit chip. The position of the growth chip is controlled by piezo-motors and micro-manipulators. There is also the possibility to isolate the growth chip from the main chamber, in order to protect the carbon nanotubes during the cleaning of the circuit chip with the argon gun or during the replacement of the circuit chip.

The growth chip is lowered using the piezo-motor stage, and the contact is detected by monitoring the current between the two external contact electrodes (at $V_{\text{bias}} = 0.5\text{--}1.5\text{ V}$) and the two inner contacts, as depicted in Fig. 1(c). After the contact, the resistance of the circuit typically ranges between $10\text{ M}\Omega$ and $100\text{ G}\Omega$. Similar to ref.^{1,22} the two external sections of the nanotube are cut by driving a large current through it (typically between $10\text{ }\mu\text{A}$ and $20\text{ }\mu\text{A}$). During this operation, all the other contacts and gates are set to a floating potential, to avoid accidental cutting of the central section. Two examples of I-V curves corresponding to the cutting of the nanotube are shown in the inset of Fig. 2(b). The two curves show very similar features, such as the current and the bias voltage at which the nanotube is cut. This shows that the nanotube has almost identical properties at two different sections a few μm away from each other, attesting the cleanliness of the transferred nanotube.

One advantage of this cutting procedure is that it acts as a local annealing of the interface between the contact metal and the nanotube, drastically lowering the contact resistance, which can reach values below $1\text{ M}\Omega$ (see Table I). The cutting step also allows us to distinguish between a single nanotube, which displays a single current drop [see Fig. 2(b)], and a bundle of nanotubes, which display multiple current drops. At this stage, it is possible to further improve the contact quality by driving a high current in the central section of the device, being careful not to exceed the cutting current. If the contacted nanotube turns out to be a bundle or display unwanted characteristics, we remove it using the Ar gun. To obtain a suspended nanotube over the gate array, we found that the spacing of the contact electrode should not exceed ten times the height difference between contact and gate electrodes.

As a last characterization study of the nanotube at room temperature, we measure the gate dependence of the current in the central section, in order to differentiate between a small gap, a semiconducting or a metallic nanotube. Finally, during the transfer of the circuit chip

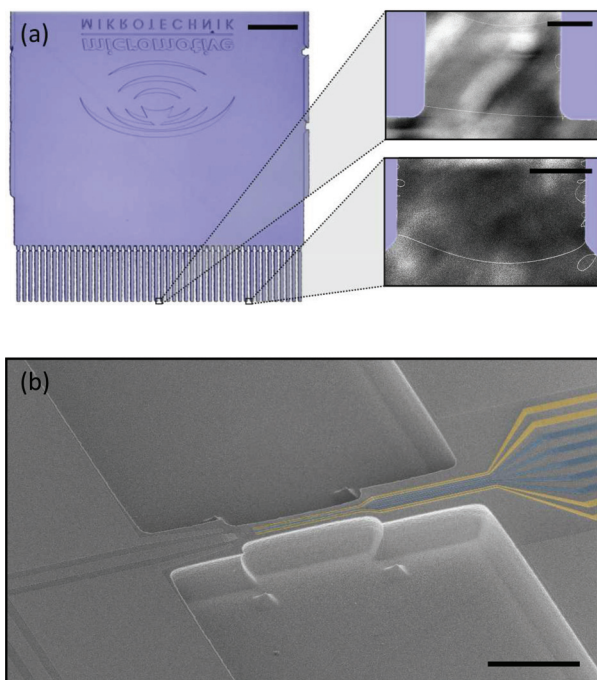


FIG. 1. (a) Optical photograph of the growth chip. It is composed of 48 cantilevers, which are visible on the bottom edge of the chip. Scale bar: $500\text{ }\mu\text{m}$. The zoom-in on the cantilevers shows isolated carbon nanotubes suspended in between two adjacent cantilevers. Scale bar: $10\text{ }\mu\text{m}$. (b) False color scanning electron micrograph of the circuit chip (device 5). The circuit, constituted of the gate electrodes (in the dark blue) and the contact electrode (in yellow), is sandwiched in between the two trenches. Scale bar: $50\text{ }\mu\text{m}$.

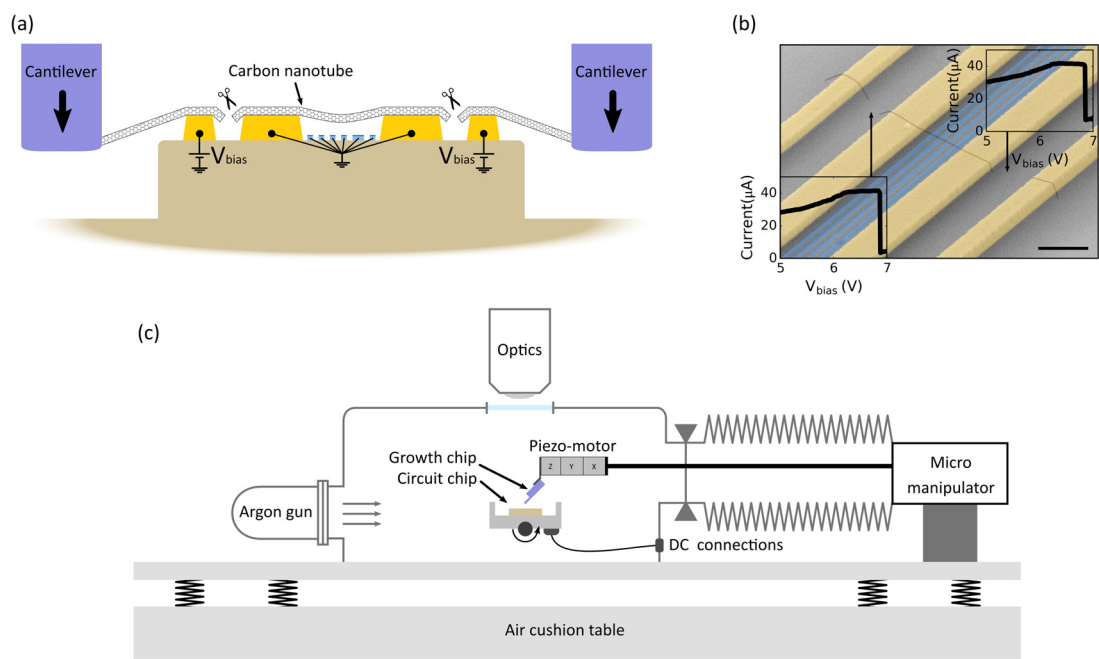


FIG. 2. Principle of the transfer technique. (a) Schematics of the carbon nanotube transfer process. The cantilevers (in purple) are lowered until the carbon nanotube makes the connection between external and internal contacts (in orange) of the circuit chip. The nanotube is then cut on both sides of the circuit by successively biasing only one of the two external contacts. (b) False color scanning electron micrograph of the region where the nanotube is integrated (Device 1). The carbon nanotube (in black) is suspended over a gate array (in dark blue). Scale bar: 1 μ m. Insets: current as a function of the bias voltage applied on the left external contact (bottom-left inset) and on the right external contact (top-right inset). (c) Schematic of the vacuum chamber used for the transfer of the carbon nanotube. The argon gun on the left side of the chamber is used for cleaning the circuit chip or to remove unwanted carbon nanotubes. The growth chip (in purple) can be locked-up in the bellows chamber on the right side of the chamber, for protecting the nanotubes during ventilation of the main chamber or the use of the argon gun.

into the cryostat, all the electrodes are grounded and the sample is maintained under a nitrogen atmosphere or a low vacuum.

The results discussed in the following are based on 8 different devices with various geometries and contact metals. The measurements have been performed at a base temperature of 20 mK (250 mK for device 8). We present in Fig. 3 a single dot stability diagram in the Coulomb blockade regime. The continuous evolution of the contact transparency as the gate voltage is swept and the absence of charge jumps are indicative of the clean environment of the nanotube.

The tunability of the circuit parameters has also been observed in double quantum dot circuits. Figure 4 presents stability diagrams of

three double quantum dots devices with different carbon nanotube electronic behaviors: small-gap in Fig. 4(a) and semiconducting in Figs. 4(b) and 4(d). The transition from double quantum dot behavior to a large single dot behavior is also visible in Fig. 4(c), showing again the weak influence of disorder on the confinement potential, which is instead dominated by electrostatic gating. The control over the interdot coupling was also demonstrated via cavity transmission measurements, where the transition from the resonant to the dispersive regime was observed (see the supplementary material). In the resonant regime, the resonator not only is a readout circuit but can also be used for manipulation and coupling of qubits (spin or topological). The capability to

TABLE I. Extensive characterization of eight samples.

	Contact metal	Geometry	Quality factor	RT resistance (M Ω)	tunnel rates $\Gamma_1/2\pi$ (GHz)	tunnel rates $\Gamma_2/2\pi$ (GHz)	Charging energy (meV)
Device 1	Au	3 gates	7400	0.08	0.5	16.5	1
Device 2	Au	3 gates	16000	7	0.3	253	7
Device 3	PdNi(25 nm)/Pd(4 nm)	5 gates	4000	2
Device 4	PdNi(25 nm)/Pd(4 nm)	5 gates	...	0.8	0.5	0.5	2
Device 5	PdNi(25 nm)/Pd(4 nm)	5 gates	...	0.23	0.5	67	1.5
Device 6	PdNi(25 nm)/Pd(4 nm)	5 gates	4000	0.55	2.2–2.4
Device 7	Nb(45 nm)/Pd(10 nm)	3 gates	...	0.2	33.6	3.7×10^{-4}	1.0
Device 8	Nb(45 nm)/Pd(10 nm)	1 gate	...	0.25–1.0	7	80	2.4

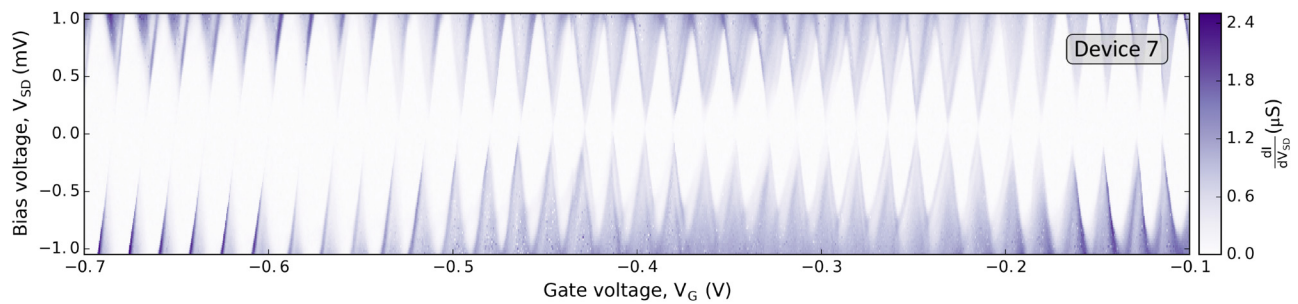


FIG. 3. Differential conductance measured in device 7 vs bias voltage V_{SD} and gate voltage V_G at $B = 0$ T. The fact that the Coulomb diamonds do not close at $V_{SD} = 0$ V is a manifestation of the superconducting contacts (see Table I). This 2D map is spanning over a large range of gate voltages in the Coulomb blockade regime (more than 30 Coulomb diamonds) and shows continuous evolution of the contact transparency, which indicates an electrostatic control of the coupling rates. From the contrast of the coulomb diamonds, one can notice that in the gate voltage region: $-0.7 \text{ V} < V_G < -0.5 \text{ V}$, we have $\Gamma_1 < \Gamma_2$, then for $-0.4 \text{ V} < V_G < -0.3 \text{ V}$, we have $\Gamma_1 \simeq \Gamma_2$, and finally for $-0.2 \text{ V} < V_G < -0.1 \text{ V}$, we have $\Gamma_1 \simeq \Gamma_2$. This observation further emphasizes the electrostatic control of the coupling rates.

tune an electronic transition in the carbon nanotube to be resonant with the microwave cavity is, therefore, crucial for cQED applications.

Importantly, we also applied this technique to multiple contact-ing materials. Table I gathers information on the eight samples investigated in this article, including the coupling rates obtained with various materials. Using this transfer technique on ferromagnetic PdNi(25 nm)/Pd(4 nm) contacts allowed us to induce local polarization of electronic spin states in the nanotube and to couple it to cavity photons.²³ This technique can also be applied to superconducting contacts. In Fig. 3, we observe a non-closing of Coulomb diamonds and shift of their apex characteristic of a superconducting contact on a quantum dot²⁴ (see Fig. S3 of the supplementary for a zoom on a specific Coulomb diamond). The size of the superconducting gap can be extracted from the distance between the apices. We find Δ of the order

of 0.35–0.4 meV with large coupling rates (7 GHz and 80 GHz), using Nb(45 nm)/Pd(10 nm) contacts.

We believe that the cleanliness of our quantum dot circuit and the compatibility with various contacting materials are partially due to the fact that the transfer process is performed under vacuum, which results in a cleaner nanotube-metal interface.

Transferring a carbon nanotube on top of circuit electrodes has proven to be a very efficient technique to integrate clean and suspended carbon nanotubes into complex circuit designs with the possibility to be part of a cQED platform. Here, we show that this approach can also be adapted to hybrid circuits containing superconducting and ferromagnetic materials using commercial cantilevers, hence enlarging the scope of this technique.^{25–30}

See the [supplementary material](#) for extended data for the resonant case and its tunability. The theory of cavity-double quantum dot coupling is briefly explained, and the corresponding measurements are shown in Fig. S2. A close-up on specific Coulomb diamonds of Fig. 3 is also presented in Fig. S3 together with a cut at $V_G = -0.48 \text{ V}$, which shows the superconducting gap.

The devices have been made within the consortium Salle Blanche Paris Centre. The authors gratefully acknowledge help from José Palomo, Aurélie Pierret, and Michael Rosticher. This work was supported by the ERC Starting Grant “CirQys,” the ERC Proof of Concept grant “QUBE,” the ANR “FunTheme,” and the Quanterra grant “SuperTop.”

DATA AVAILABILITY

The authors declare that the main data supporting the findings of this study are available within this article. Extra data are available from the corresponding authors upon reasonable request.

REFERENCES

- ¹J. Weissman, M. Honig, S. Pecker, A. Benyamini, A. Hamo, and S. Ilani, “Realization of pristine and locally tunable one-dimensional electron systems in carbon nanotubes,” *Nat. Nanotechnol.* **8**, 569–574 (2013).
- ²T. Sharf, J. W. Kevek, T. DeBorde, J. L. Wardini, and E. D. Minot, “Origins of charge noise in carbon nanotube field-effect transistor biosensors,” *Nano Lett.* **12**, 6380–6384 (2012).

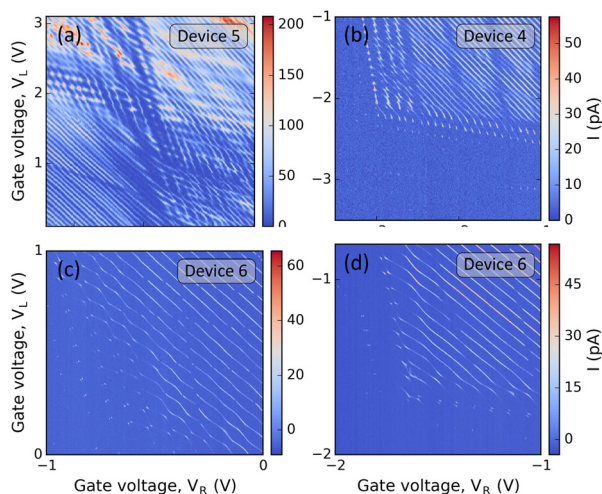


FIG. 4. Double quantum dot stability diagram measured on three devices. (a) Device 5 displays a narrow-gap behavior, and one can see the four different charge distributions depending on whether the two dots are filled with electrons or holes. Carbon nanotubes (b) (Device 4) and (d) (device 6) are semiconducting, and only the electron–electron filling distribution is visible in the gate–gate plane. (c) For other gate voltage parameters of device 6, we observed a smooth transition from a single dot behavior (parallel lines in the top right region) to a double quantum dot behavior (anti-crossing in the bottom left region).

- ³V. Sazonova, Y. Yaish, H. Üstünel, D. Roundy, T. A. Arias, and P. L. McEuen, "A tunable carbon nanotube electromechanical oscillator," *Nature* **431**, 284–287 (2004).
- ⁴B. Witkamp, M. Poot, and H. S. J. van der Zant, "Bending-mode vibration of a suspended nanotube resonator," *Nano Lett.* **6**, 2904–2908 (2006).
- ⁵B. Lassagne, D. Garcia-Sanchez, A. Aguasca, and A. Bachtold, "Ultrasensitive mass sensing with a nanotube electromechanical resonator," *Nano Lett.* **8**, 3735–3738 (2008).
- ⁶A. Cottet and T. Kontos, "Spin quantum bit with ferromagnetic contacts for circuit QED," *Phys. Rev. Lett.* **105**, 160502 (2010).
- ⁷A. Cottet, M. C. Dartiailh, M. M. Desjardins, T. Cubaynes, L. C. Contamin, M. Delbecq, J. J. Vienneot, L. E. Bruhat, B. Douçot, and T. Kontos, "Cavity QED with hybrid nanocircuits: From atomic-like physics to condensed matter phenomena," *J. Phys.: Condens. Matter* **29**, 433002 (2017).
- ⁸C. C. Wu, C. H. Liu, and Z. Zhong, "One-step direct transfer of pristine single-walled carbon nanotubes for functional nanoelectronics," *Nano Lett.* **10**, 1032–1036 (2010).
- ⁹M. Muoth and C. Hierold, "Transfer of carbon nanotubes onto microactuators for hysteresis-free transistors at low thermal budget," in *IEEE 25th International Conference on Micro Electro Mechanical Systems (MEMS)* (IEEE, 2012), pp. 1352–1355.
- ¹⁰F. Pei, E. A. Laird, G. A. Steele, and L. P. Kouwenhoven, "Valley-spin blockade and spin resonance in carbon nanotubes," *Nat. Nanotechnol.* **7**, 630–634 (2012).
- ¹¹V. Ranjan, G. Puebla-Hellmann, M. Jung, T. Hasler, A. Nunnenkamp, M. Muoth, C. Hierold, A. Wallraff, and C. Schönenberger, "Clean carbon nanotubes coupled to superconducting impedance-matching circuits," *Nat. Commun.* **6**, 7165 (2015).
- ¹²S. Blien, P. Steger, A. Albang, N. Paradiso, and A. K. Hüttel, "Quartz tuning-fork based carbon nanotube transfer into quantum device geometries," *Phys. Status Solidi B* **255**, 1800118 (2018).
- ¹³S. Blien, P. Steger, N. Hüttner, R. Graaf, and A. K. Hüttel, "Quantum capacitance mediated carbon nanotube optomechanics," *Nat. Commun.* **11**, 1636 (2020).
- ¹⁴I. Shapir, A. Hamo, S. Pecker, C. P. Moca, Ö. Legeza, G. Zarand, and S. Ilani, "Imaging the electronic Wigner crystal in one dimension," *Science* **364**, 870–875 (2019).
- ¹⁵A. Benyamini, A. Hamo, S. V. Kusminskiy, F. von Oppen, and S. Ilani, "Real-space tailoring of the electron-phonon coupling in ultraclean nanotube mechanical resonators," *Nat. Phys.* **10**, 151–156 (2014).
- ¹⁶I. Khivrich, A. A. Clerk, and S. Ilani, "Nanomechanical pump-probe measurements of insulating electronic states in a carbon nanotube," *Nat. Nanotechnol.* **14**, 161–167 (2019).
- ¹⁷Y. Wen, N. Ares, F. J. Schupp, T. Pei, G. A. D. Briggs, and E. A. Laird, "A coherent nanomechanical oscillator driven by single-electron tunnelling," *Nat. Phys.* **16**, 75–82 (2020).
- ¹⁸J. J. Vienneot, J. Palomo, and T. Kontos, "Stamping single wall nanotubes for circuit quantum electrodynamics," *Appl. Phys. Lett.* **104**, 113108 (2014).
- ¹⁹J. H. Choi, J. Lee, S. M. Moon, Y.-T. Kim, H. Park, and C. Y. Lee, "A low-energy electron beam does not damage single-walled carbon nanotubes and graphene," *J. Phys. Chem. Lett.* **7**, 4739–4743 (2016).
- ²⁰W. K. Wong, A. Nojeh, and R. F. W. Pease, "Parameters and mechanisms governing image contrast in scanning electron microscopy of single-walled carbon nanotubes," *Scanning* **28**, 219–227 (2006).
- ²¹J. J. Vienneot, M. C. Dartiailh, A. Cottet, and T. Kontos, "Coherent coupling of a single spin to microwave cavity photons," *Science* **349**, 408–411 (2015).
- ²²P. G. Collins, "Engineering carbon nanotubes and nanotube circuits using electrical breakdown," *Science* **292**, 706–709 (2001).
- ²³T. Cubaynes, M. R. Delbecq, M. C. Dartiailh, R. Assouly, M. M. Desjardins, L. C. Contamin, L. E. Bruhat, Z. Leghtas, F. Mallet, A. Cottet, and T. Kontos, "Highly coherent spin states in carbon nanotubes coupled to cavity photons," *npj Quantum Inf.* **5**, 47 (2019).
- ²⁴L. E. Bruhat, J. J. Vienneot, M. C. Dartiailh, M. M. Desjardins, T. Kontos, and A. Cottet, "Cavity photons as a probe for charge relaxation resistance and photon emission in a quantum dot coupled to normal and superconducting continua," *Phys. Rev. X* **6**, 021014 (2016).
- ²⁵M. Mergenthaler, A. Nersisyan, A. Patterson, M. Esposito, A. Baumgartner, C. Schönenberger, G. A. D. Briggs, E. A. Laird, and P. J. Leek, "Realization of a carbon-nanotube-based superconducting qubit," preprint [arXiv:1904.10132](https://arxiv.org/abs/1904.10132) (2019).
- ²⁶M. M. Desjardins, L. C. Contamin, M. R. Delbecq, M. C. Dartiailh, L. E. Bruhat, T. Cubaynes, J. J. Vienneot, F. Mallet, S. Rohart, A. Thiaville, A. Cottet, and T. Kontos, "Synthetic spin-orbit interaction for Majorana devices," *Nat. Mater.* **18**, 1060–1064 (2019).
- ²⁷K. E. Khosla, M. R. Vanner, N. Ares, and E. A. Laird, "Displacemon electromechanics: How to detect quantum interference in a nanomechanical resonator," *Phys. Rev. X* **8**, 021052 (2018).
- ²⁸W. Qin, A. Miranowicz, G. Long, J. Q. You, and F. Nori, "Proposal to test quantum wave-particle superposition on massive mechanical resonators," *npj Quantum Inf.* **5**, 58 (2019).
- ²⁹P. Stadler, W. Belzig, and G. Rastelli, "Ground-state cooling of a mechanical oscillator by interference in Andreev reflection," *Phys. Rev. Lett.* **117**, 197202 (2016).
- ³⁰O. Lesser, G. Shavit, and Y. Oreg, "Topological superconductivity in carbon nanotubes with a small magnetic flux," *Phys. Rev. Res.* **2**, 023254 (2020).