

# Yielding and shear banding in amorphous solids



**Srikanth Sastry**

Jawaharlal Nehru Centre for  
Advanced Scientific Research  
Bengaluru, India

**Yielding versus depinning in disordered systems**  
**ENS, Oct 22 – 24, 2018**

# Outline

## Introduction

- Mechanical response of amorphous solids
- Oscillatory shear deformation of a model glass
- The Yielding transition
- Avalanches
- Shear banding
- Summary



P. Leishangthem  
(JNCASR)



Anshul D. S. Parmar  
(JNCASR/TIFR)



Saurabh Kumar  
(JNCASR/Koln)

Earlier work: In Collaboration with:

**Giuseppe Foffi (EPFL/Orsay)**

**Davide Fiocco (EPFL/Google/Frontiers)**



*Fiocco, Foffi, Sastry, **Phys Rev E** 88, 020301(R) (2013), **Phys Rev Lett** 112 025702 (2014), **JPCM** 27, 194130 (2015)*

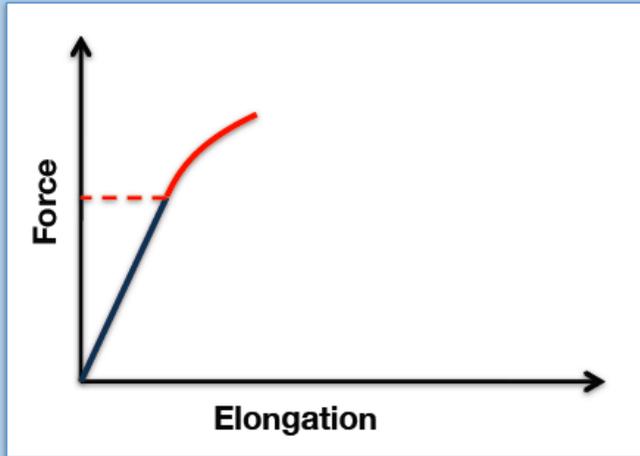
*- Leishangthem, Parmar, Sastry "The yielding transition in amorphous solids under oscillatory shear deformation" **Nature Comm.** (2017)*

*- Parmar, Kumar, Sastry. Strain localisation above the yielding point in cyclically deformed glasses. [arXiv:1806.02464](https://arxiv.org/abs/1806.02464)*

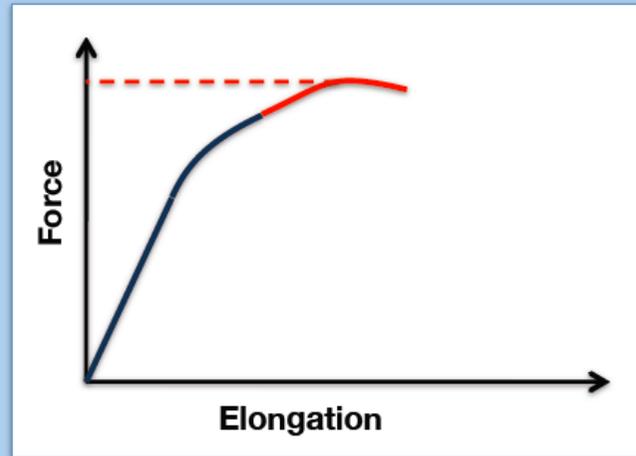
---

# Characterizing mechanical behavior of solids

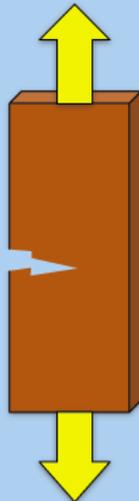
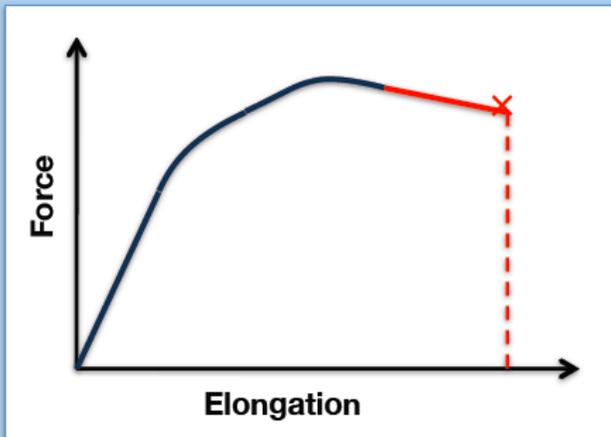
- Yield Stress – Onset of Irreversibility



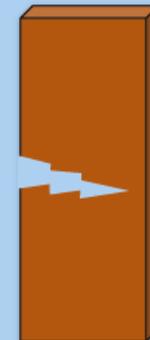
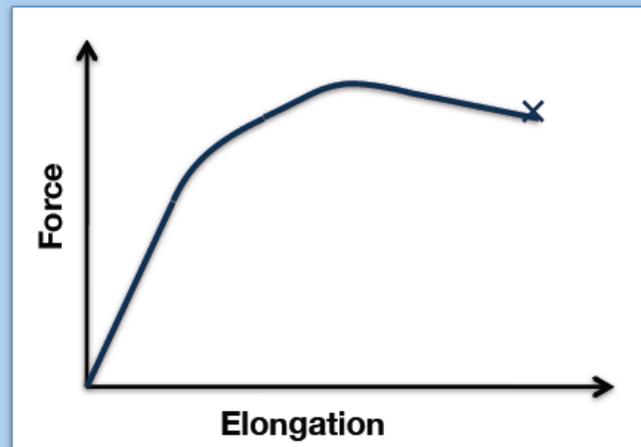
- Strength – Maximum stress attainable



- Ductility – Strain to failure



- Toughness – Energy expended per unit crack advance

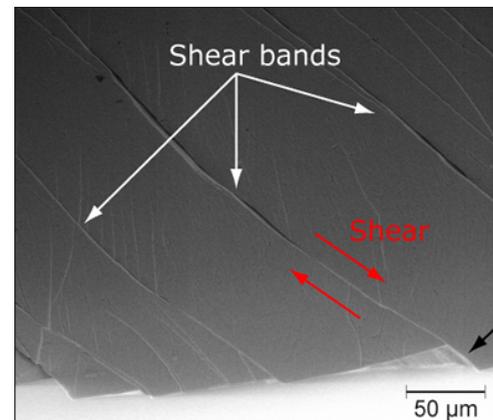
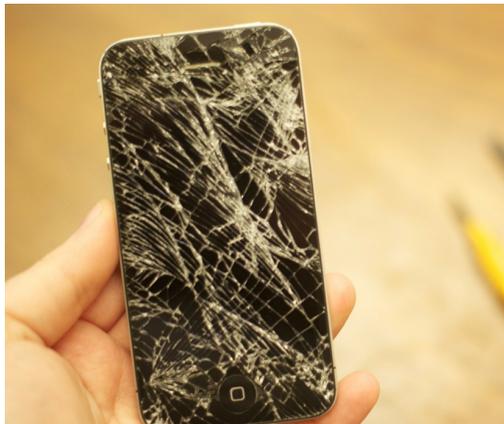
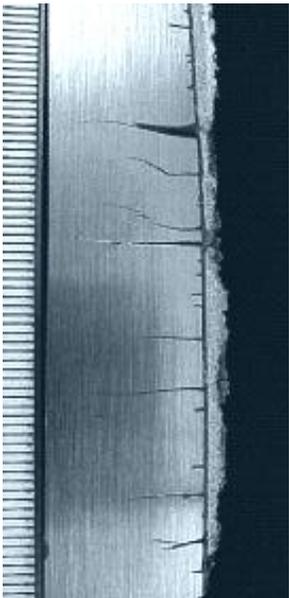


# Yielding and Mechanical Failure

External stresses can cause (amorphous) solids to fail when deformed beyond a point of “**mechanical failure**”.

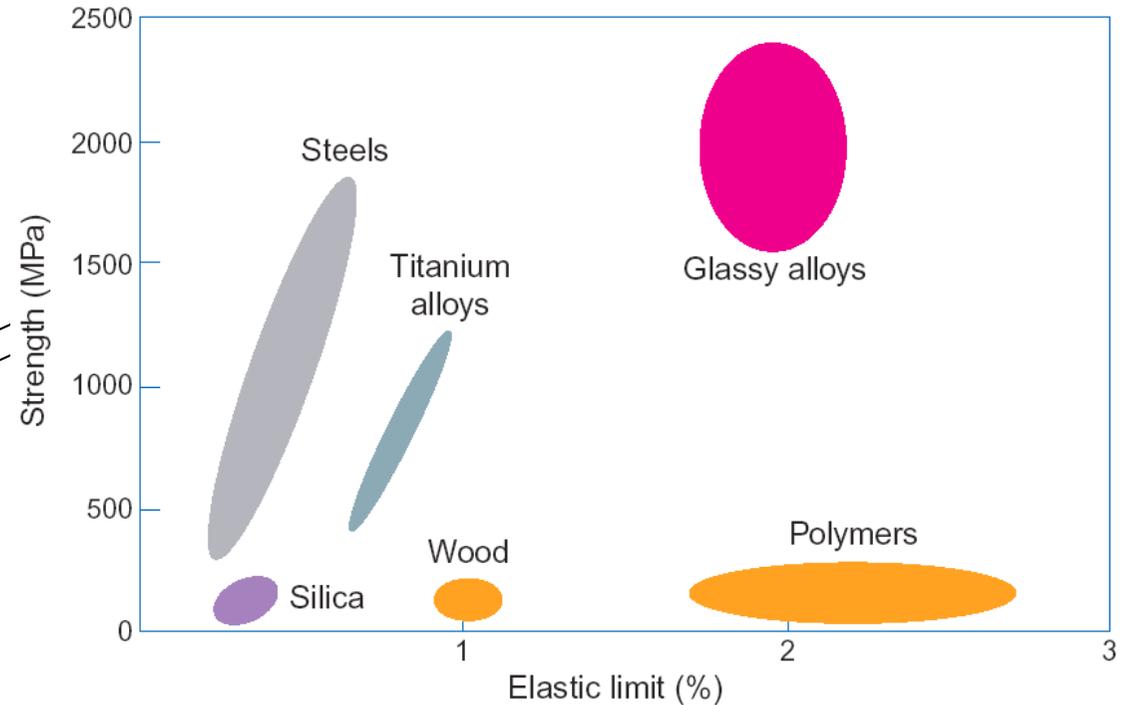
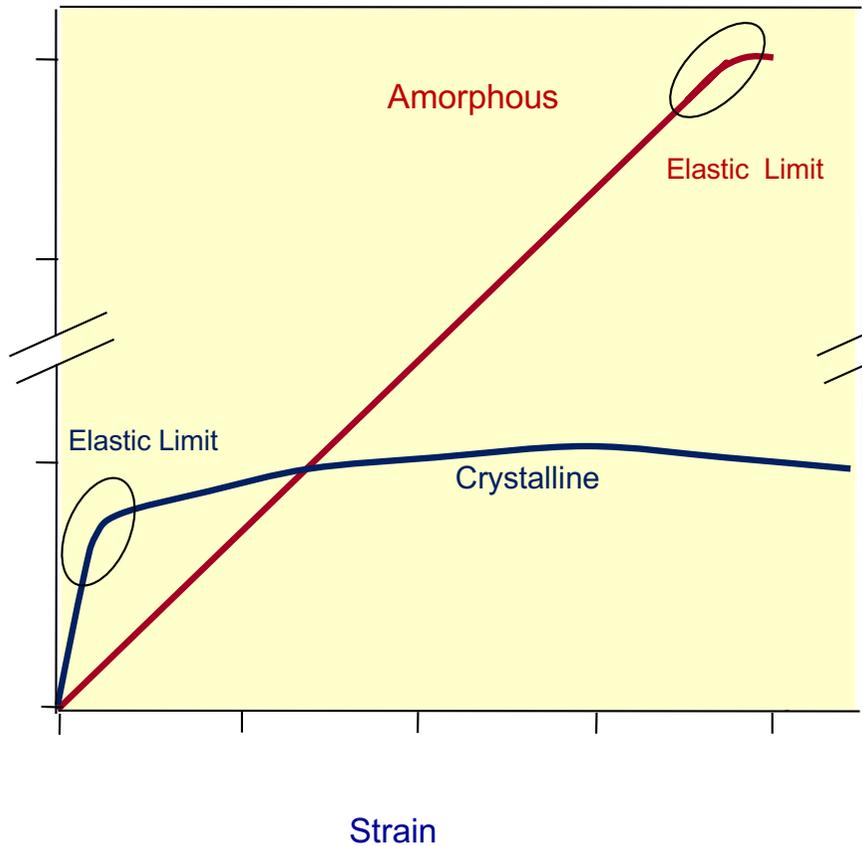
Important for materials performance, and characterization.

Gorilla glass, metallic glasses, some recent important examples.



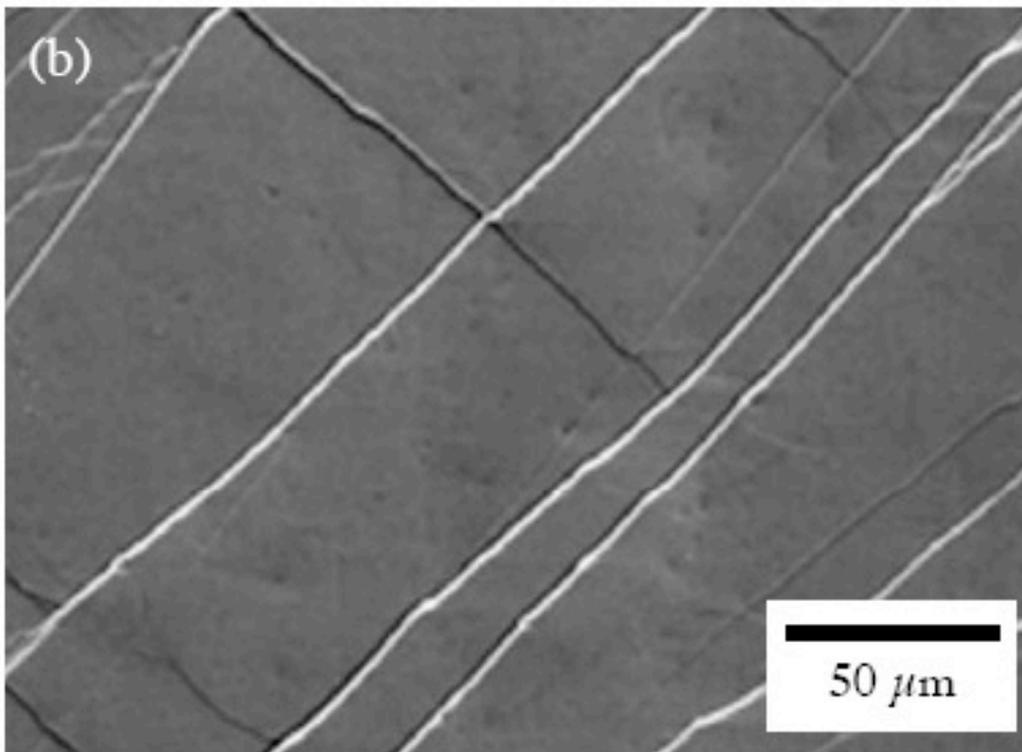
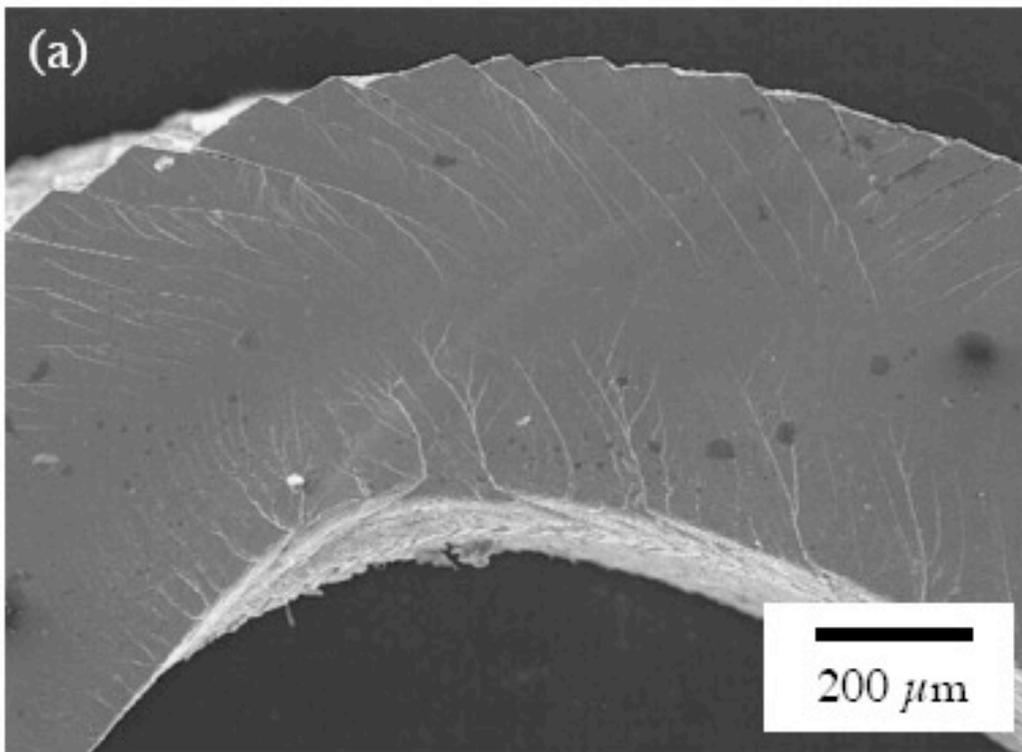
**What is the microscopic description?**

# Amorphous Solids: Example - Metallic Glasses



Metallic glasses exhibit large yield strengths, and large yield strains.

Thanks: U Ramamurty



But exhibit brittle failure accompanied by shear localization (slip steps)..

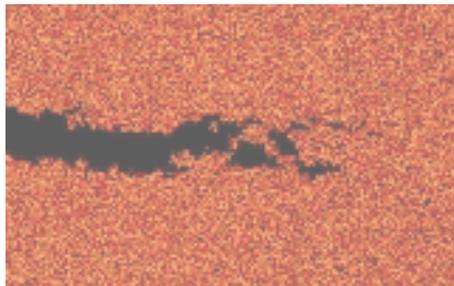
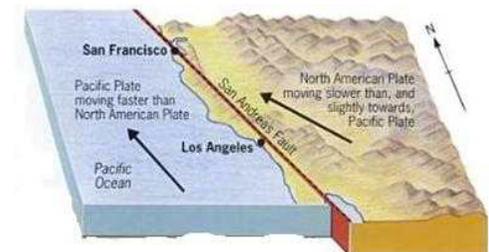
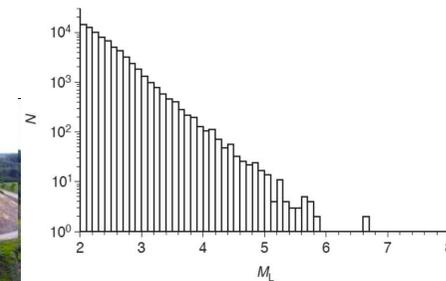
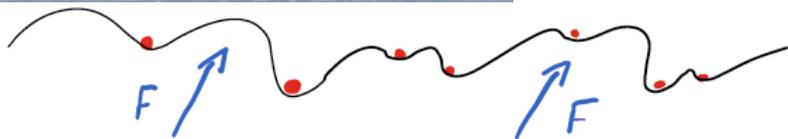
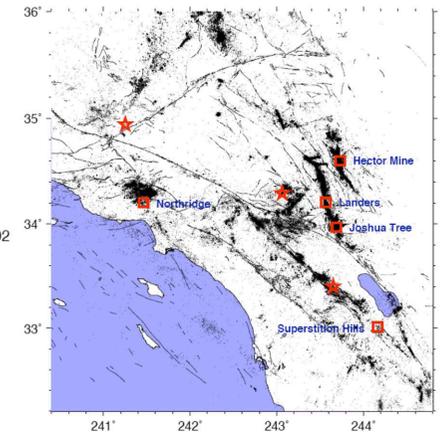
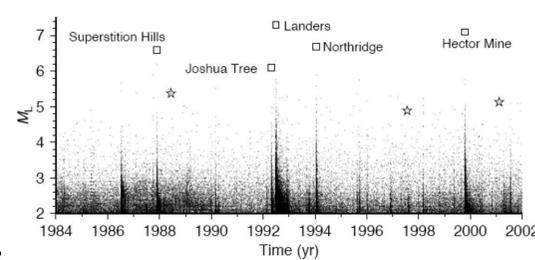
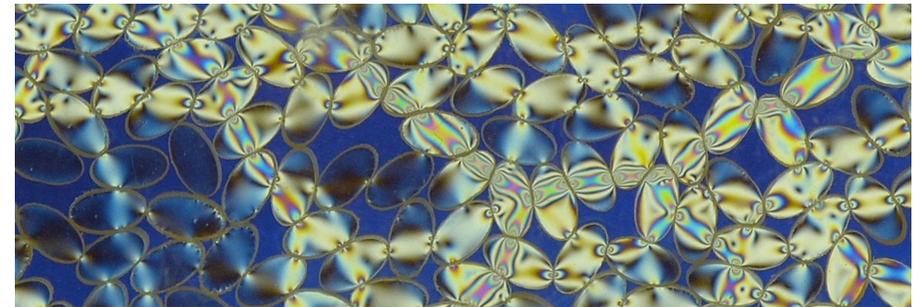
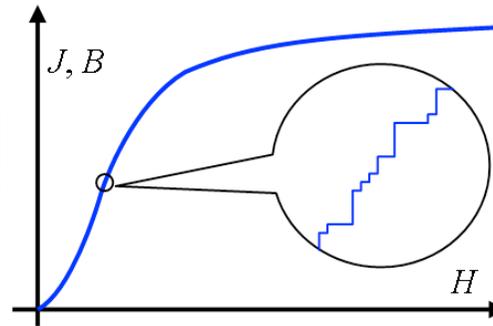
Parameters governing failure, and potentially their control, is critical for their applications.

From: Schuh, Hafnagel, Ramamurty (2007)

# Yielding and Failure in disordered matter: Other examples

Related physics of intermittent, plastic, response in many disordered systems:

- Barkhausen noise in magnets
- Front propagation in disordered media
- Granular matter
- Landslides and avalanches
- Earth quakes

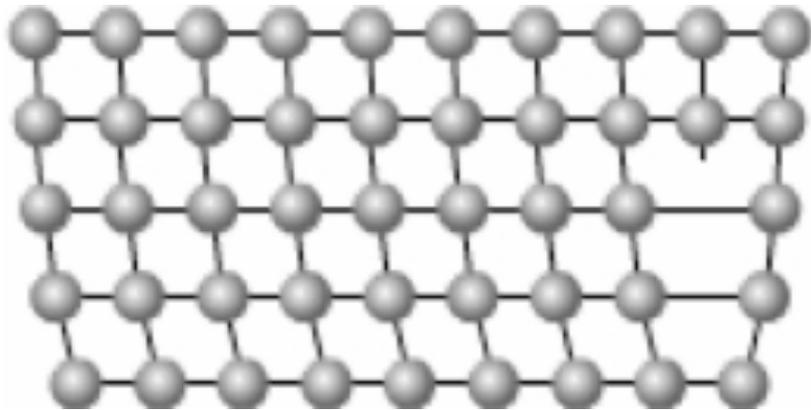


# Mechanical behavior of solids

Response to mechanical stress is a fundamental characteristic of a solid.

Elastic and plastic responses – Elastic moduli, yield stress, strength of materials..

The process by which a crystalline solid responds to stress is understood in terms of the movement of dislocations – defects that are well defined and observable.



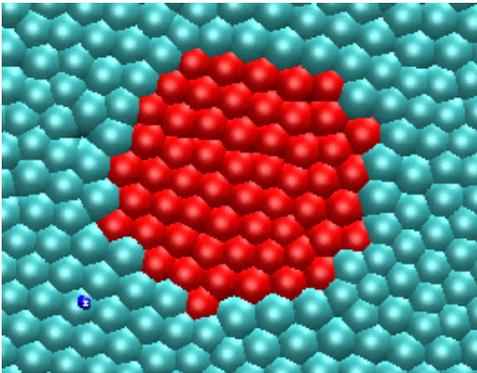
Thanks: Surajit Sengupta



The movies have been provided for teaching purposes by Professor Hideharu Nakashima ([ageigz@mbox.nc.kyushu-u.ac.jp](mailto:ageigz@mbox.nc.kyushu-u.ac.jp)) of Kyushu University, Japan.

# Mechanical behavior of amorphous solids

- **Amorphous solids** do not have regular translational order, and correspondingly no equivalent of dislocation defects.

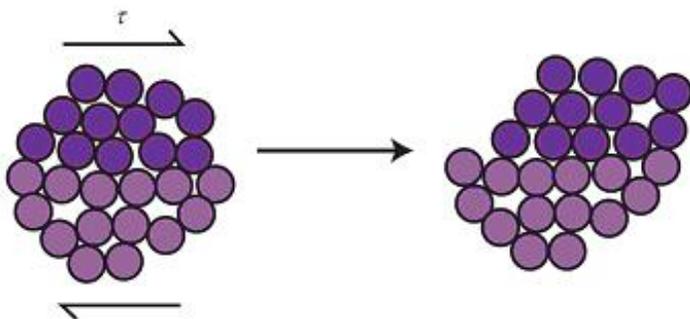


From: Peter Sollich

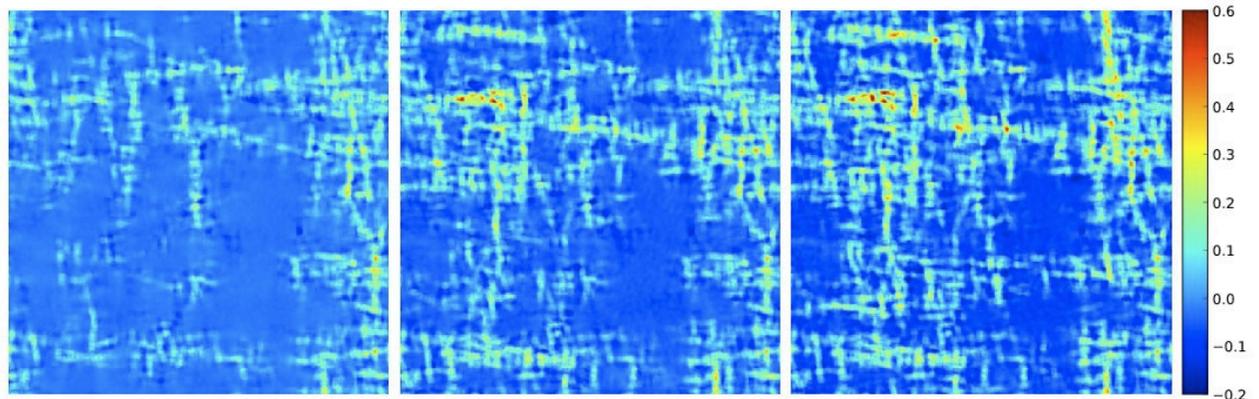
- Deformation mechanism involves localized rearrangements, and avalanches of correlated (or triggered) deformation events.

- What is the description of the transition to the state of plastic flow?

- Wide range of materials, and responses.



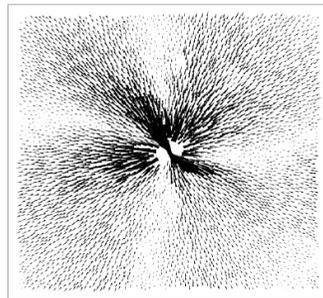
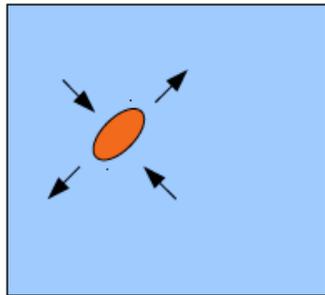
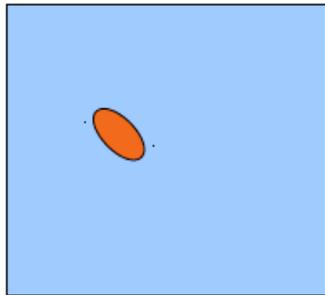
From Schuh et al 2007 [from Argon, Spaepen]



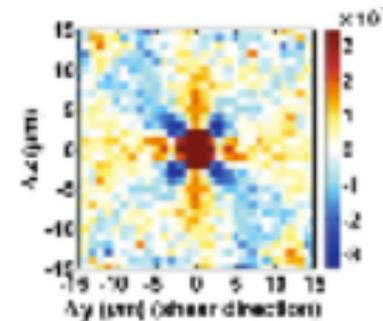
[Lemaitre, Caroli 2009]

# Plastic rearrangements and interactions

A local transformation embedded in an elastic medium leads to long range strain/stresses (Eshelby 1957), interactions:



PRE 74, 016118 (2006)



Jensen et al PRE 2014

$$G(r, r') \sim \frac{\cos(4\theta)}{|r - r'|^d}$$

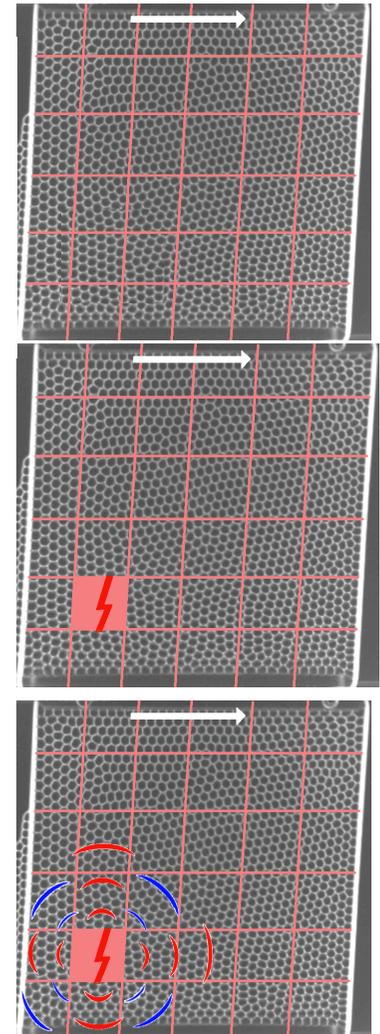
Quadrupolar propagator for stress redistribution.

Observed in simulations and experiments.

Generation of stress field from a local plastic event.

Redistributed stresses trigger other events.

**Avalanches, yielding..**



Barrat

# Crackling Noise: RFIM

Many systems crackle; when pushed slowly, they respond with discrete events of a wide range of sizes.

*Examples:* Earthquakes [motion of tectonic plates], candy wrapper [on crumpling] or **a magnetic material under changing field.**

Field on magnetic domain under an external field ( $H(t)$ ):

$$H(t) + \sum_j S_j + h_i$$

$h_i$  : a random field for each domain with variance "R".

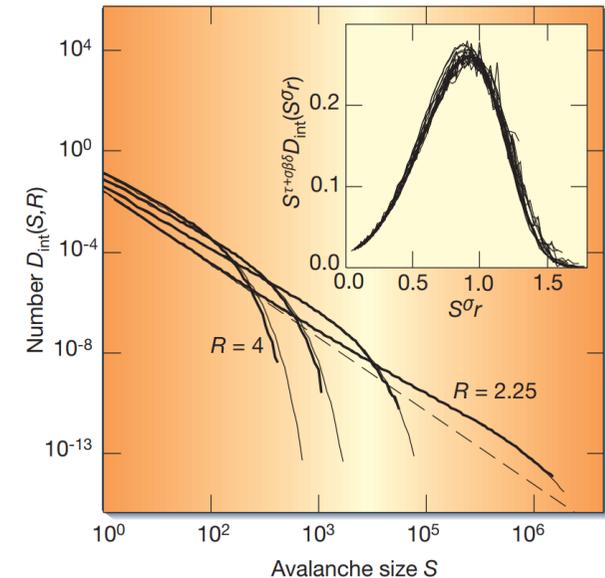
Flipping of spins can initiate avalanches.

If "R" is large (or  $h \gg J$ ) most domains flip independently: all the avalanches are small (small popping noise).

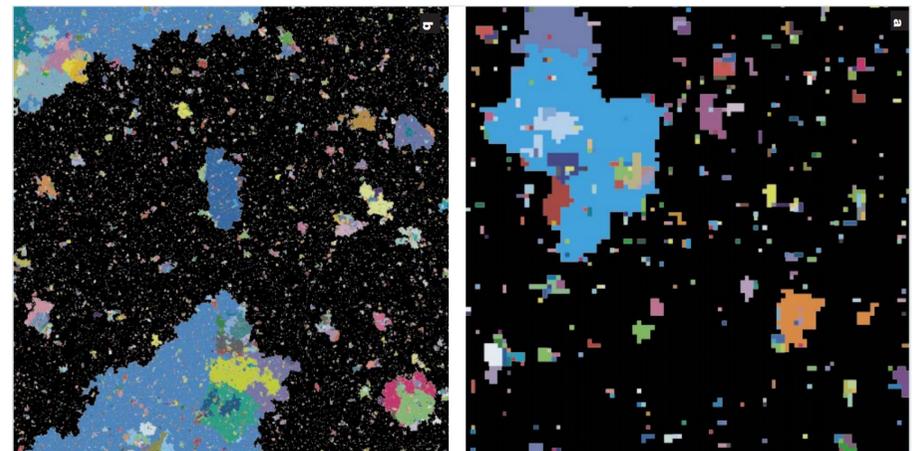
For sufficiently large  $H(t)$ : large scale avalanches.

**A R – H phase diagram with a critical point.**

Sethna et al Nat Rev. 2001



Cross-sections of the avalanches during the magnetization, various color: avalanches



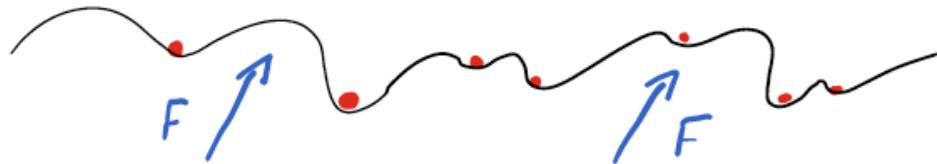
# Depinning

Depinning of an interface moving in an interface (fluid, etc) a problem of broad interest.

Pinning arises from disorder. Depinning as a result of external drive.

Depinning model

$$\partial_t h(x, t) = \nu \partial_x^2 h - \partial_h V(x, h(x, t)) + f$$



- Pinning is induced by quenched disorder ( $V$ ).
- Stops the motion of driven elastic manifolds for applied forces  $\mathbf{f}$  below a critical value  $f_c$ .
- The avalanche size defined as the area swept by the interface in a burst.
- Avalanche size diverges at depinning.

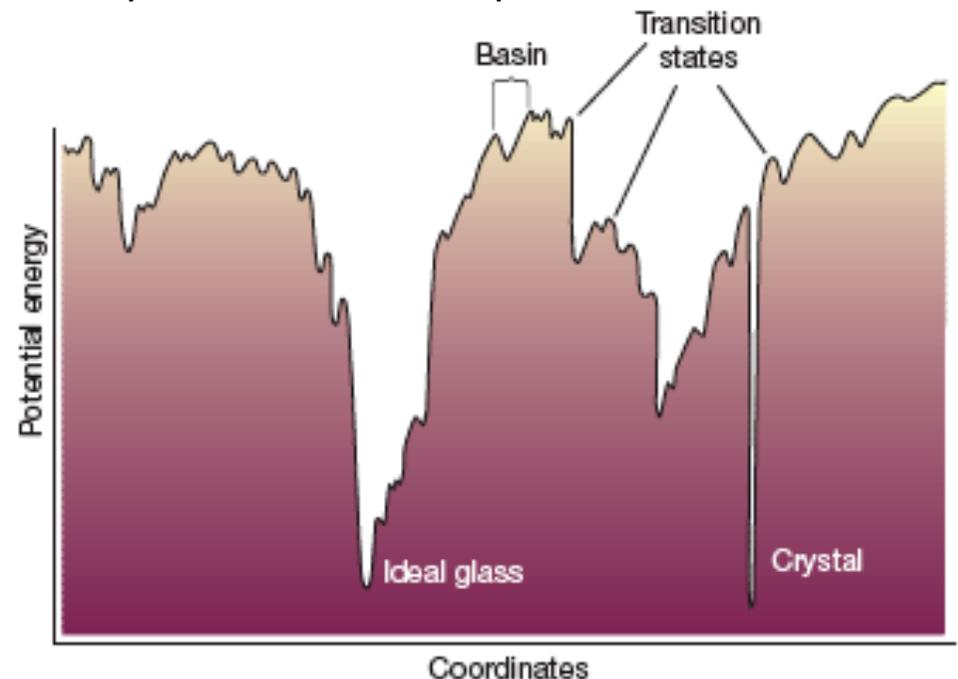
# What we do: Approach and Questions

- Mechanical response to affine, shear deformations studied by analyzing the modification of the landscape, and transitions between minima, through the **AQS** (Athermal Quasi Static) protocol.
- Cyclic shear deformation of model glasses.
- Amplitude of shear deformation the only control parameter.
- Limitation: No information on the interplay between relaxation processes and instabilities triggered by deformation. But useful insights obtained.
- Simulations done also at finite shear rate and temperatures for comparison.

Focus for this talk:

The Yielding Transition.

Shear banding.



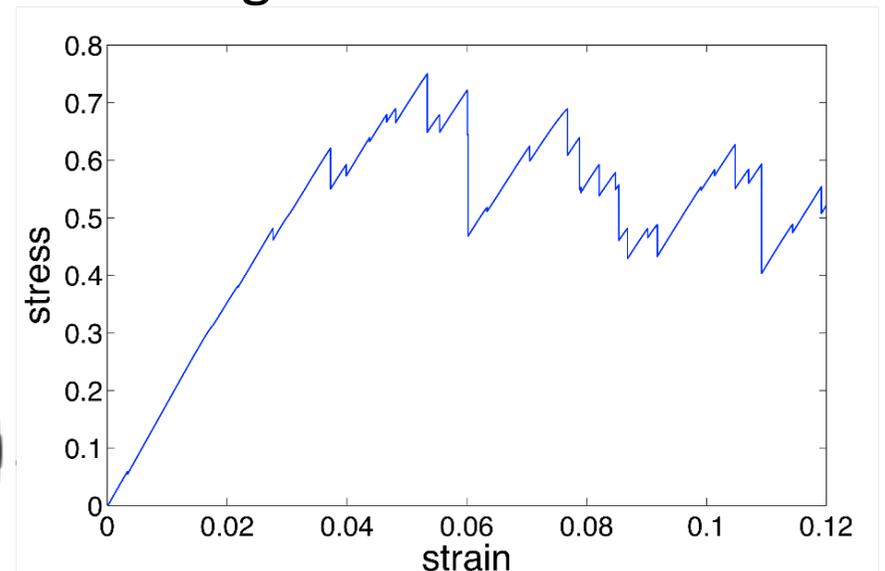
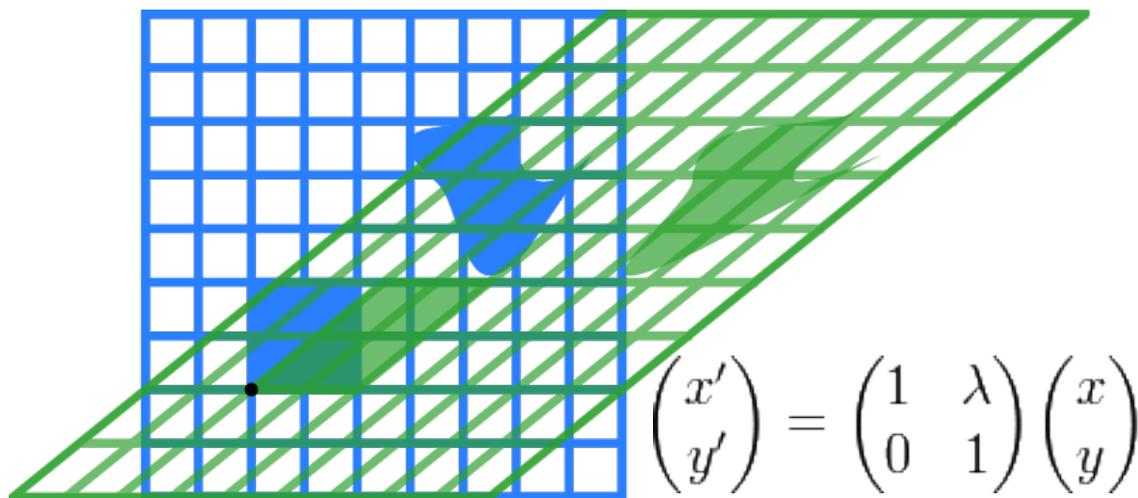
# Questions/Issues

- The character of **avalanches** before yielding and after yielding.
- The nature of the **yielding transition**. Is it a sharply defined transition?
- Strain localization above yielding.
- Relaxation/annealing effects under cyclic loading.

Computer simulations of deformation of model glasses in search of some answers

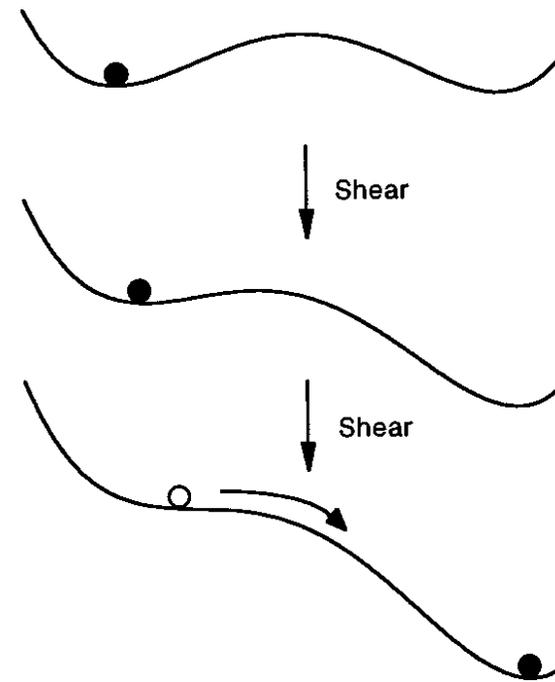
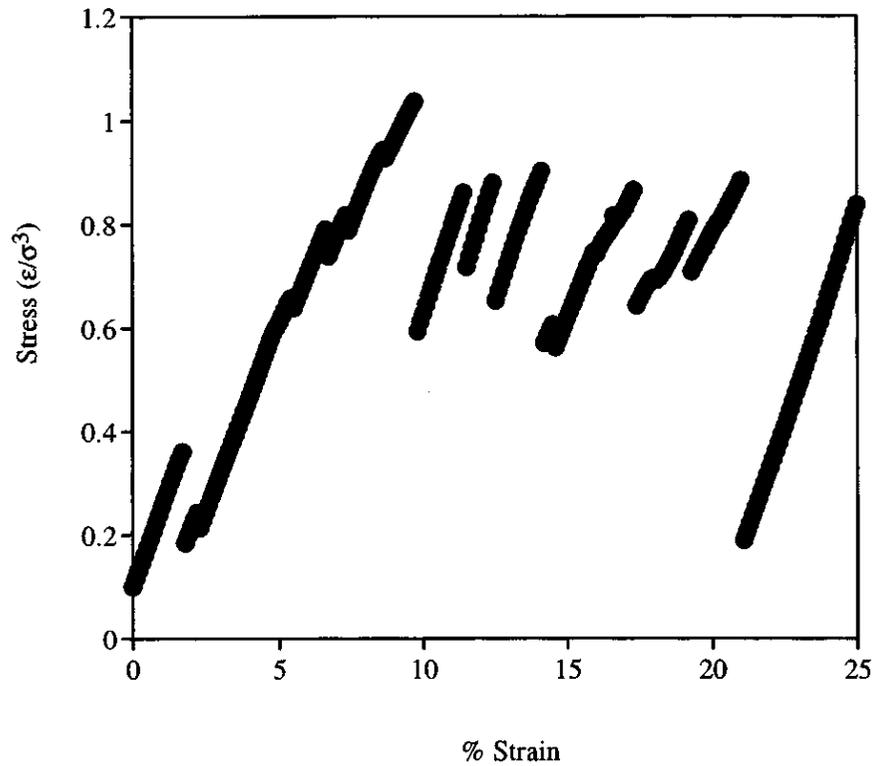
# Athermal Quasi Static Deformation

1. Subject energy minimum structures to shear deformation.
2. Minimize the resulting deformed structure subject to suitable (Lees-Edwards) boundary conditions.
3. Deformation strain increased quasi-statically.
4. The procedure produces a sequence of configurations that are always energy minima.
5. Continuous change of energies interrupted by discontinuous change.
6. Discontinuous changes correspond to rearrangements.



# Mechanical deformation of model glasses

Discontinuous changes in stress from destabilization of minima.

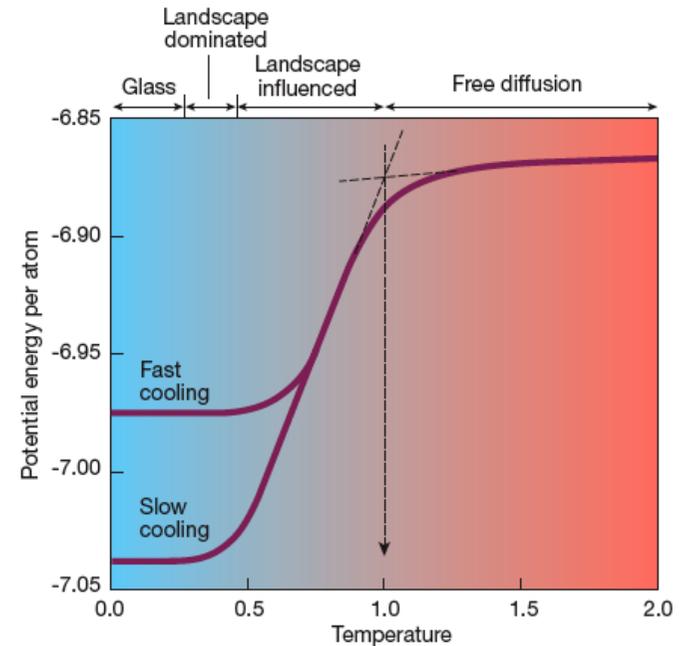
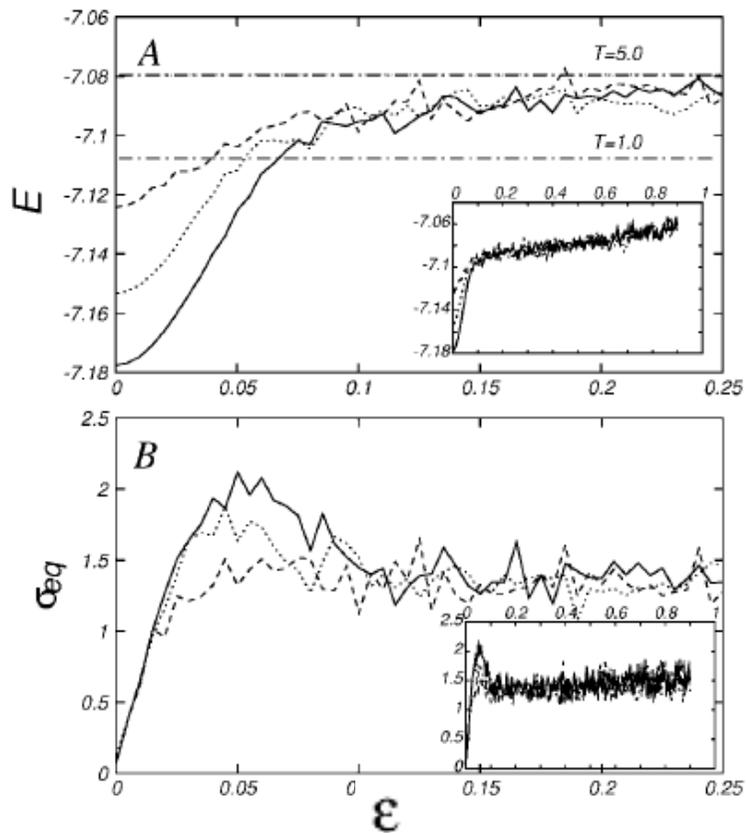


Malandro and Lacks 1999

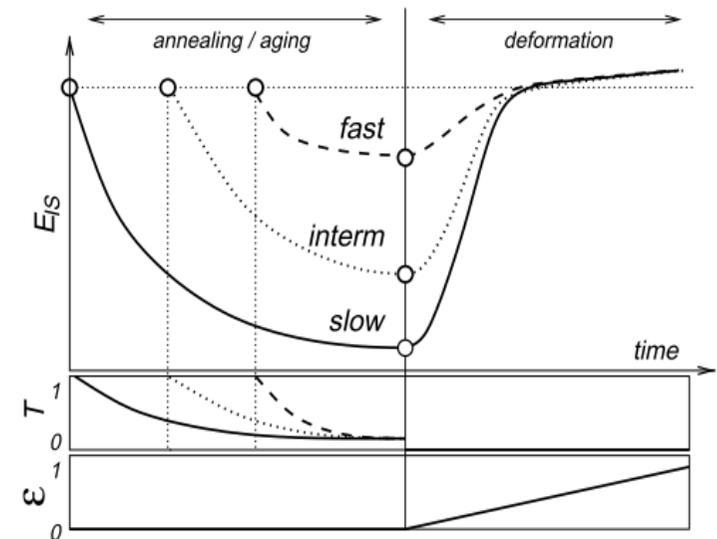
# Mechanical deformation leading to yielding..

Early studies:

1. Deeper energy minima are sampled at lower temperatures.
2. Upon shearing, the energy of the minima rises to those at very high temperatures.
3. Deformation induced 'rejuvenation'

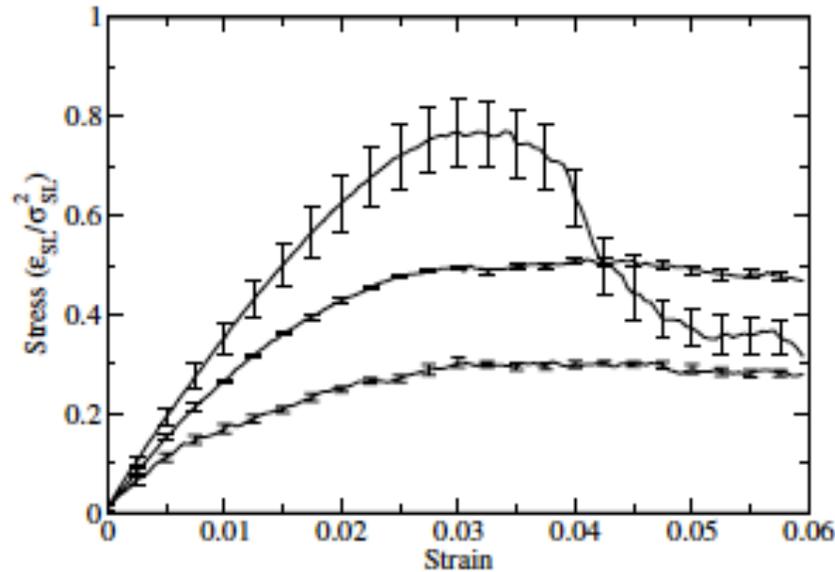


Sastry et al Nature 1998



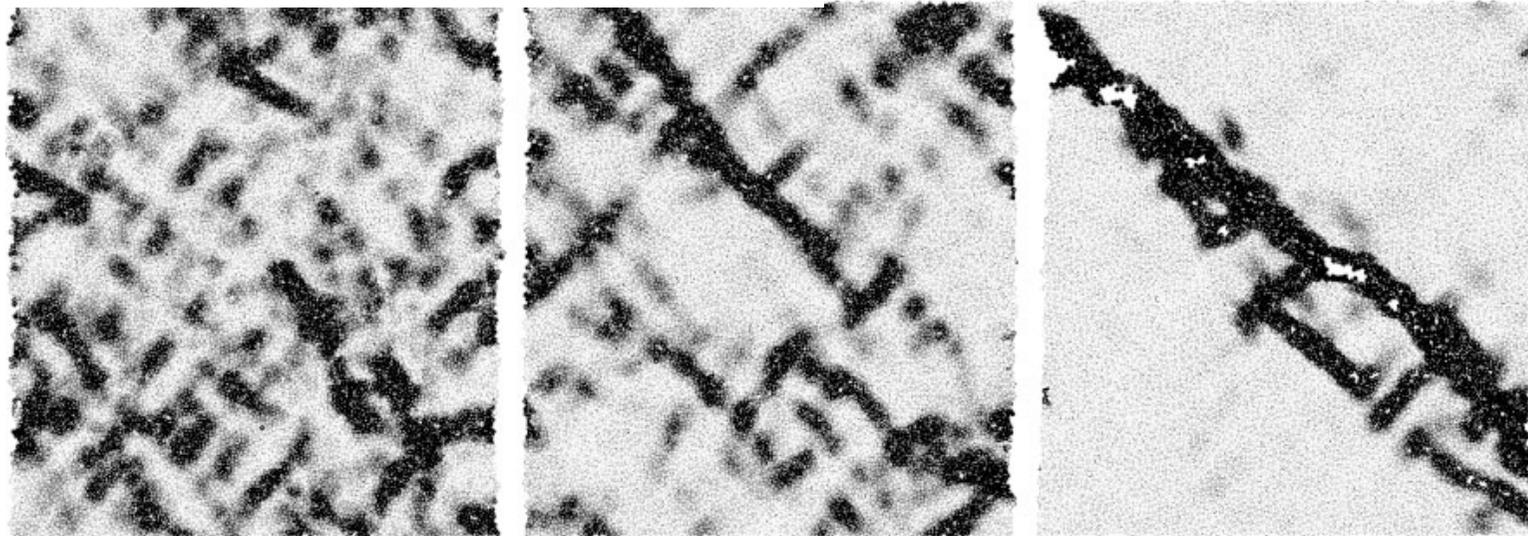
Utz, Debenedetti and Stillinger and Utz, 2000

# Strain Localization



Strain localisation depends on degree of structural relaxation prior to shear deformation.

Shi and Falk (2005)

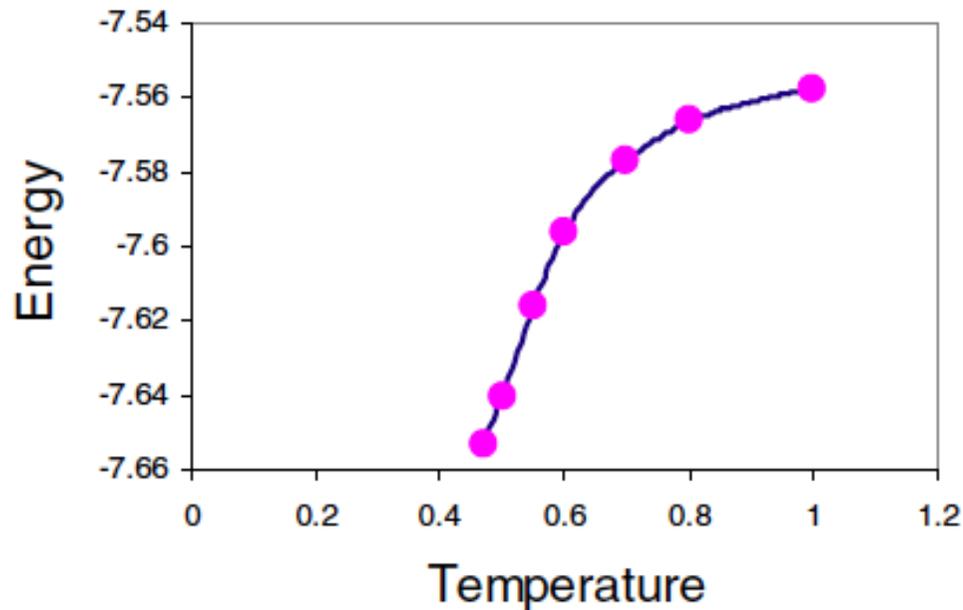


Related: Transient shear bands in flow (Srivastav et al 2016, Vasisht et al 2017)

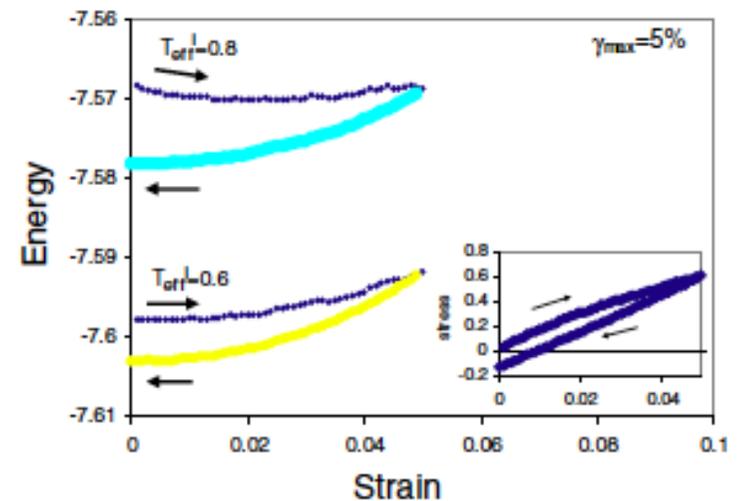
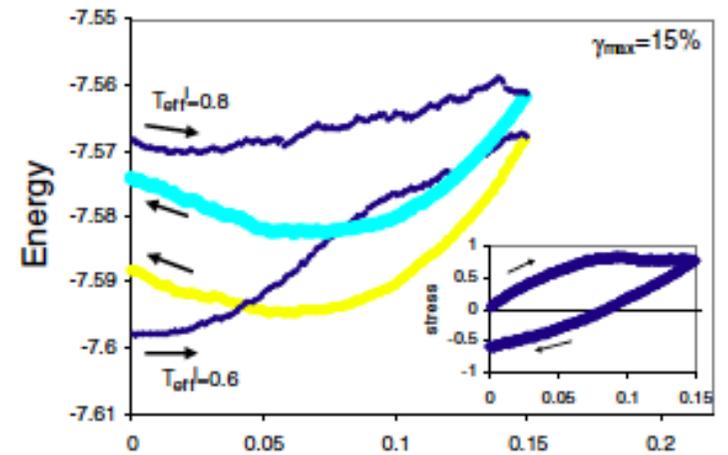
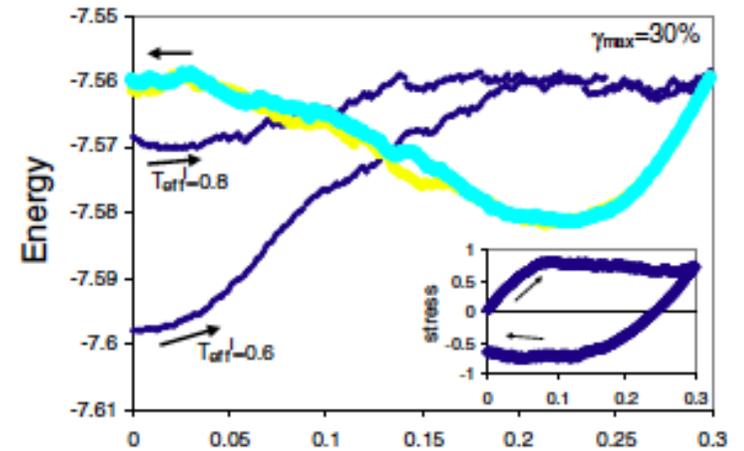
Somewhat more complicated picture when a forward and back cycle of strain is applied:

The nature of change (aging or rejuvenation) depends on amplitude of strain and the initial state..

What happens if this cycle is repeated?



Lacks and Osborne, PRL 2004.



# Cyclic Deformation: Schematic

Shear deformation modifies the potential energy landscape and destabilizes the system, eventually leading to irreversible rearrangements.



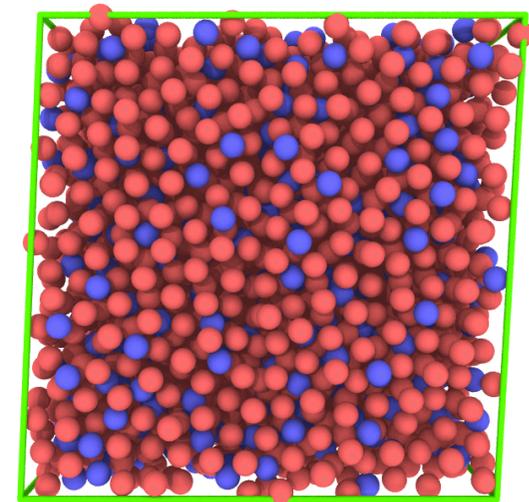
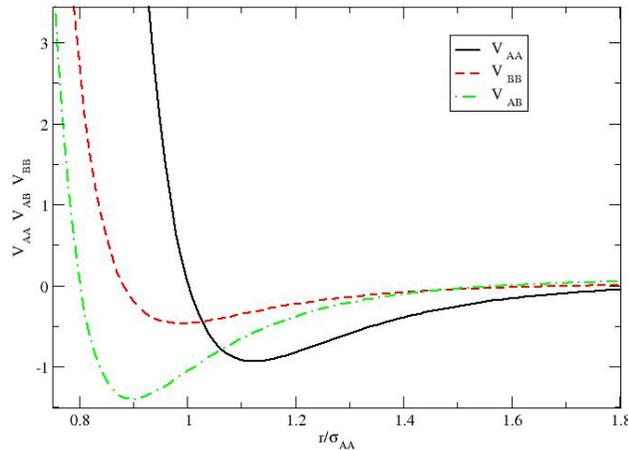
How does such deformation modify the properties of the glasses?

# Simulations of oscillatory strained binary Lennard-Jones (LJ) solids

Kob-Andersen binary glass forming model liquid.

(Constant Volume AQS)

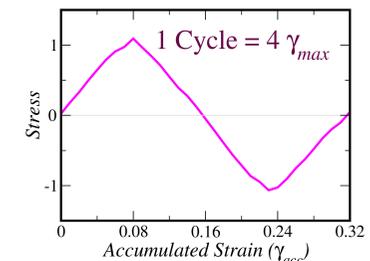
$$\phi_{ij}(r) = 4\epsilon_{ij} \left( \frac{\sigma_{ij}^{12}}{r^{12}} - \frac{\sigma_{ij}^6}{r^6} \right)$$



Different system sizes: 2000, 4000, 8000, 16000, 32000, 64000, 128000 and 256000.  $0 \rightarrow \gamma_{max} \rightarrow 0 \rightarrow -\gamma_{max} \rightarrow 0$

Two starting temperatures: **1** (poorly) & **0.466** (well annealed glass)

Cyclic shear for range of  $\gamma_{max}$  values with strain step  $d\gamma = 2 \times 10^{-4}$ .



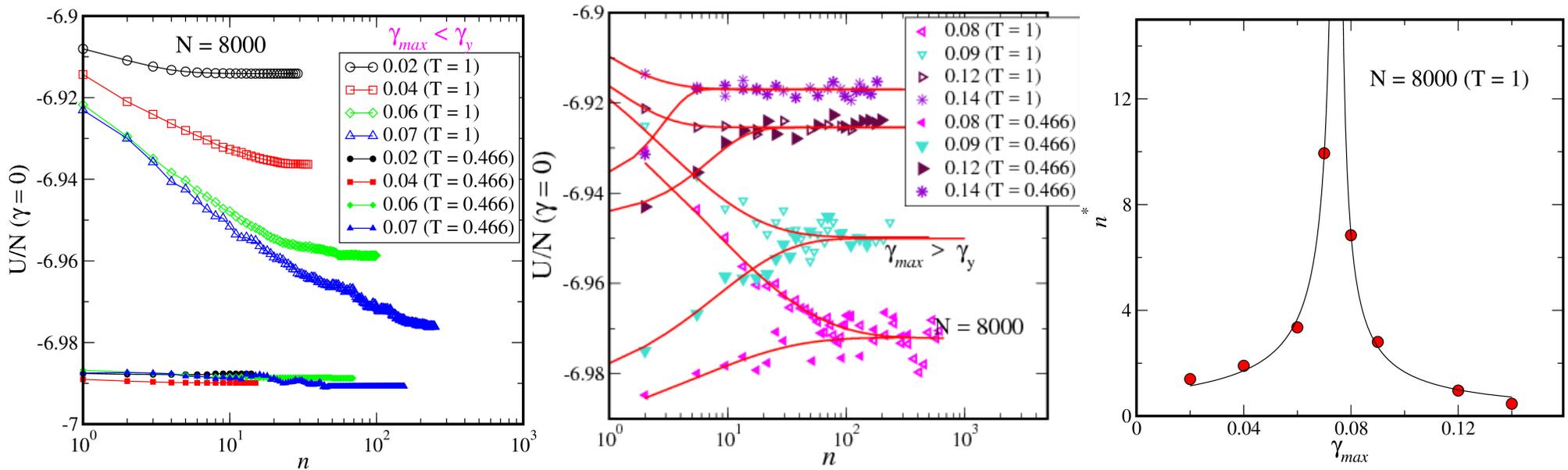
Each  $\gamma = 0$  configuration labeled with the *accumulated* strain  $\gamma_{acc} = \sum d\gamma$

These **stroboscopic** configurations were used to compute various quantities, i.e. energy, MSD etc ...

# Potential Energy vs. Cycle Number

The potential energy per particle reaches a plateau that

- (a) Depends on  $\gamma_{max}$  **only** at large values of  $\gamma_{max}$ .
- (b) Depends on  $\gamma_{max}$  and initial state for small  $\gamma_{max}$ .



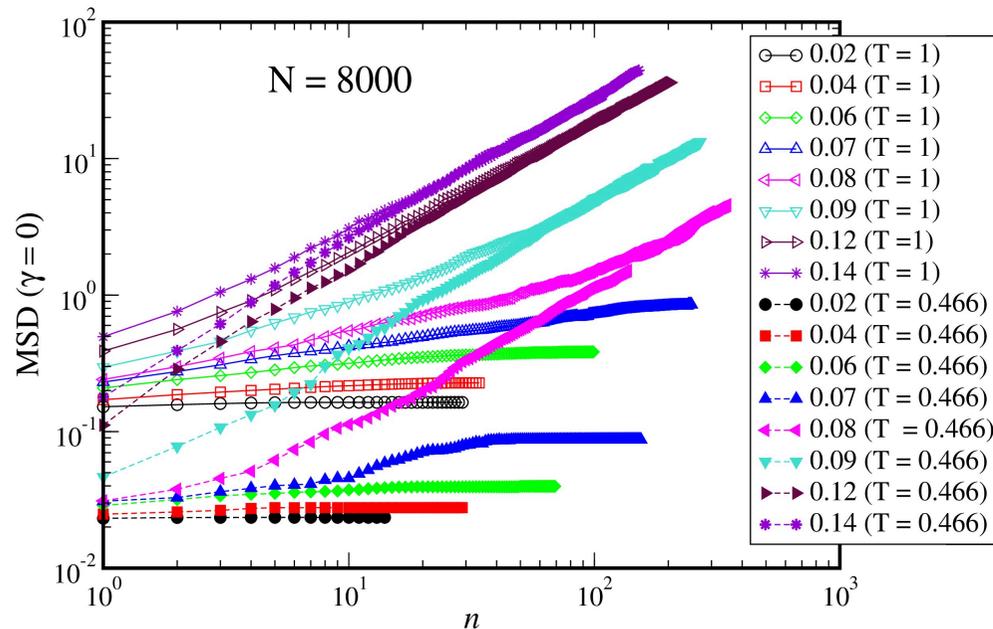
- Aging/rejuvenation depends on strain amplitude and initial annealing on the glasses.
- Relaxation to the steady state **becomes more sluggish as  $\gamma_y$  is approached.**

**Change in behavior across a critical strain amplitude  $\gamma_c$**

# Mean Squared Displacement vs. Cycle #: Diffusion Coefficient

Depending on  $\gamma_{\max}$  systems are either diffusive or non-diffusive.

In the diffusive regime, asymptotic slopes depend only on  $\gamma_{\max}$ .



$$\text{MSD} = D\gamma_{acc}^*$$

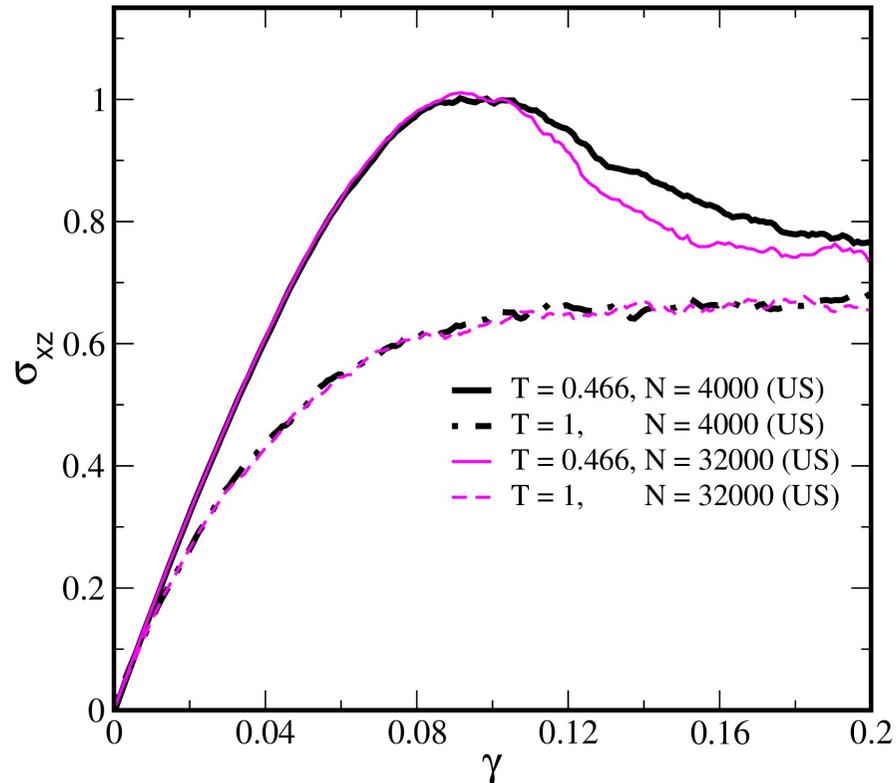
The diffusion coefficient vanishes below a finite value of  $\gamma_{\max}$

[How? Will be addressed later]

Critical  $\gamma_{\max}$  a function of system size... but approach finite value asymptotically.

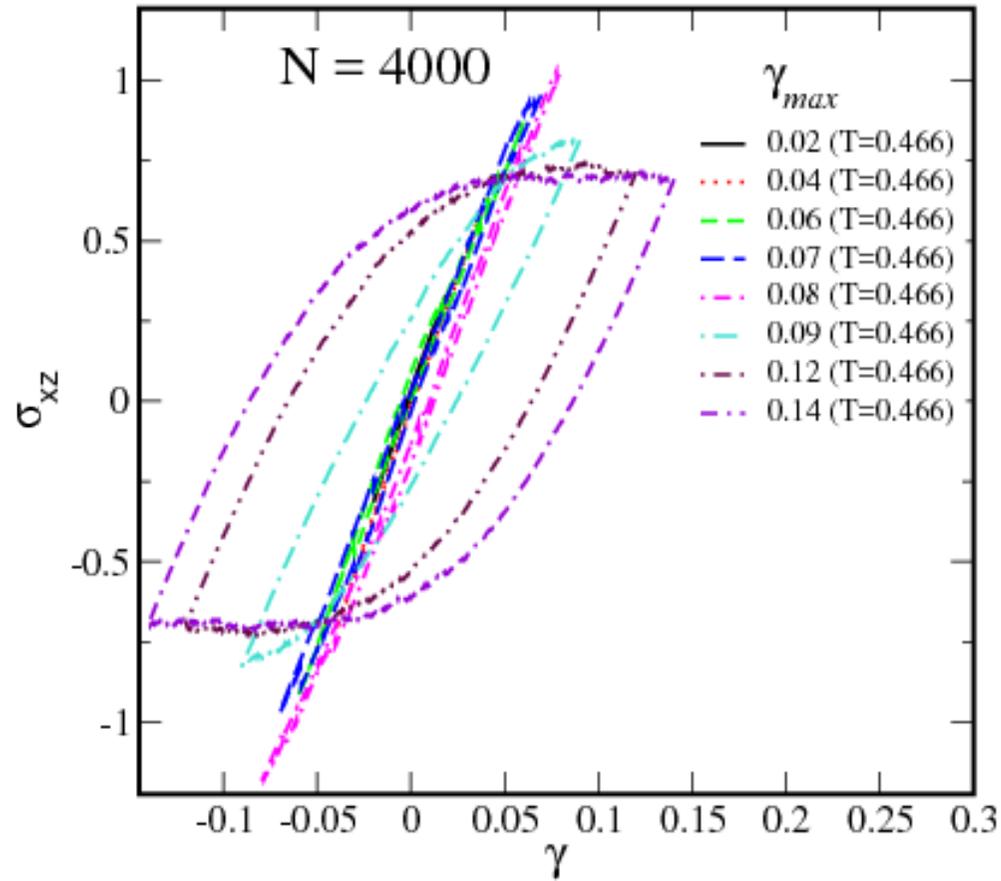
**Non-equilibrium transition from localized to diffusive regimes!**

# Yielding Uniform Shear



There is **no sharp point** to identify from uniform shearing.  
Significant **sample** (annealing) **dependence**.

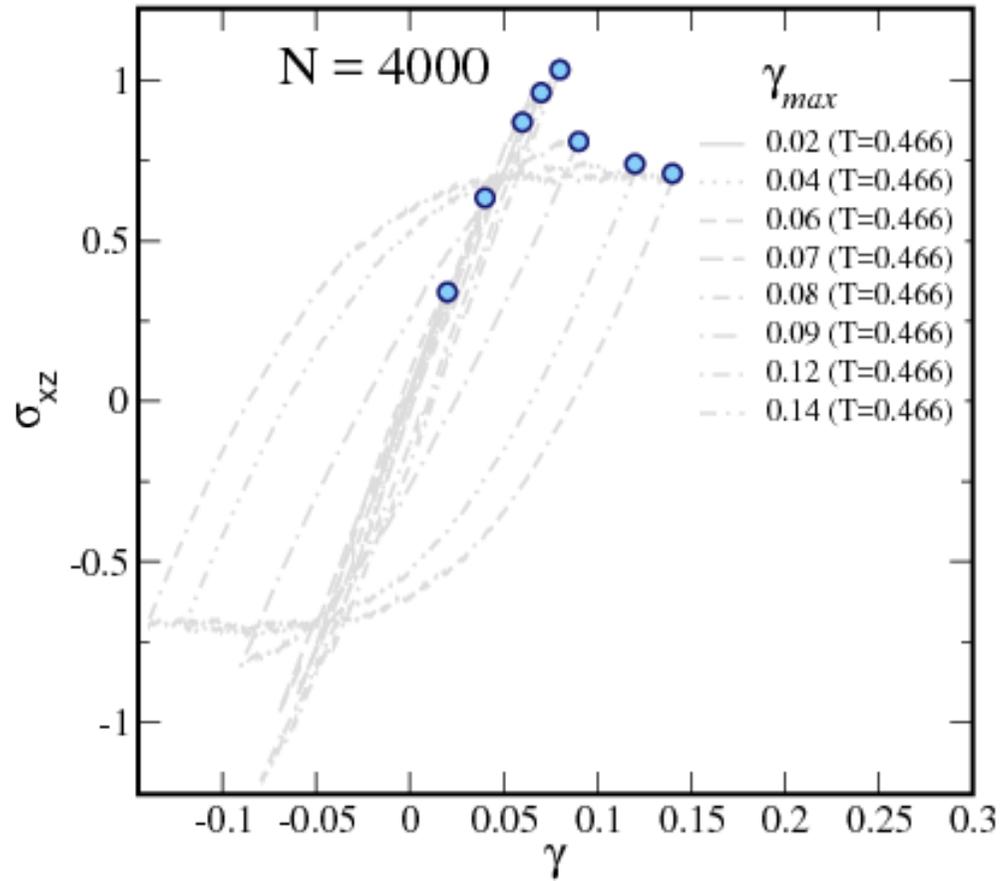
# Yielding: Oscillatory shear



Oscillatory shear provides a better characterization of the transition.

We focus on maximum stress during cyclic shear..

# Yielding: Oscillatory shear



Oscillatory shear provides a better characterization of the transition.

We focus on maximum stress during cyclic shear..

In **uniform shearing** (US), there is a clear difference in the stress-strain curves depending on the prior annealing of the glasses.

For **cyclic shear** (CS) qualitatively the same behavior regardless of initial sample conditions.

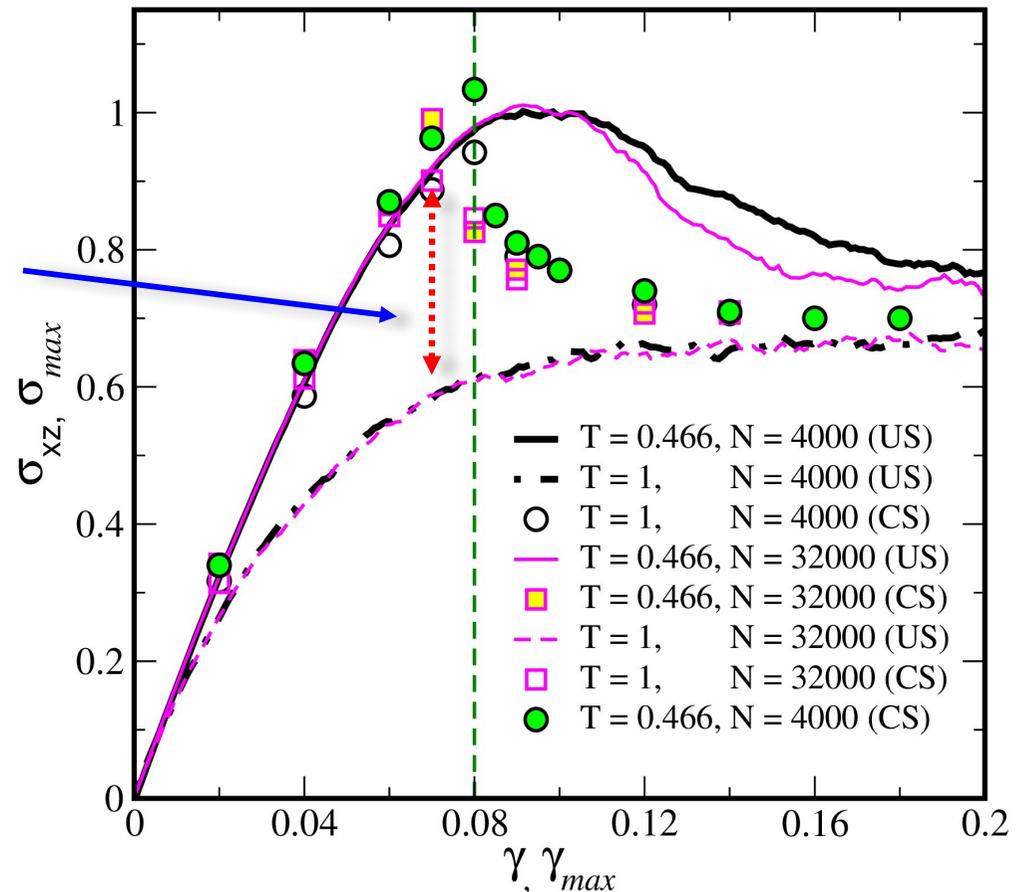
**A sharply defined, discontinuous yielding transition.**

**work hardening !!**

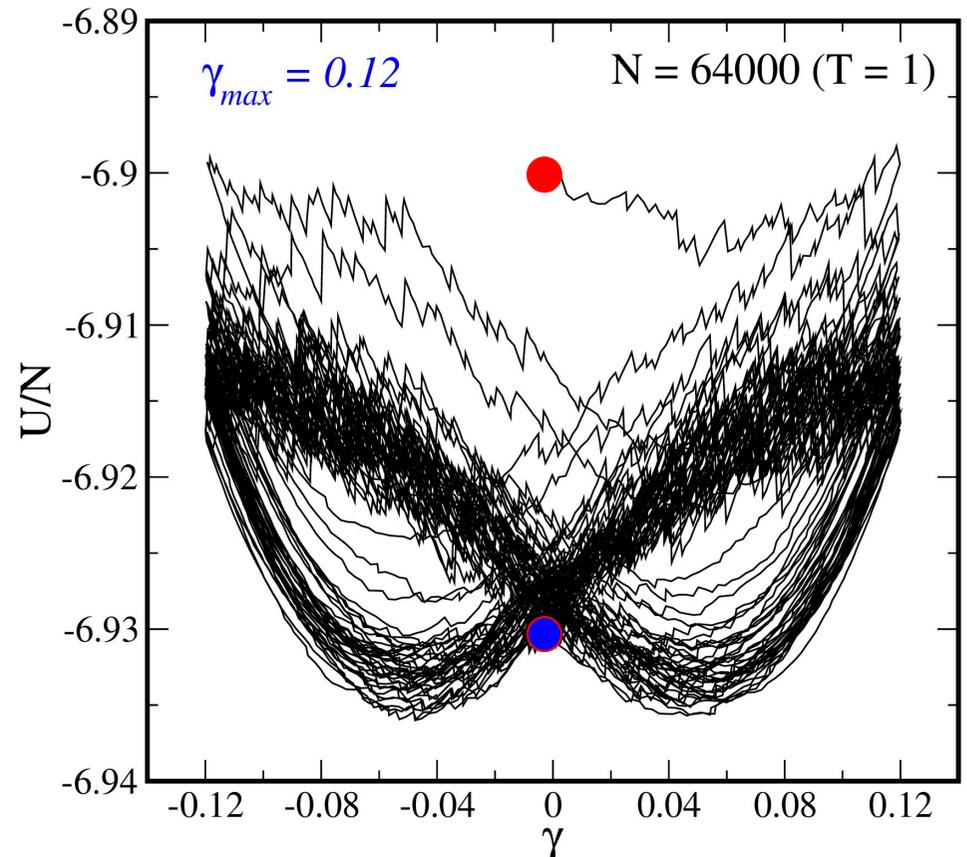
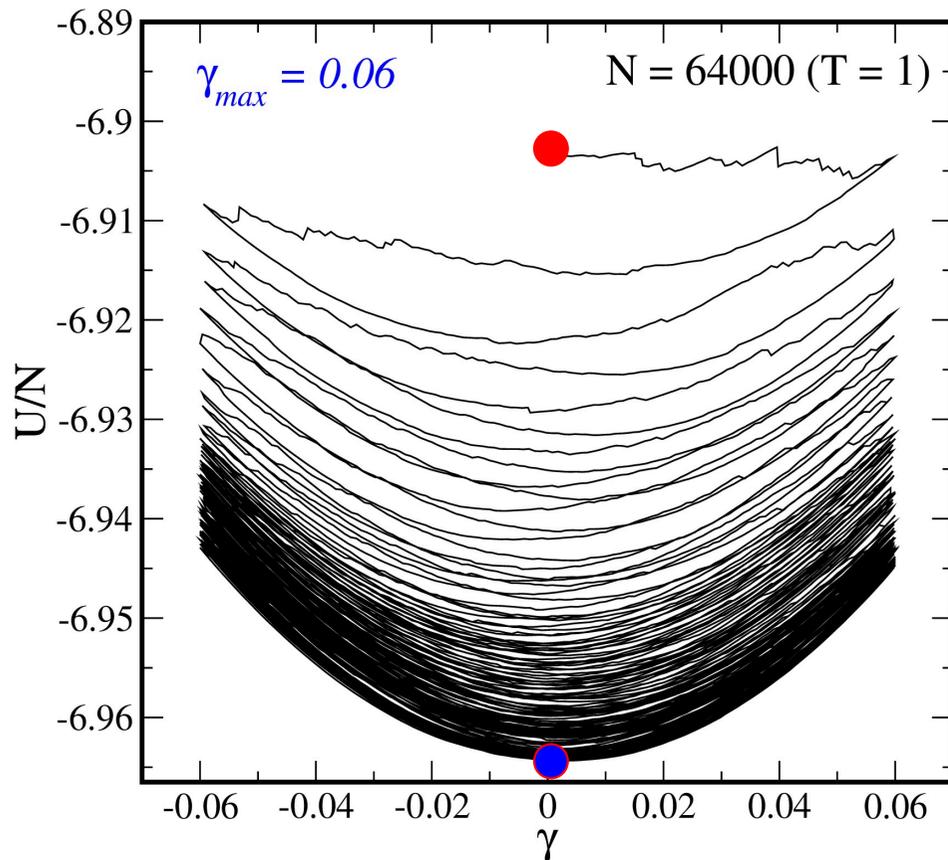
**Oscillatory shear provides clear signature of yielding.**

**Pointers for thermo-mechanical processing of metallic glasses?**

Ref: Jaiswal et al 2016, Parisi et al 2017, Osawa et al 2018, Popovic et al 2018..

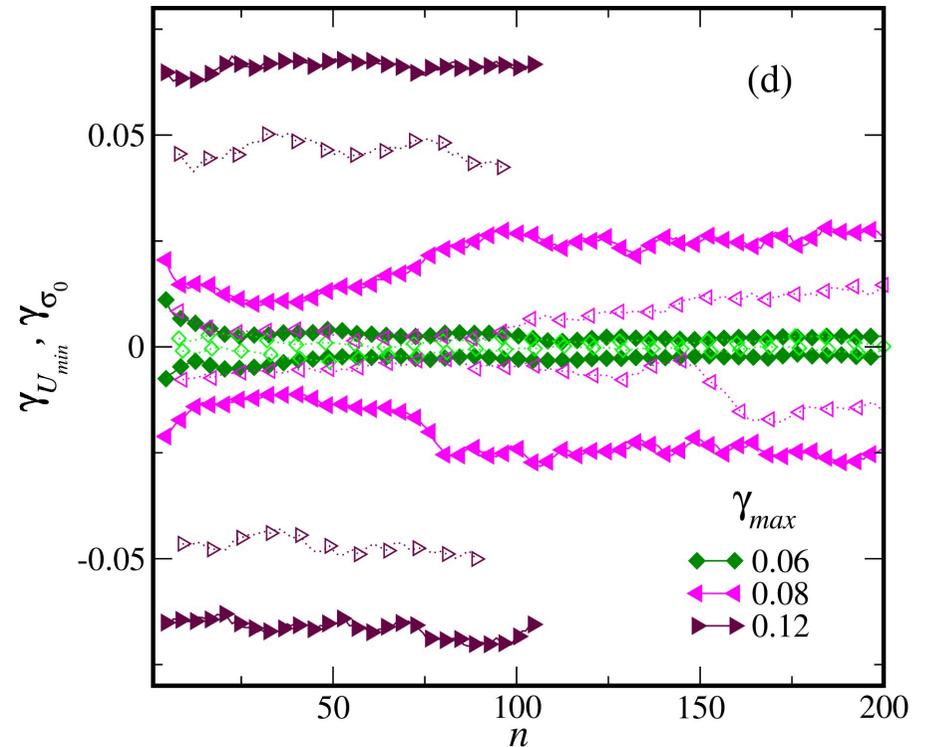
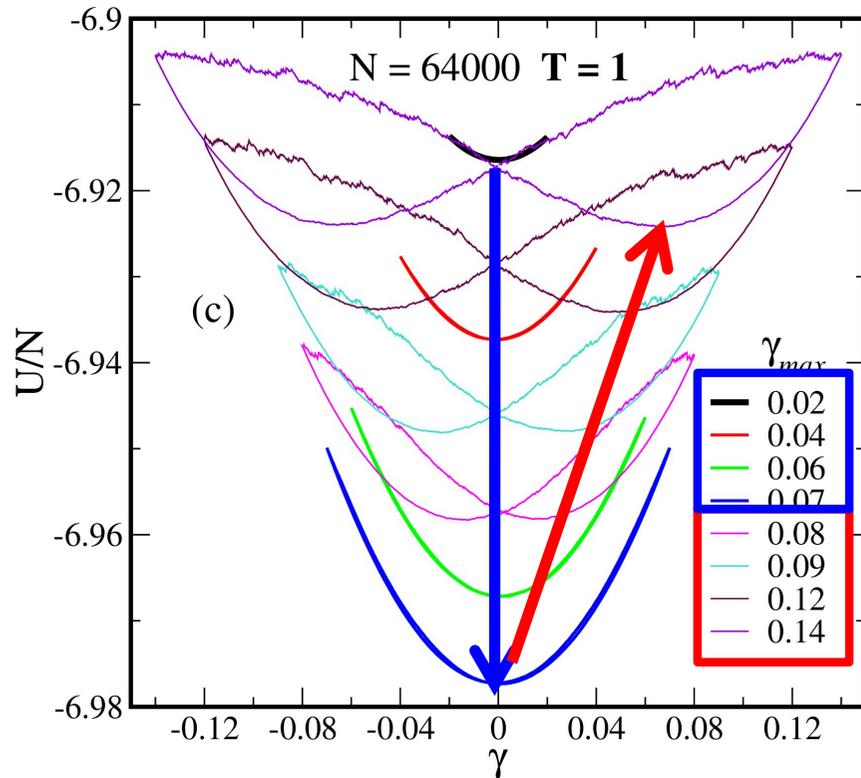


# Evolution of energy with cycles



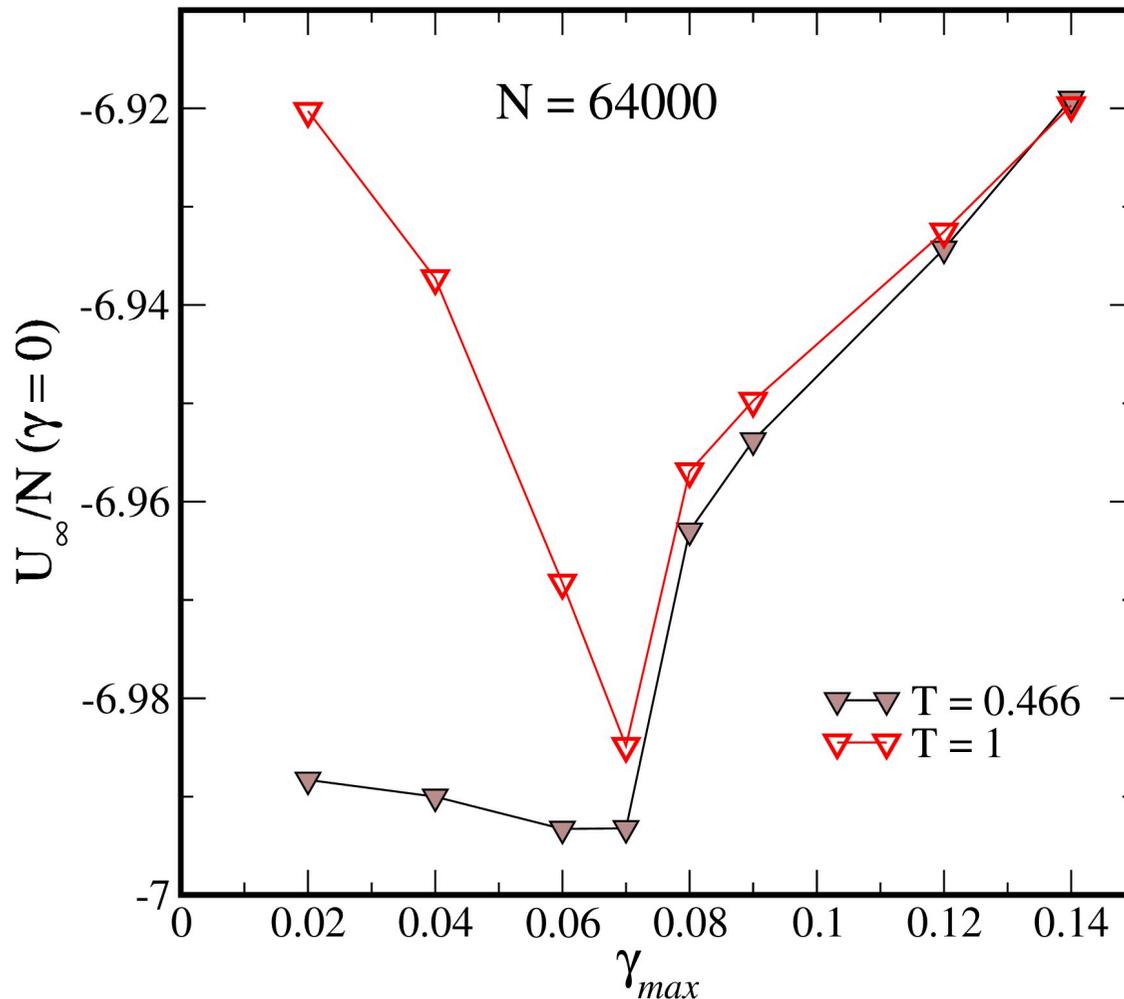
- For  $\gamma_{max} = 0.06$  (i.e.  $\gamma_{max} < \gamma_y$ ) the energy approaches to a single minimum at  $\gamma = 0$ , but
- Energy bifurcates into two minima (for  $\gamma_{max} = 0.12$  (i.e.  $\gamma_{max} > \gamma_y$ )) at finite strain.

# Evolution of energy with cycles



- Energy vs. strain in the steady states, displaying a bifurcation in the strain corresponding to minima in energy at the yielding transition between  $\gamma_{max} = 0.07$  and  $0.08$ .
- Evolution of strain values for energy minimum and zero-stress states indicate transition.

# Asymptotic energy vs. strain amplitude

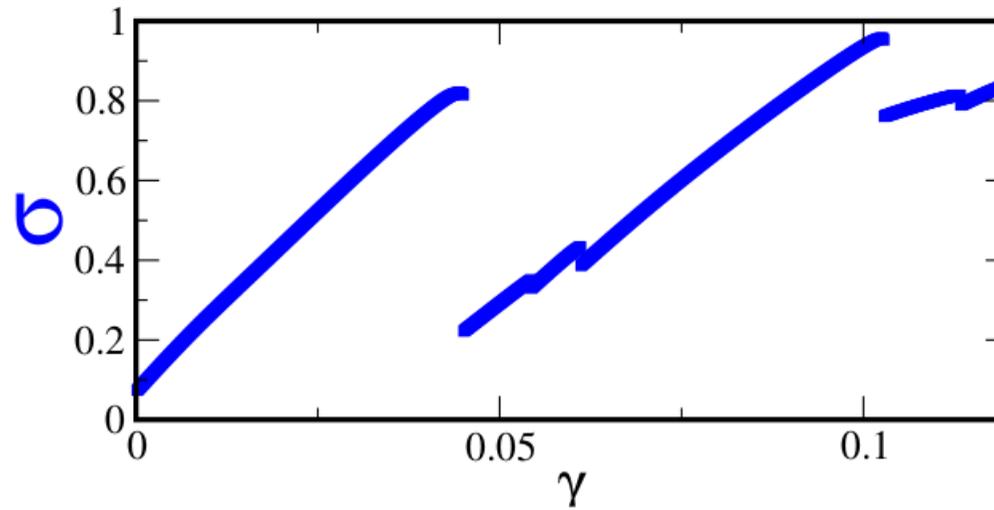


Better annealing close to the transition.

What is the nature of the energy increase beyond the yielding transition?

Energies of stroboscopic configurations decrease with  $\gamma_{max}$  till the yield strain is reached, after which they increase with  $\gamma_{max}$ .

# Avalanches

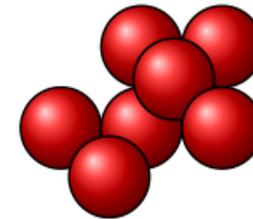


Plastic Instabilities:

Stress Released

Energy Dissipated

Clusters of  
Active Particles

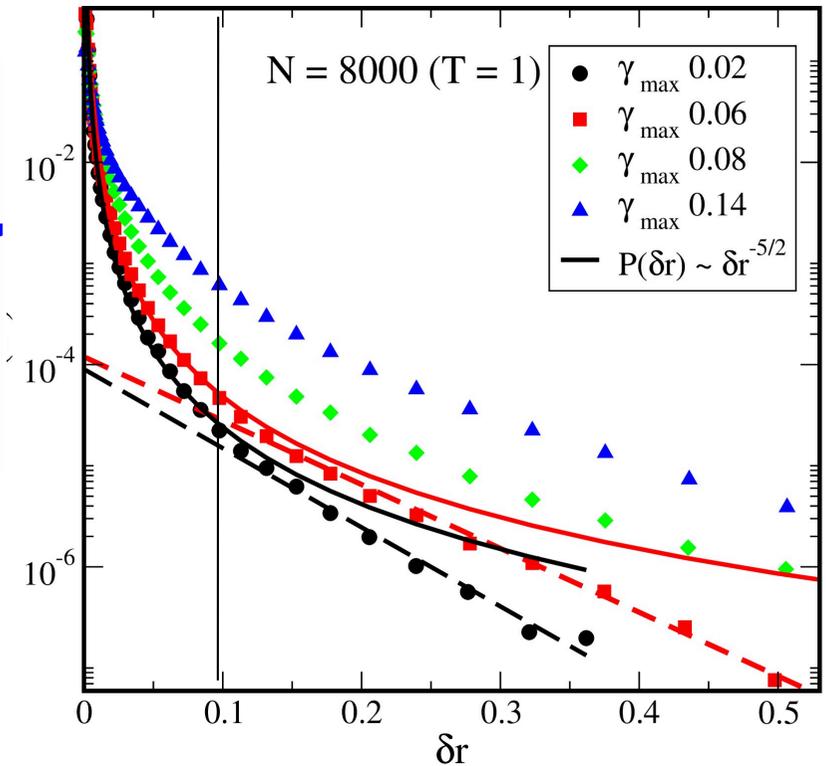
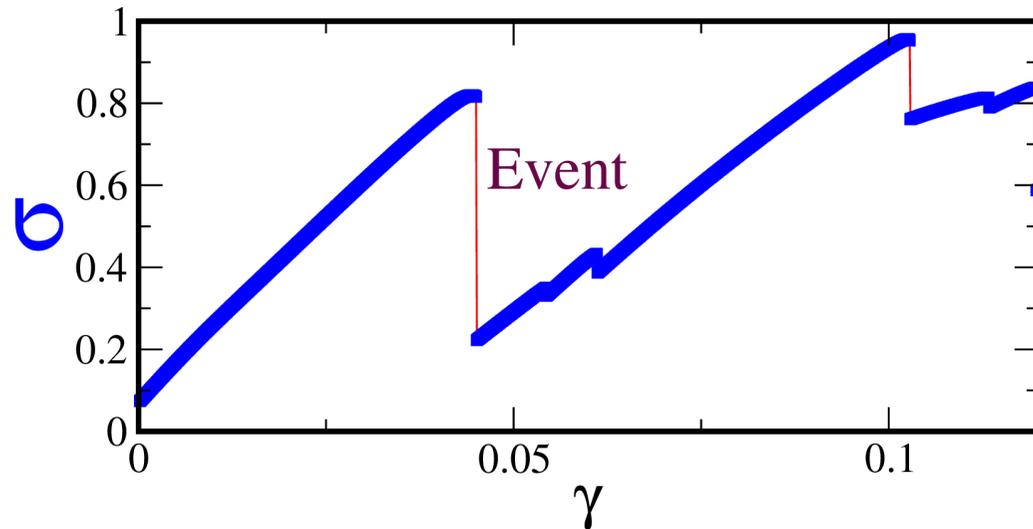


# Avalanches

Definitions:

**Active particles:** Particle is active if moved beyond  $0.1\sigma_{AA}$  cutoff in the event.

**Clusters:** Active particles with connectivity of  $1.4\sigma_{AA}$  (the first coordination shell)

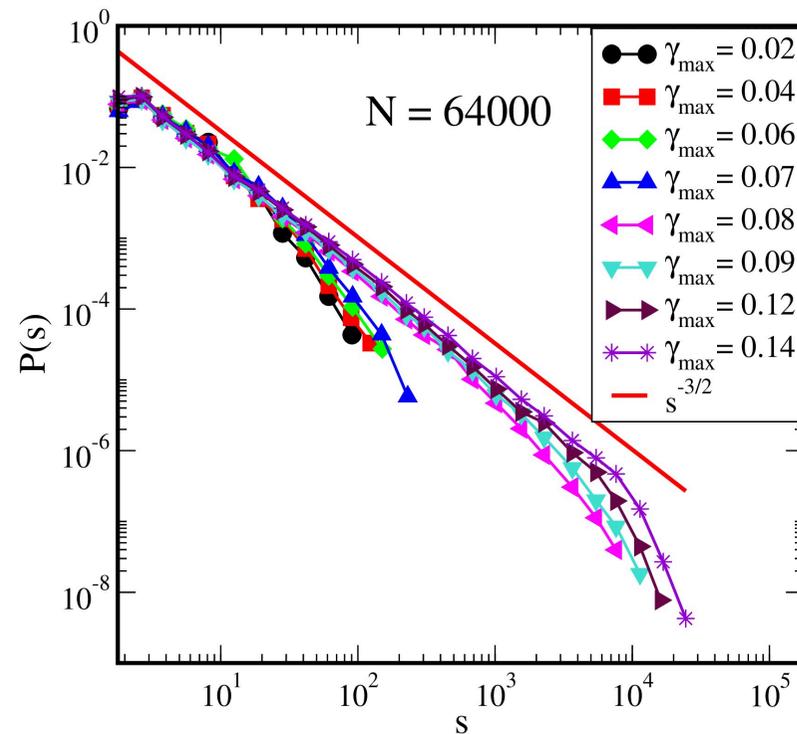
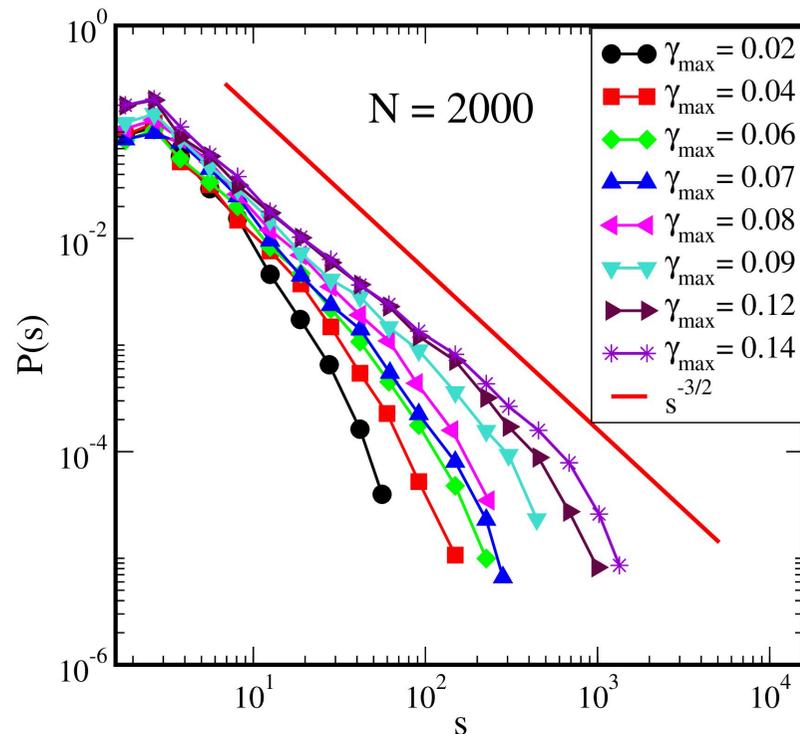


We study statistics of:

(a) **Cluster sizes** (b) **Energy drops**

How do events (avalanches) evolve with amplitude of strain?  
Do they capture the approach to the yielding transition?

# Distributions of Avalanche Sizes



- For  $N = 2000$  displaying a power law with a cutoff that grows with  $\gamma_{\max}$  but **does not indicate sharp changes** at yielding.
- For  $N = 64000$  displaying a **sharp increase** in the cutoff size across the yielding transition. **Power law regime with power 3/2.**

**System size effects are important!!**

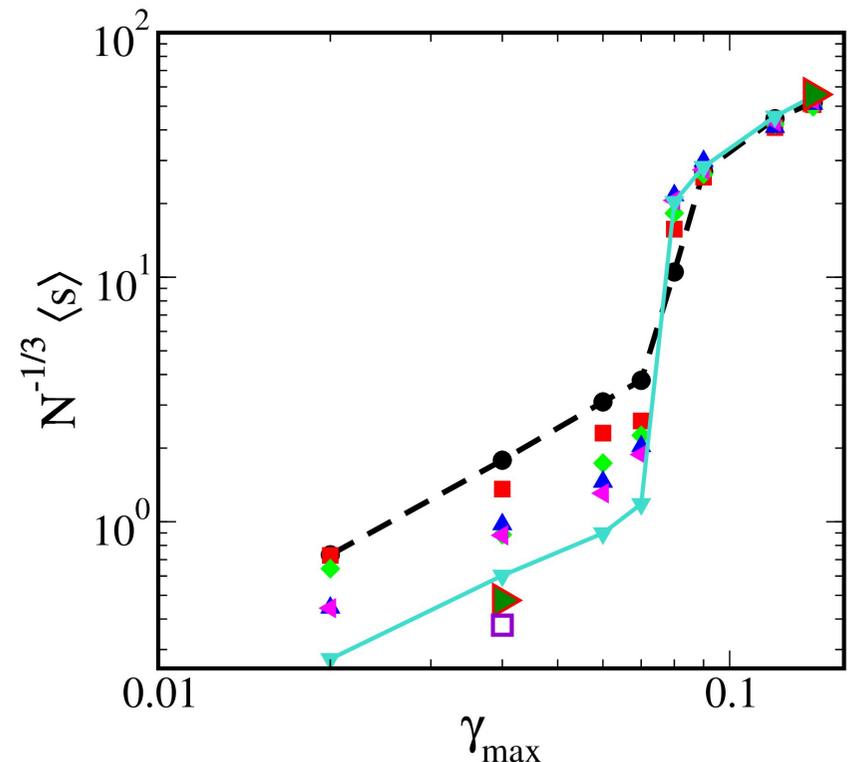
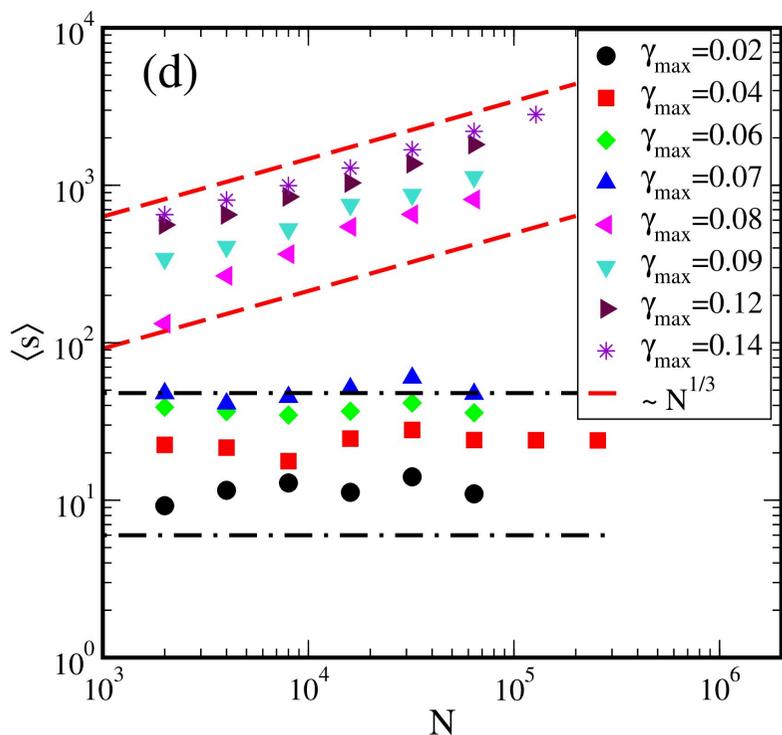
# Avalanche size vs. Strain Amplitude System Size Effects

**Above the yield point** mean avalanche size shows clear size dependence  $\sim N^{1/3}$

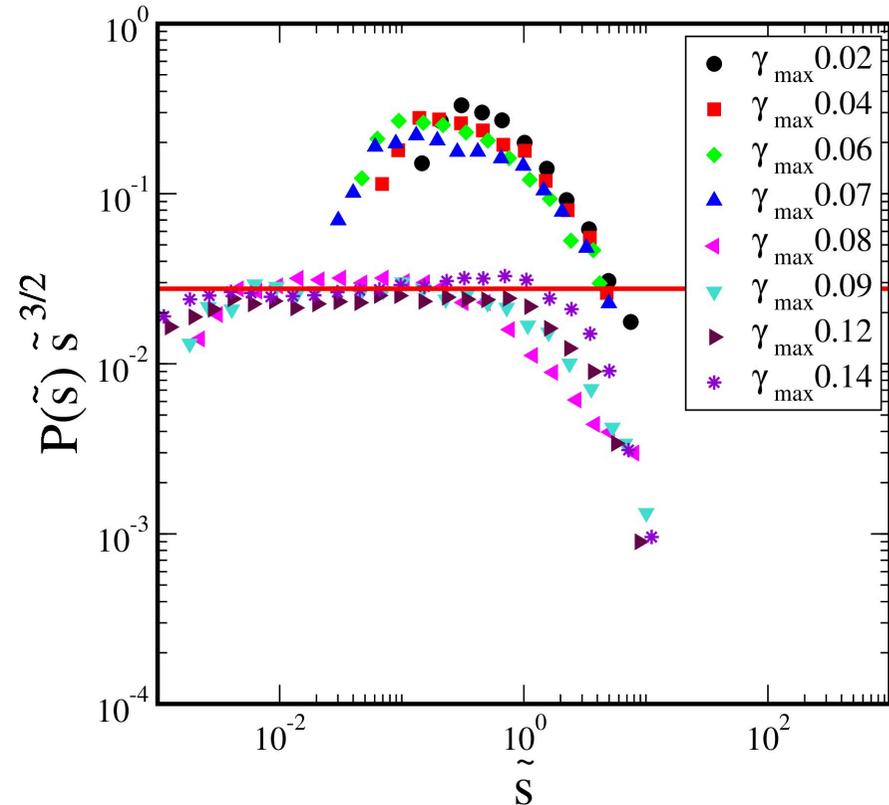
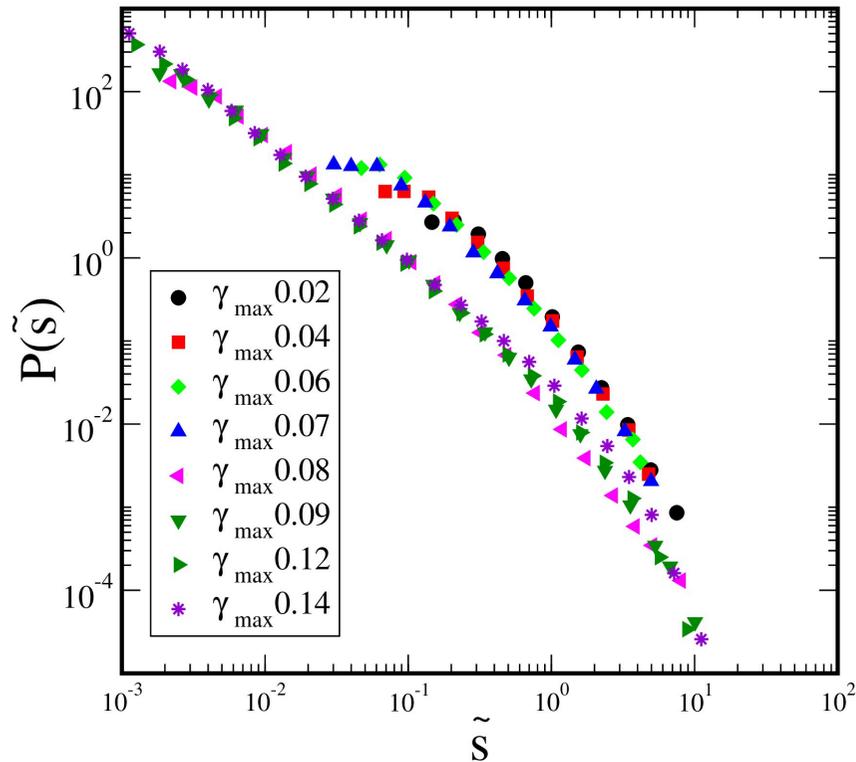
Consistent with uniform shear results for the case of plastic flow states.  
[Lerner and Procaccia 2009]

**Below the yield point**, avalanche sizes show **weak system dependence** that saturate.

## Mean Cluster Sizes Vs. N



# Scaled distributions

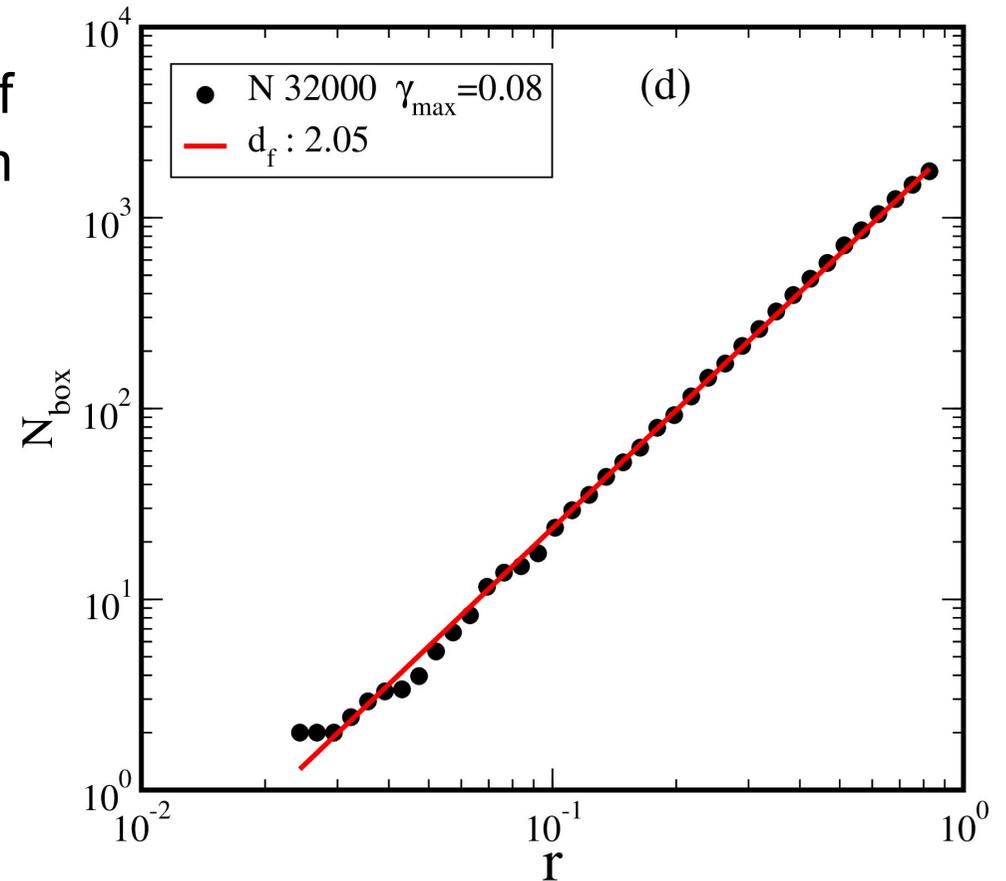
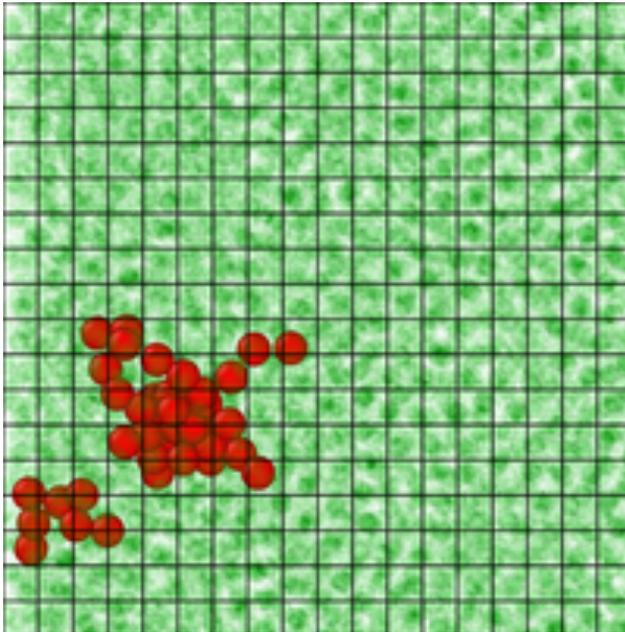


- Scaled cluster size ( $= s/\langle s \rangle$ ) distributions exhibit data collapse separately for  $\gamma_{\max} < \gamma_y$  and  $\gamma_{\max} > \gamma_y$ .
- Distributions for  $\gamma_{\max} < \gamma_y$  do not display a power law regime, whereas  $\gamma_{\max} > \gamma_y$  do.

**Qualitatively different avalanche distributions!**

# Fractal dimension

Fractal dimension of the clusters of the “active particles” estimated from box counting.



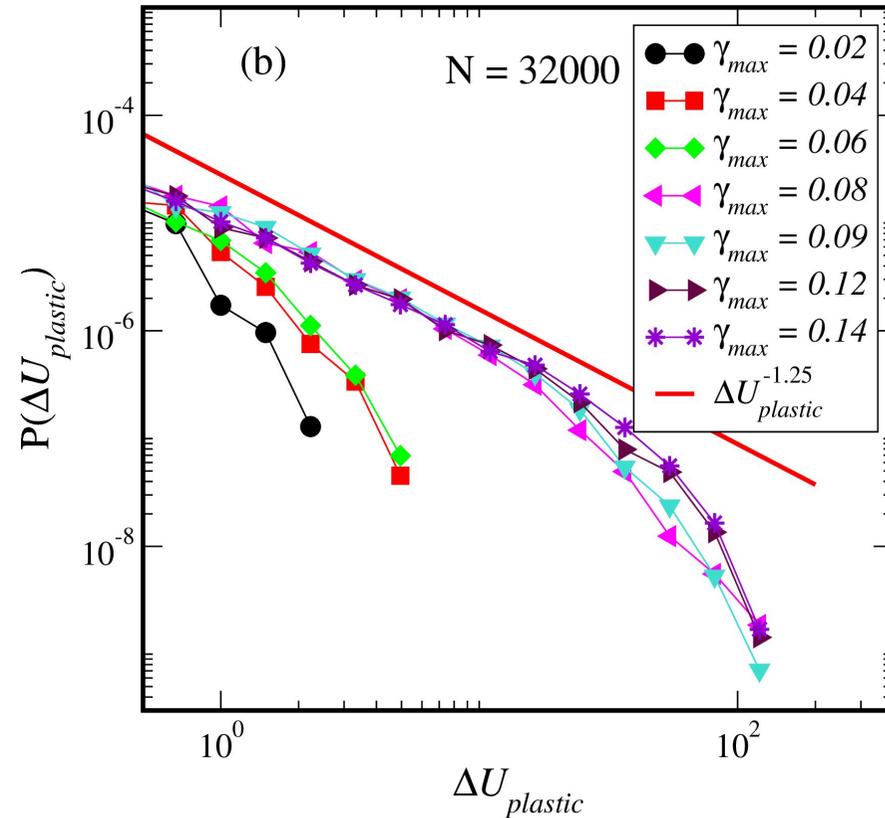
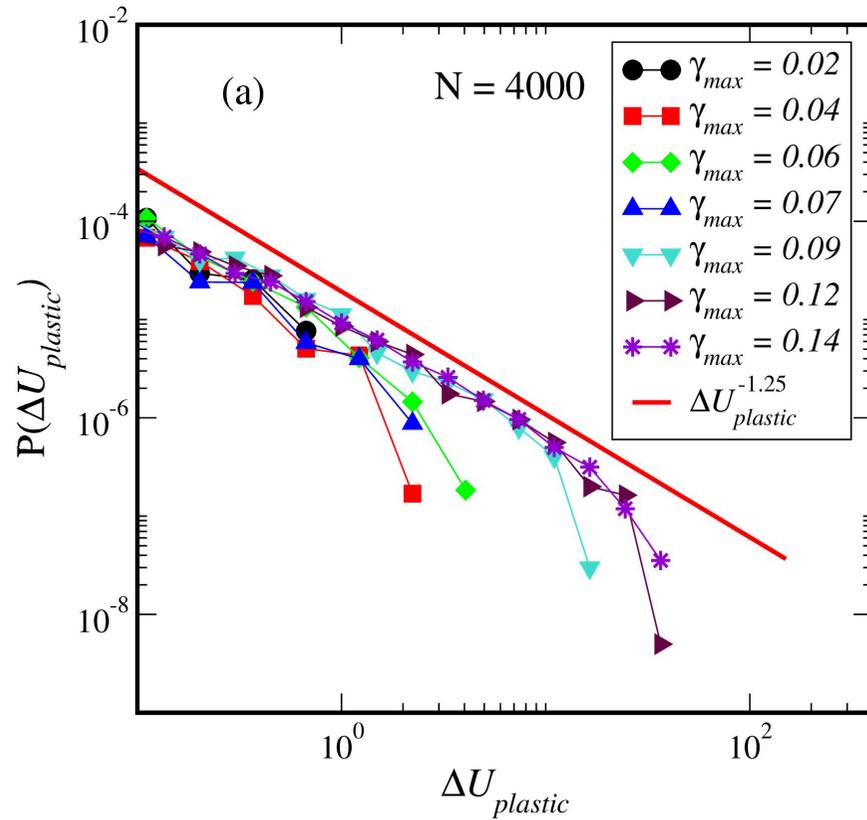
A log-log plot of the number of occupied boxes ( $N_{\text{box}}$ ) is shown vs. the magnification  $r$ .

The slope results in an estimated fractal dimension  $d_f \sim 2$ .

Not consistent with size scaling exponent of 1.

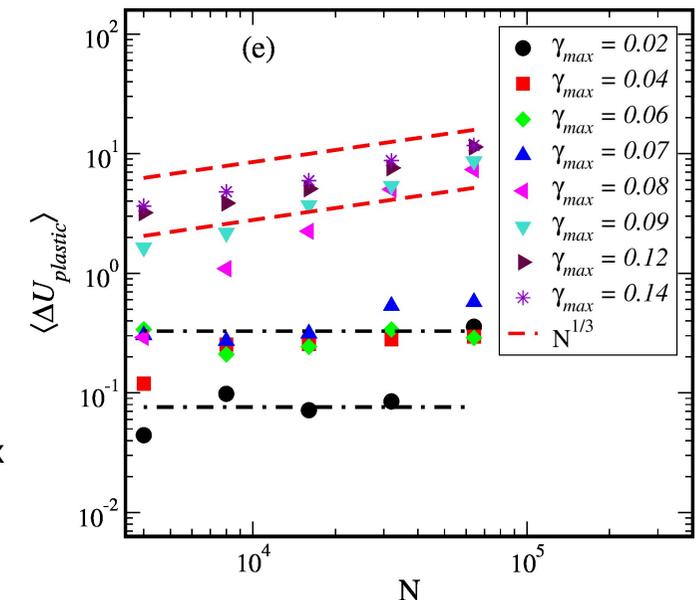
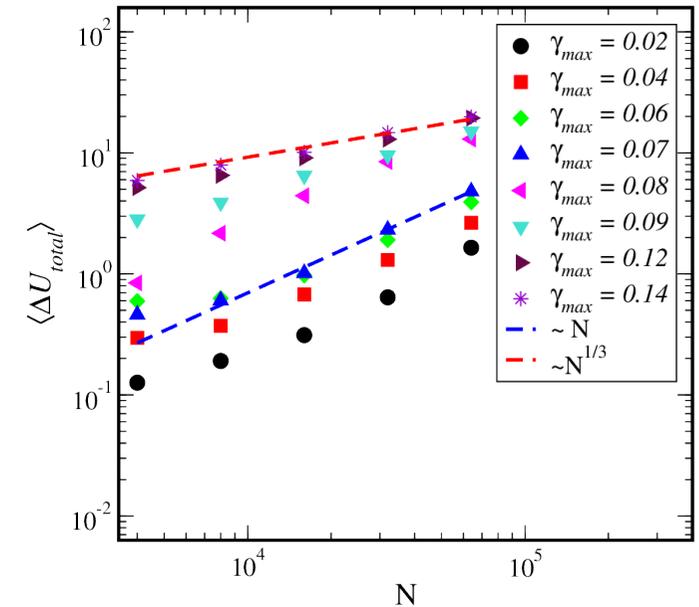
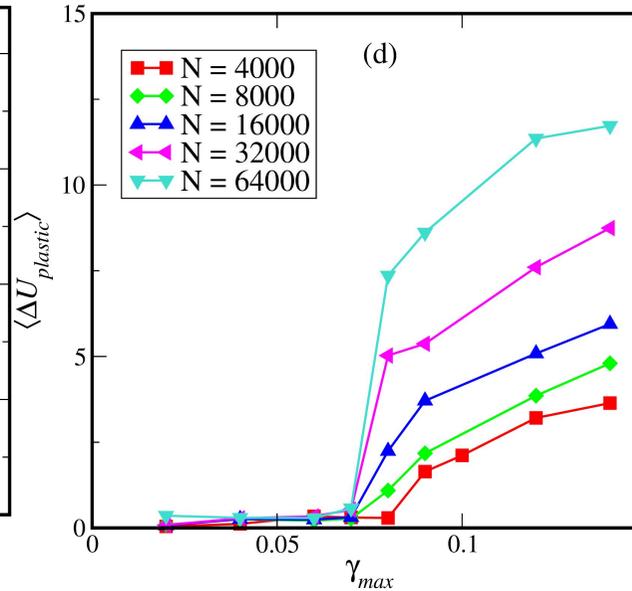
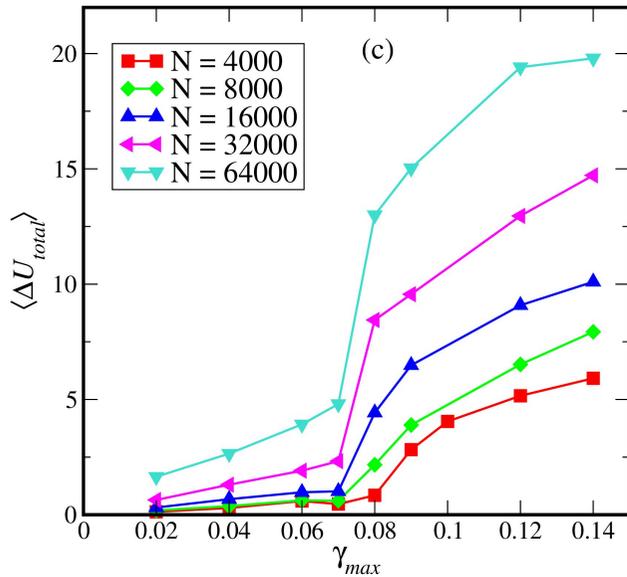
Anisotropies? [Maloney & Robbins PRL 2009]

# Avalanche: Energy drops



Energy drops distribution shows power law distribution with cutoff and the cutoff increases with  $\gamma_{max}$ .

# Avalanche: Energy drops

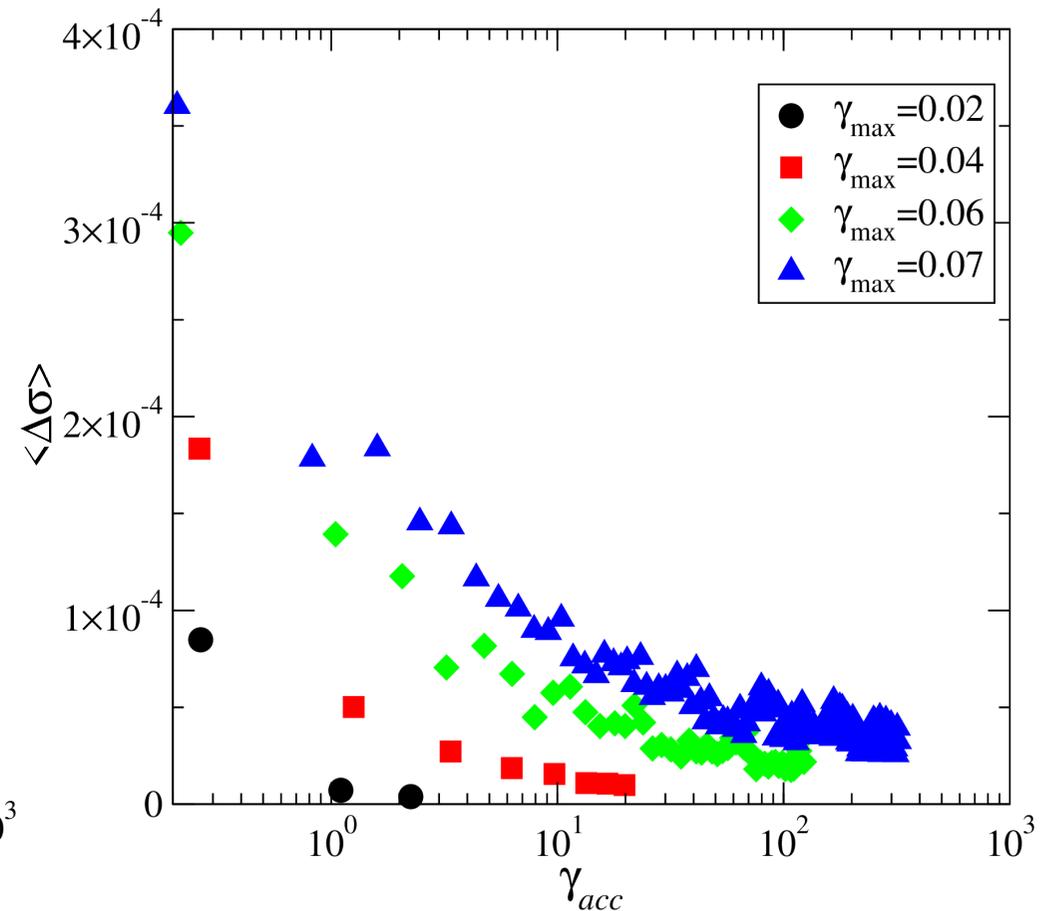
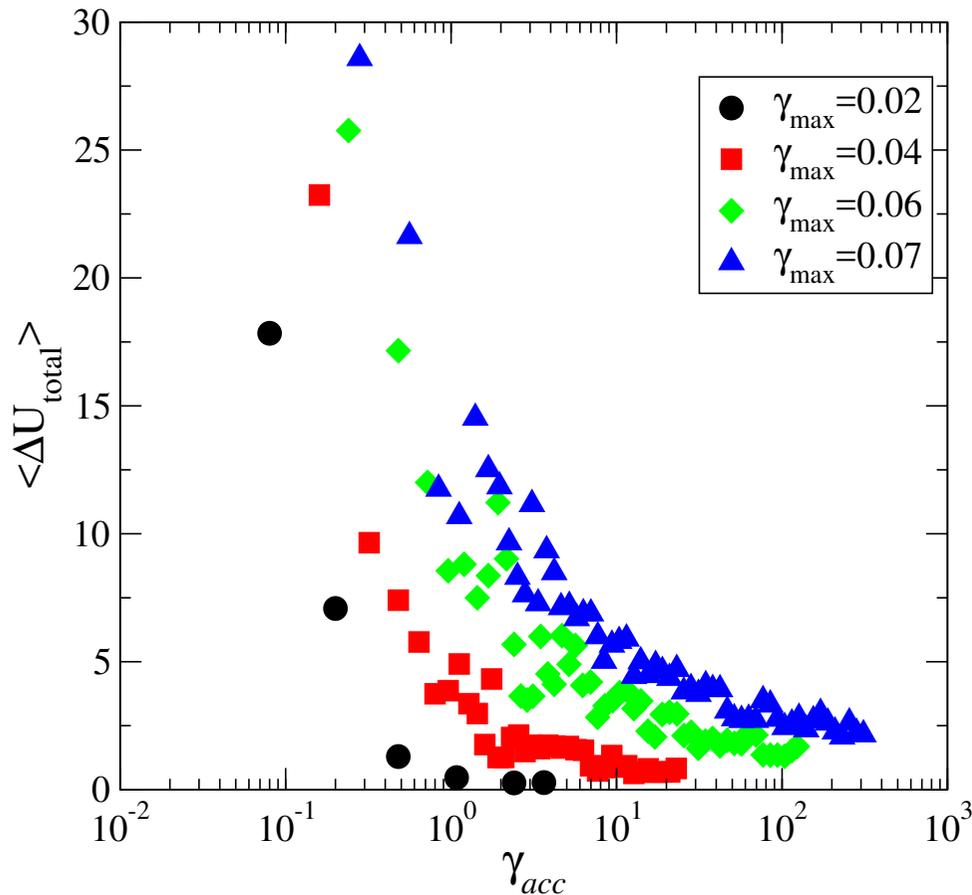


- Mean energy drops vs.  $\gamma_{max}$ , indicating a sharp change at  $\gamma_y$ .
- Mean energy drops considering only plastic regions show no system size dependence below  $\gamma_y$ .
- Mean energy drop (plastic component) vs. system size  $N$  shows no significant size dependence for  $\gamma_{max} < \gamma_y$  but a clear  $N^{1/3}$  dependence above.

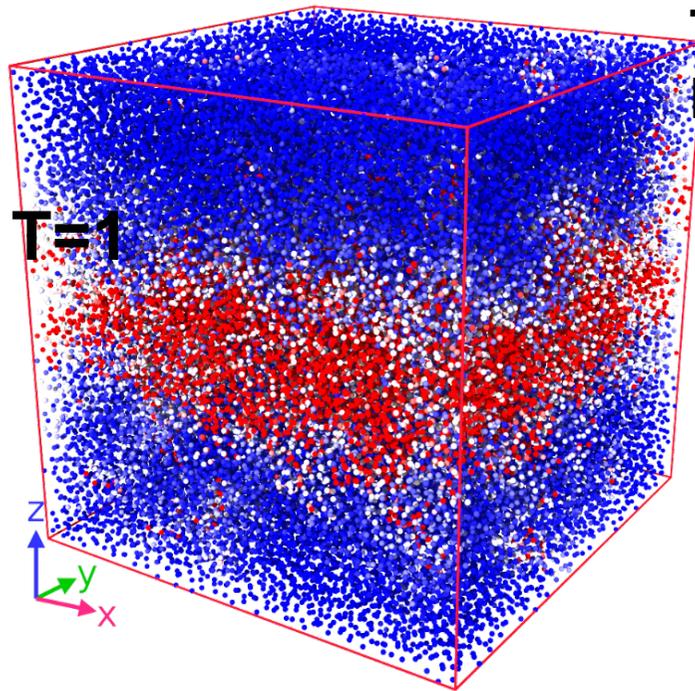
# Avalanches during the transient

Although the avalanche sizes are small below yield in the steady state.  
Large organizations occur during the transient.

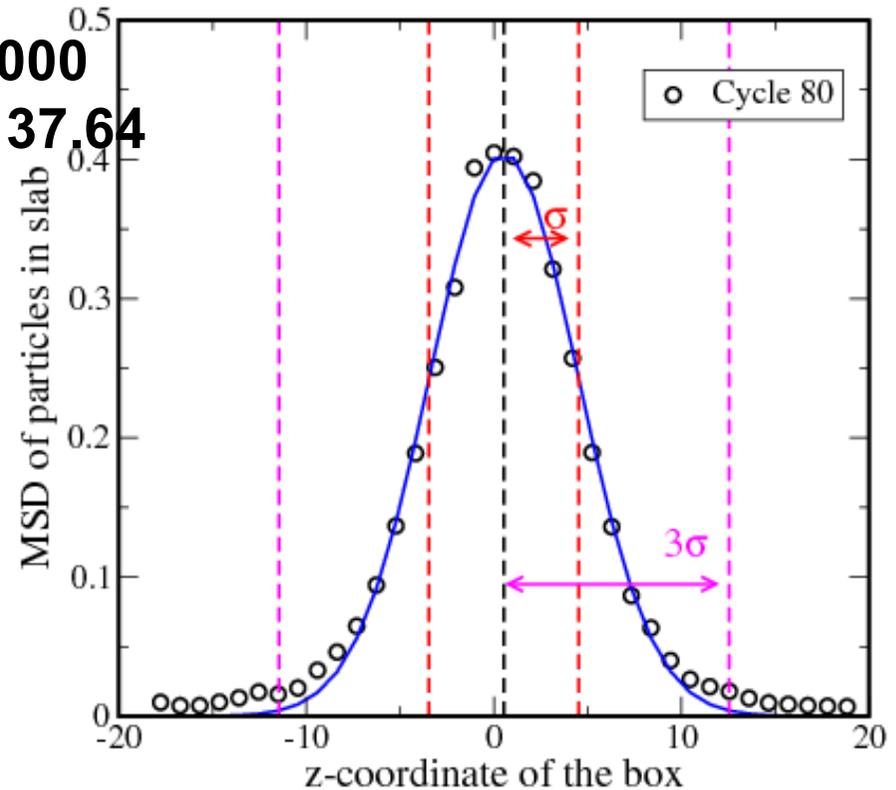
Larger and longer lived for larger strain amplitudes.



# Shear Banding



$T=1$ ,  $N=64000$   
Box size = 37.64



- A configuration from the steady states for  $\gamma_{\max}=0.09$ .
- Color code: MSD values

- The slabwise MSD of particles as function of the height of the band.

Yielding is accompanied by the formation of shear bands.

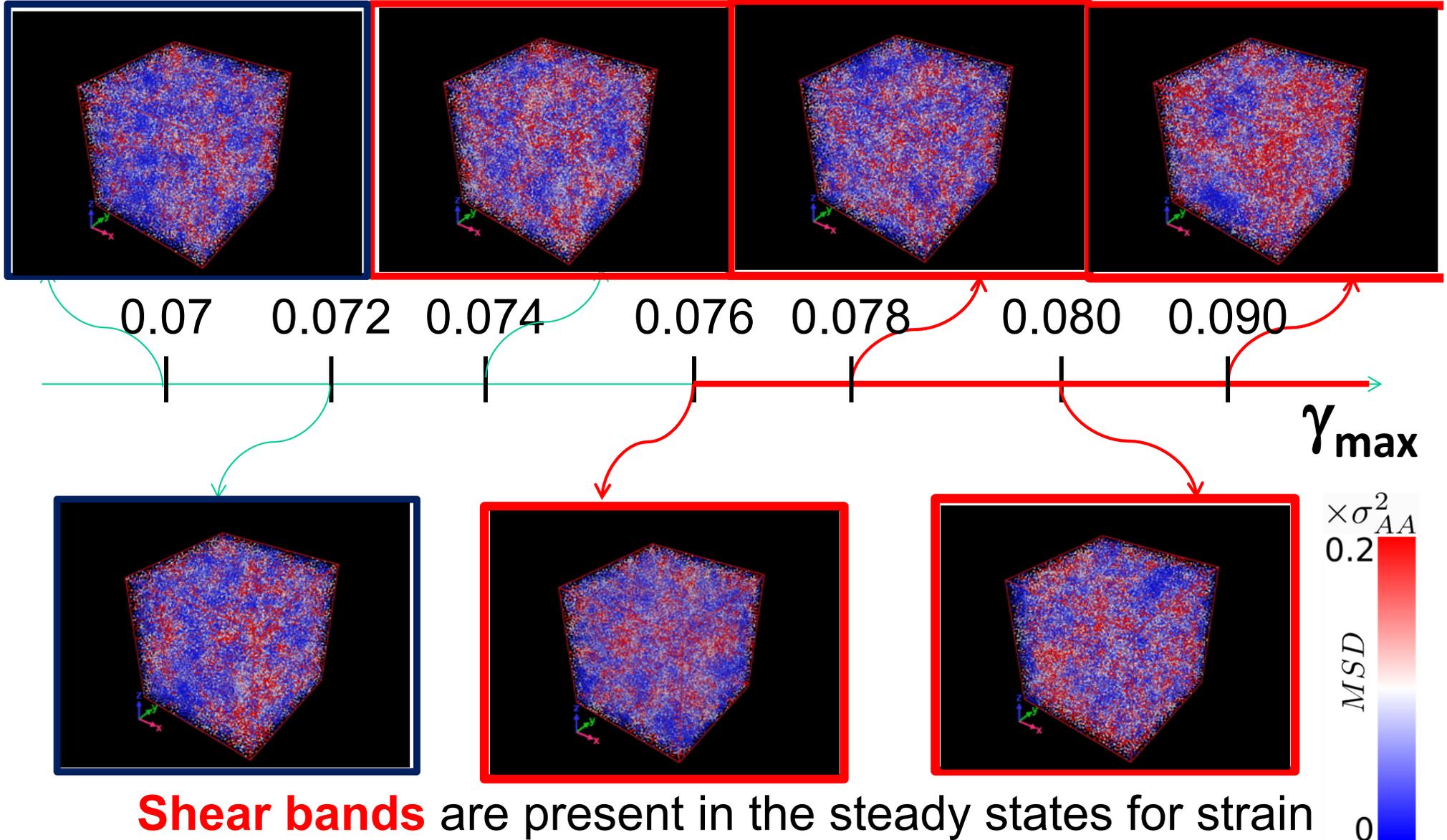
Displacements per cycle, energy etc indicate presence of shear bands  $\sim 20$  particle diameter!

Probed by different initializations (from liquid/poorly annealed glass, strain amplitude below (.07), above (.08) yield value).

# Across the yielding transition

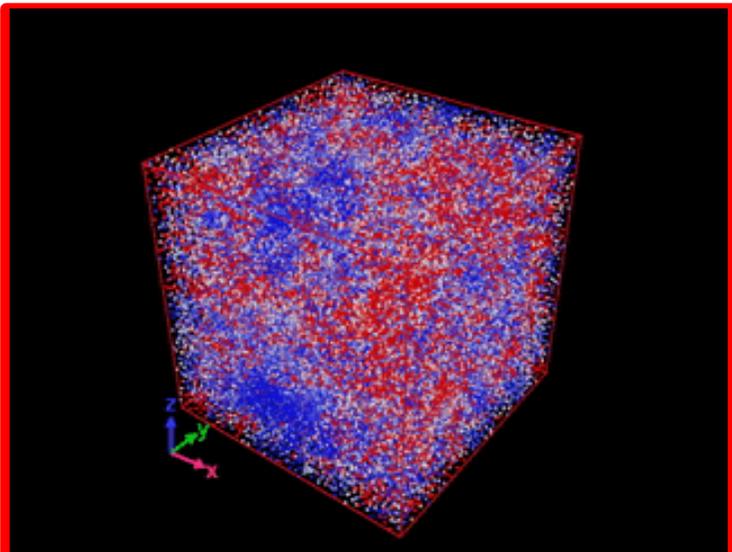
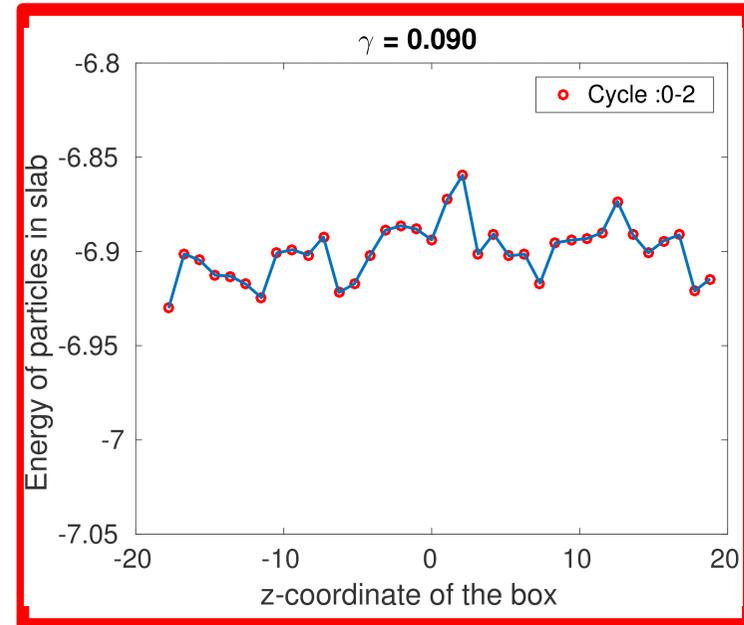
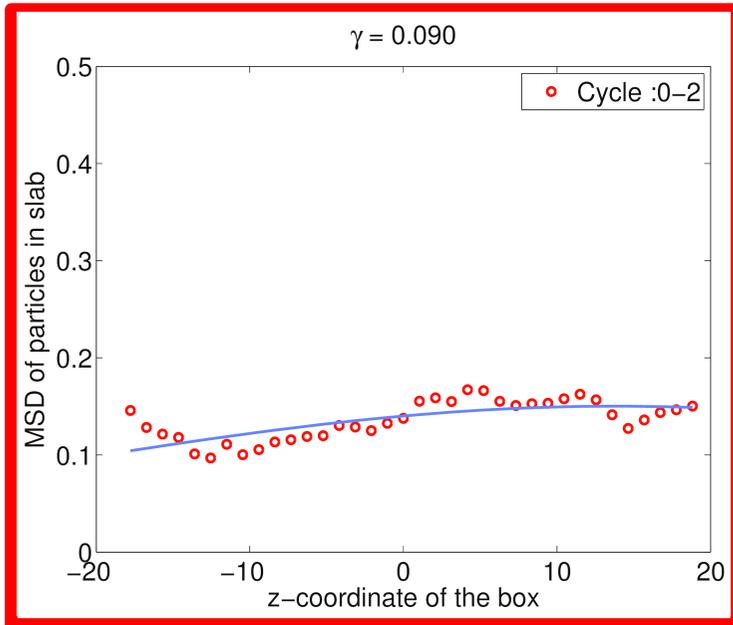
N=64000

From poorly annealed glass



**Shear bands** are present in the steady states for strain amplitudes above yielding value  $\sim 0.07$

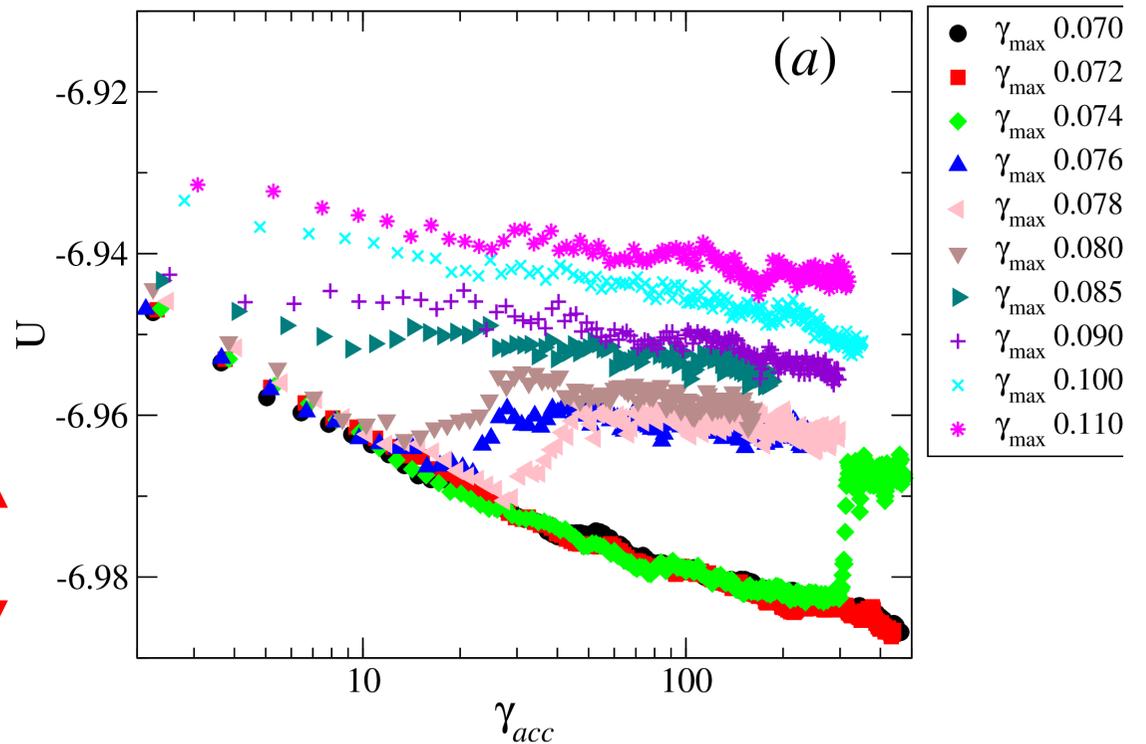
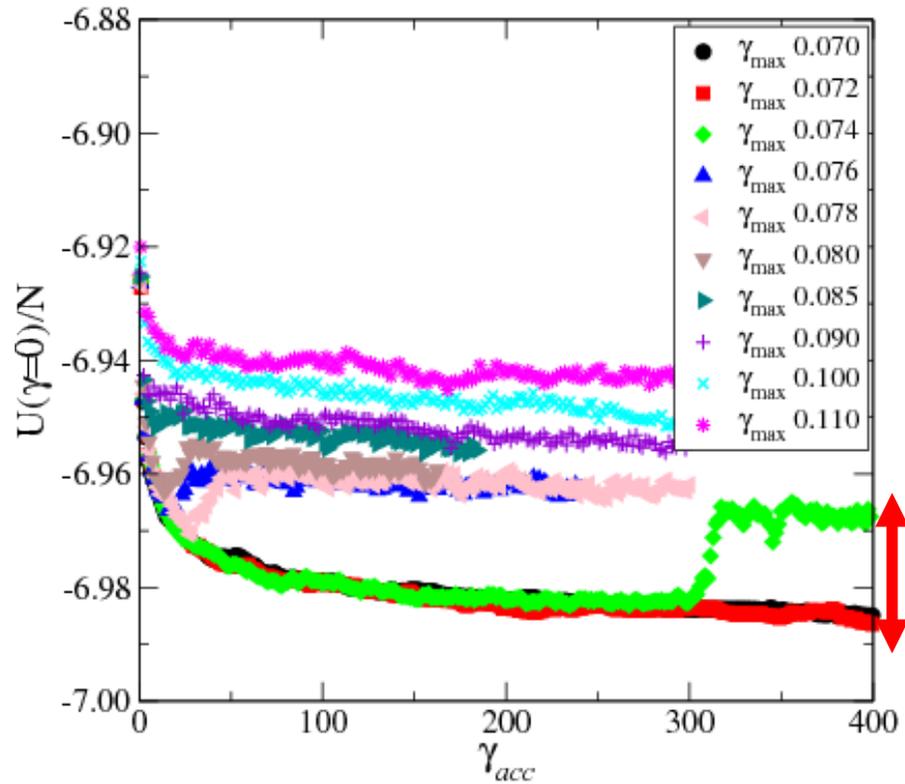
# Evolution towards steady states



- Shear bands are present in the steady states.
- The energy and MSD profile changes in the presence of the shear band.

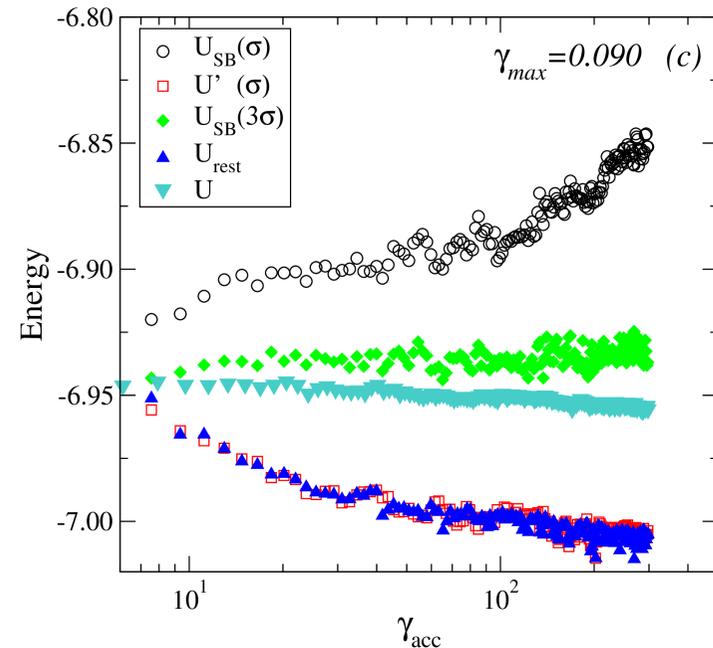
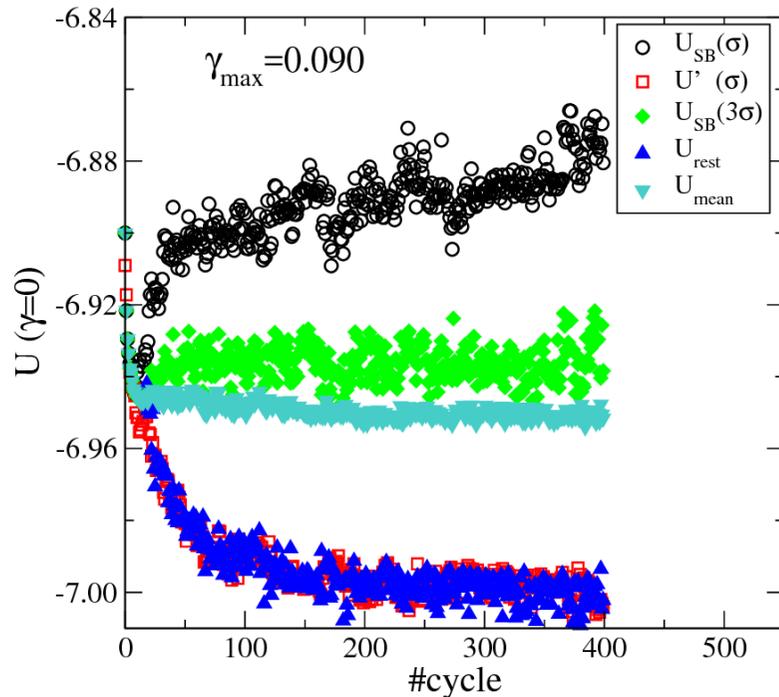
$$\gamma_{\max} = 0.09$$

# Evolution towards steady states



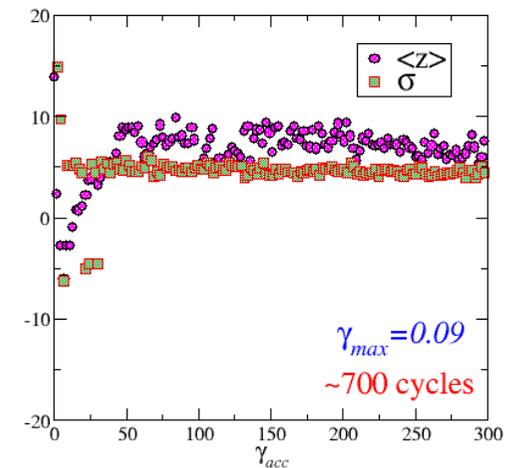
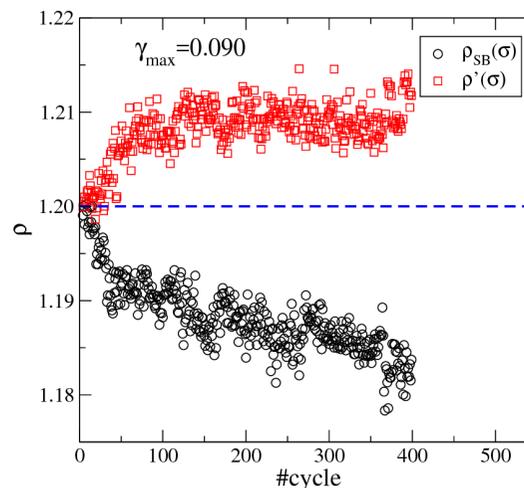
- Energies show monotonic change with cycles below yielding.
- Logarithmic relaxation!
- Abrupt change when yielding (and shear banding) happens.

# Evolution towards steady states

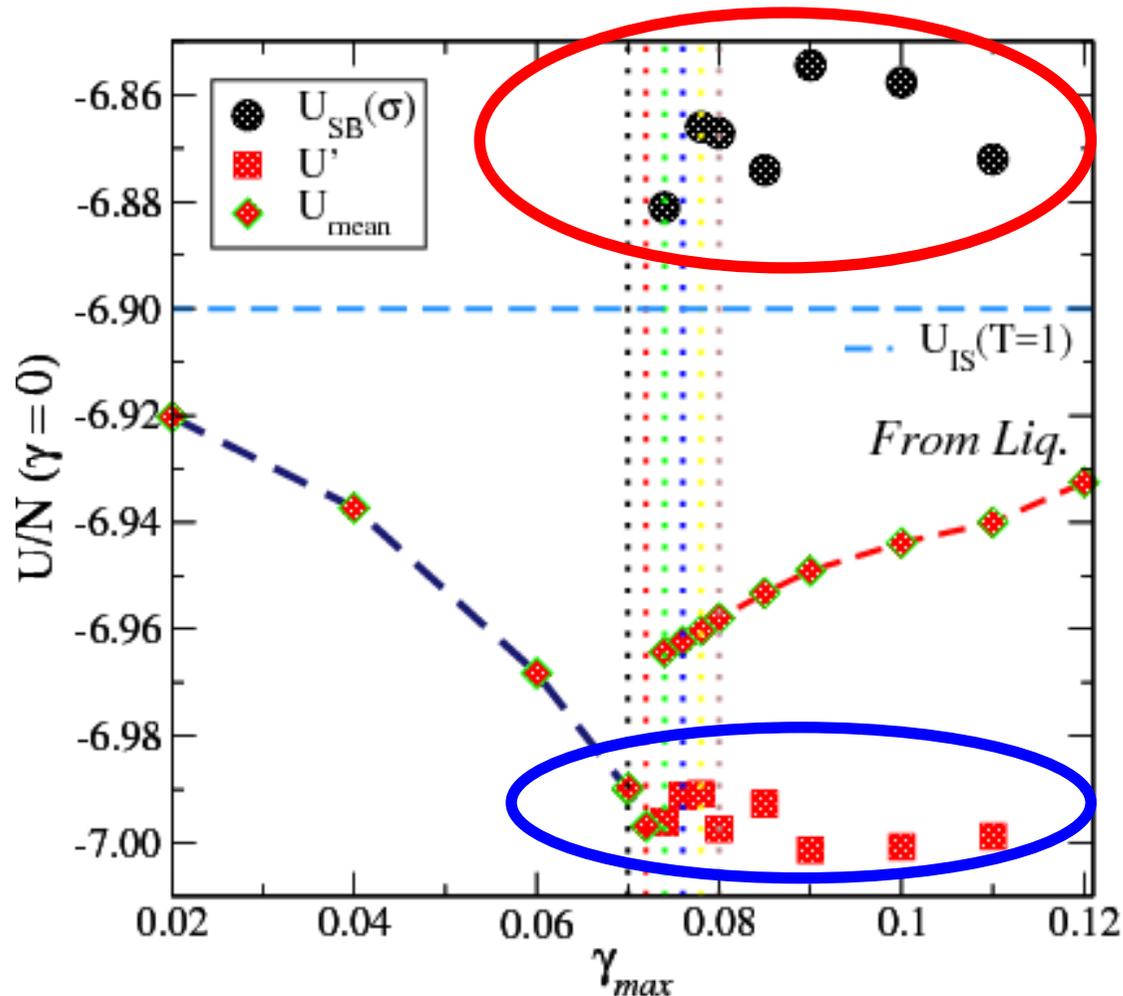


Shear band  $\gamma_{\max}=0.09$

- The energy and mobility of the particles in the shear band are higher than particles outside the shear band.
- The shear band is “fluid like”, and density of the band is less than the bulk of the system.
- Shear bands are mobile but the width is stationary.



# Annealing above Yielding



Energy in shear band ~ highest temperature glasses.

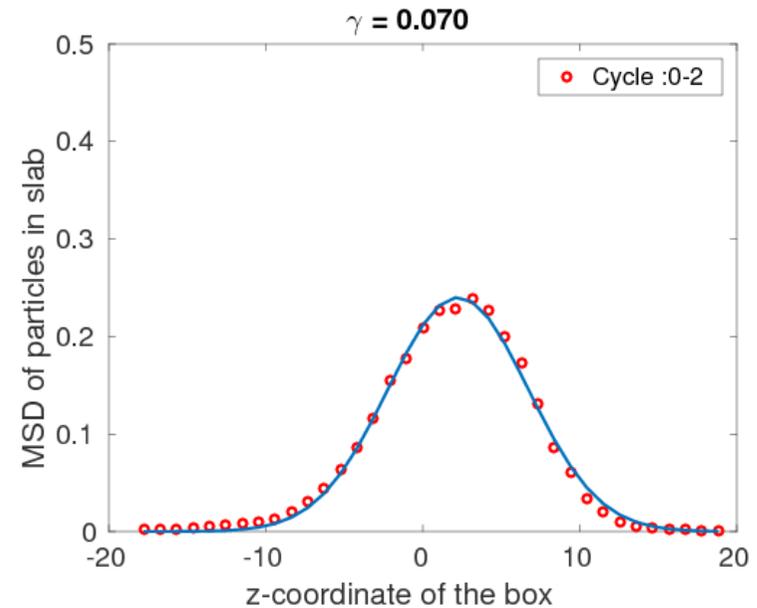
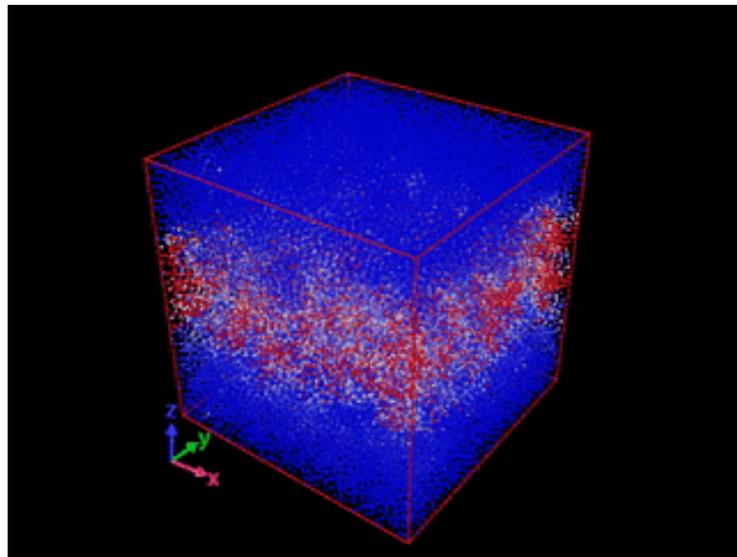
Finite jump in the mean energy.

Continued annealing outside the shear band

- Particle in the shear band access energies at the “top of the landscape”
- The rest of the system continues to be annealed beyond yield point.
- Location of yielding point exhibits some initial condition dependence.

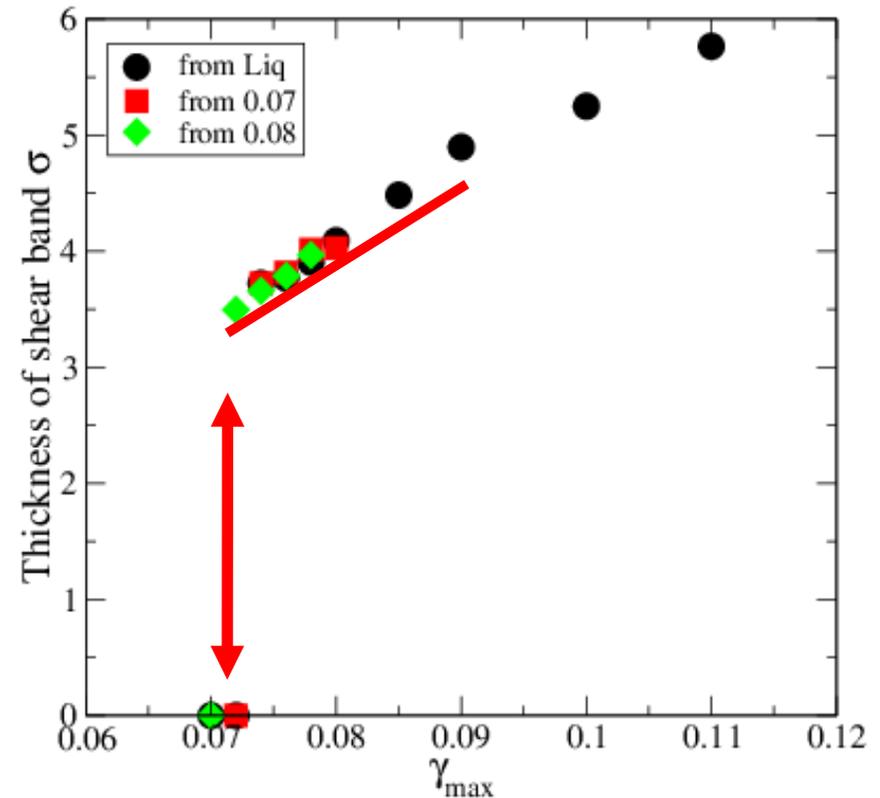
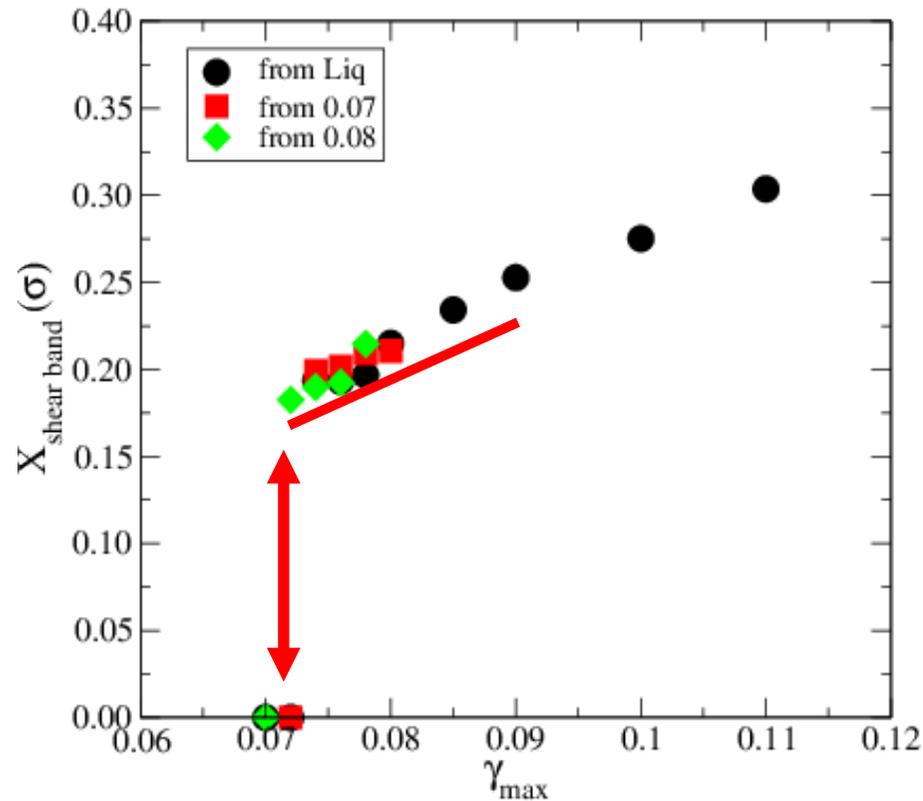
# Vanishing shear band

From  $\gamma_{\max} = 0.08$ , @  $\gamma_{\max} = 0.07$



The shear band vanishes after a large number of strain cycles at smaller amplitude.

# Characteristics of the shear band



The **fraction** of the particles, **width** of the shear band changes in **discontinuous manner** at the critical amplitude.

How does diffusivity change?

$\gamma_{\text{max}}^c$

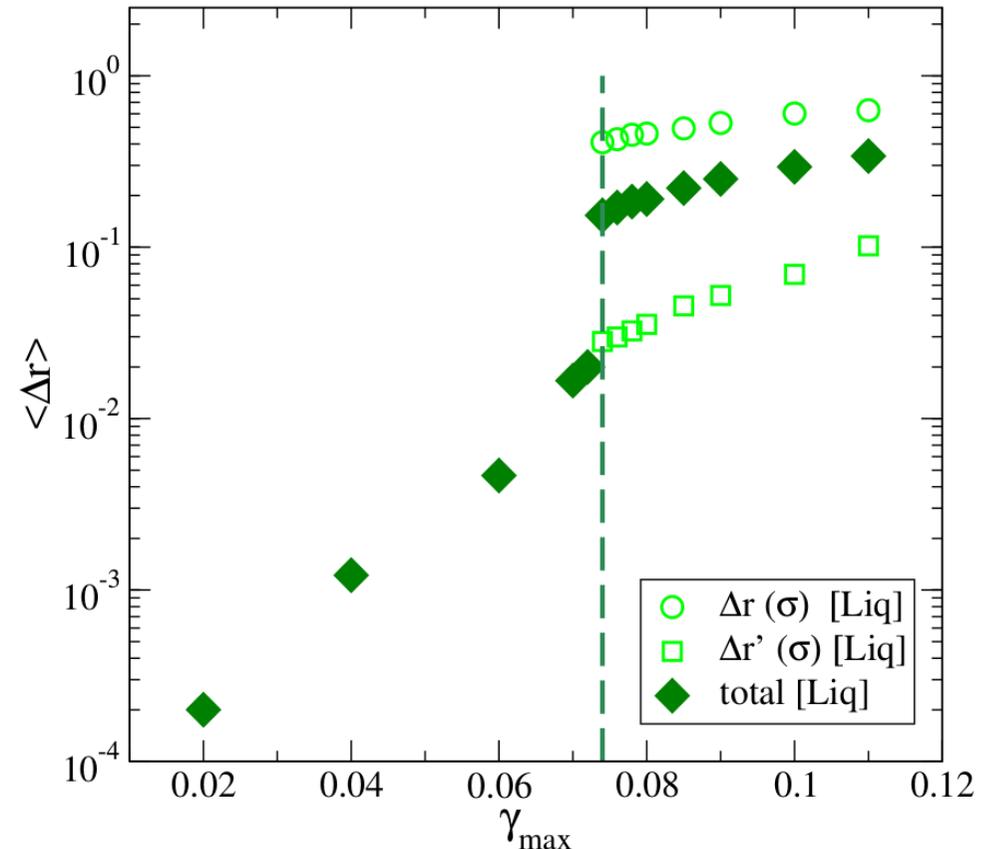
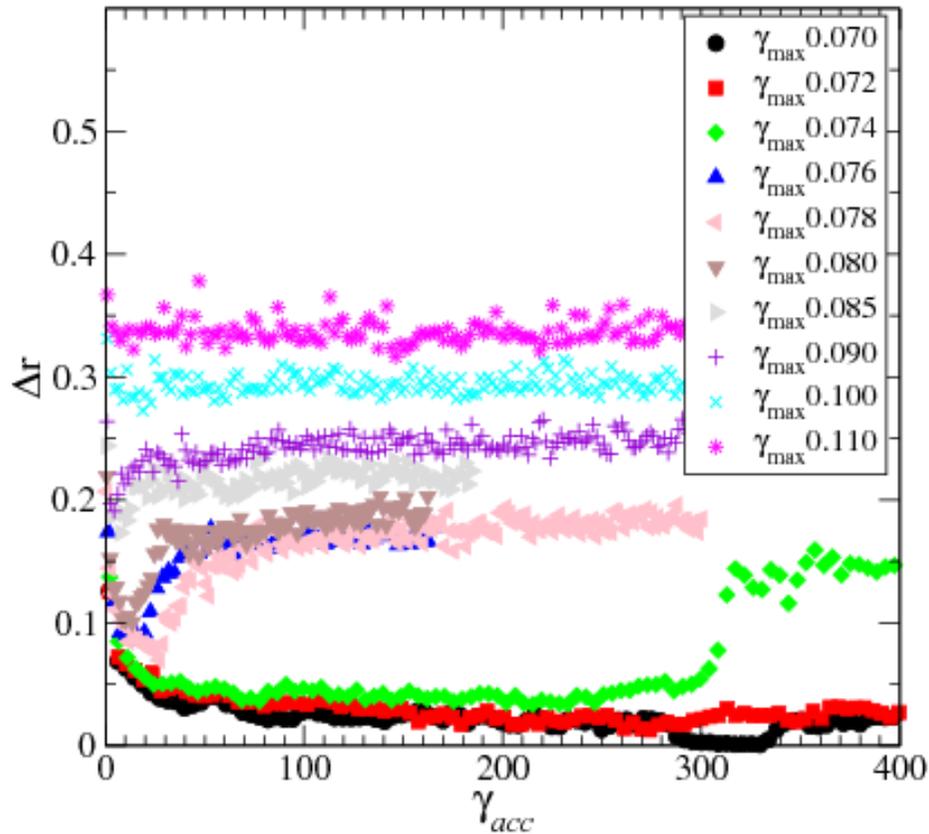
0.074

0.072

# Microscopic dynamics associated with shear banding

The average particle displacement after one deformation cycle:  $j$  is the index of cycle.

$$\Delta r = \frac{l}{N} \sum |\vec{r}(j+1) - \vec{r}(j)|$$

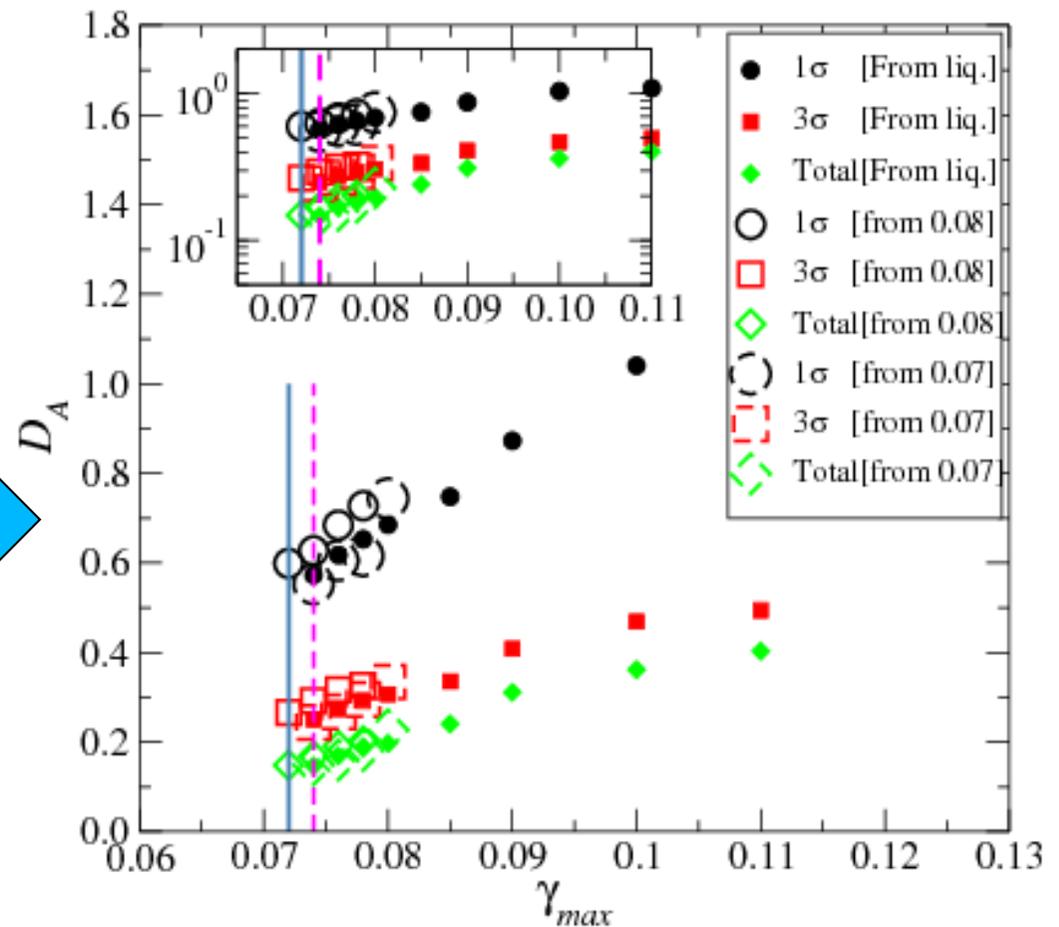
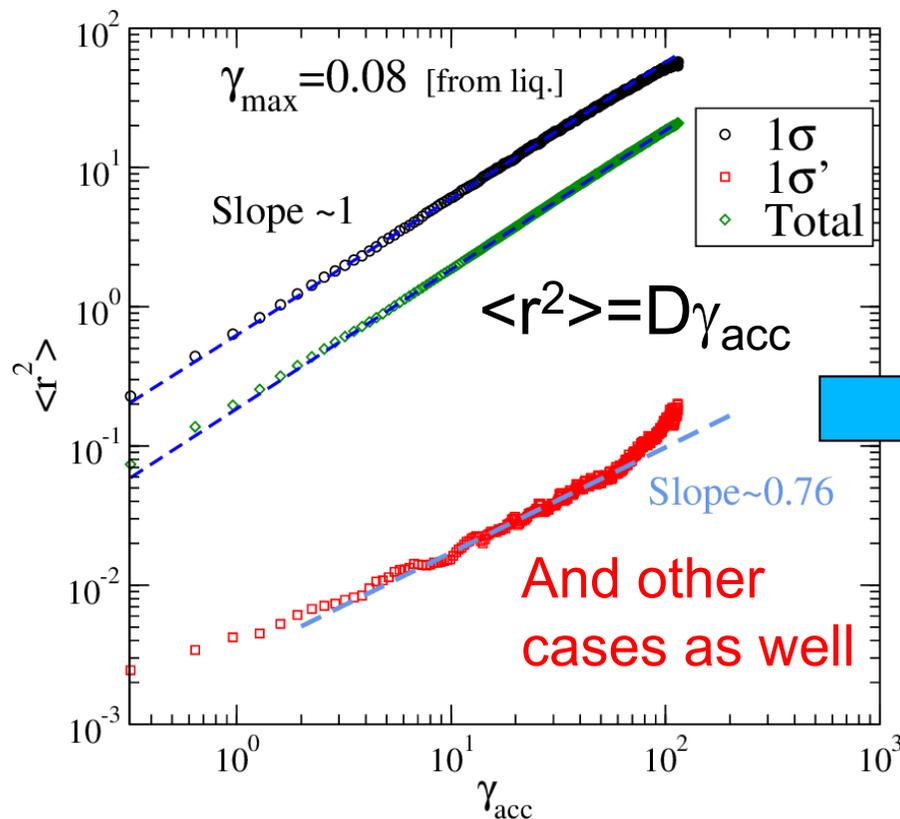


Across the yielding transition, the averaged particle displacement changes in a discontinuous manner.

Outside the shear band and below yielding, movement finite but very small.

# Characteristics of the shear band - Diffusivity

The diffusion of **total system**, **most mobile particles** ( $1\sigma$ ), the shear band ( $3\sigma$ ) and least mobile particles ( $3\sigma'$ ) estimated.



The **diffusion coefficient** changes in **discontinuous manner** at the yielding strain amplitude.  
Sub-diffusive behaviour outside the shear band.

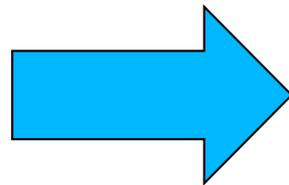
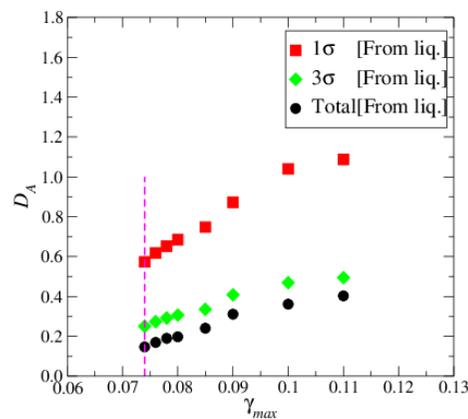
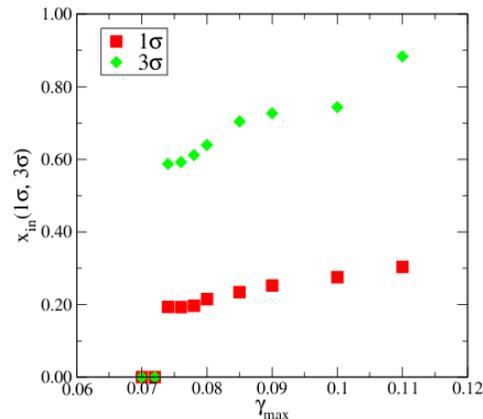
$\gamma_{max}^c$

0.074

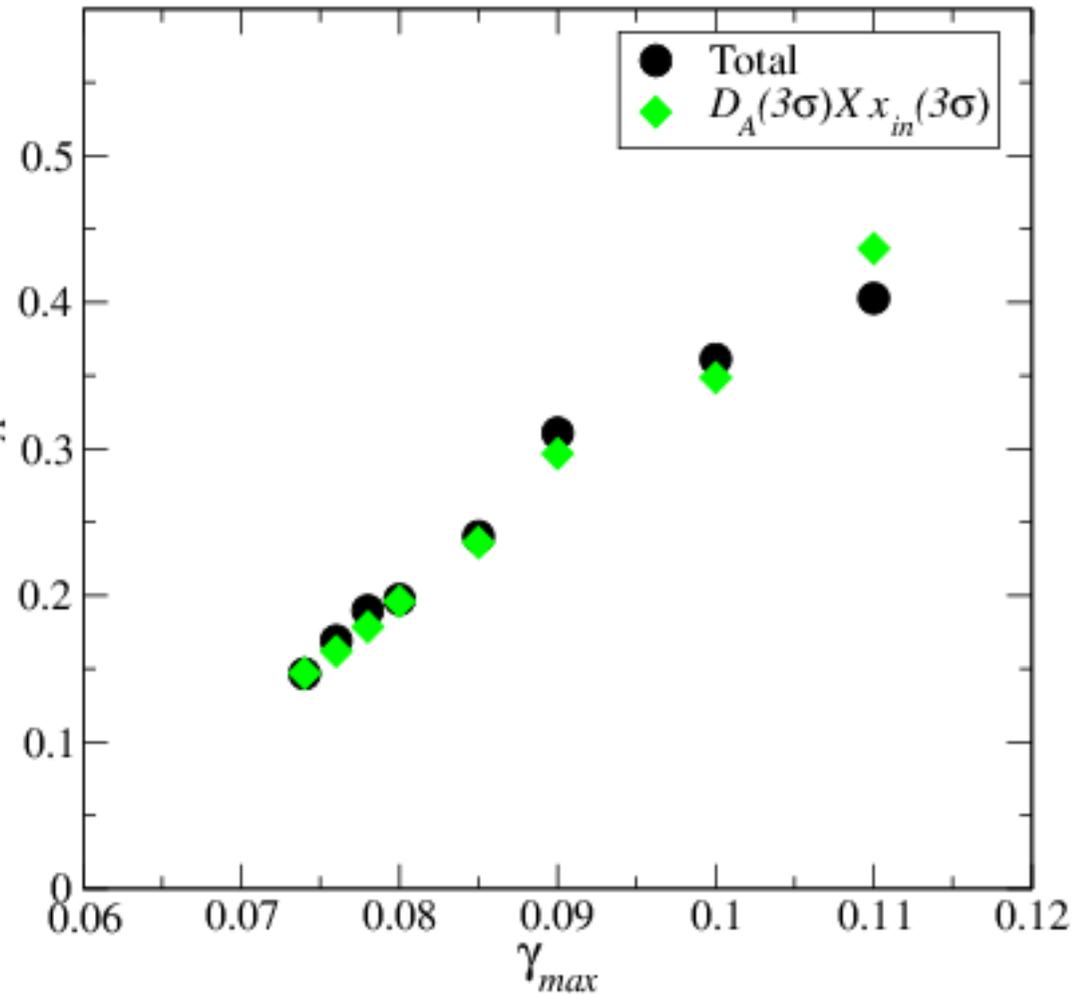
0.072

# Characteristics of the shear band - Diffusivity

Above the yielding transition, the finite diffusion is an outcome of the shear band.



$D_A$



The **diffusion coefficient** of the mobile particles changes in **discontinuous manner** at the critical amplitude.

# Finite Temperature and Shear Rates

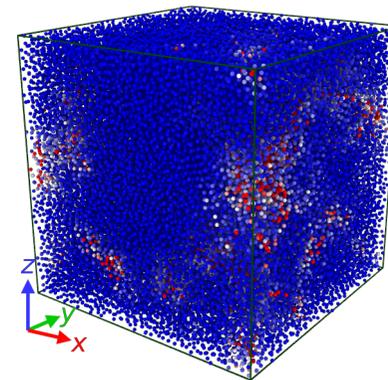
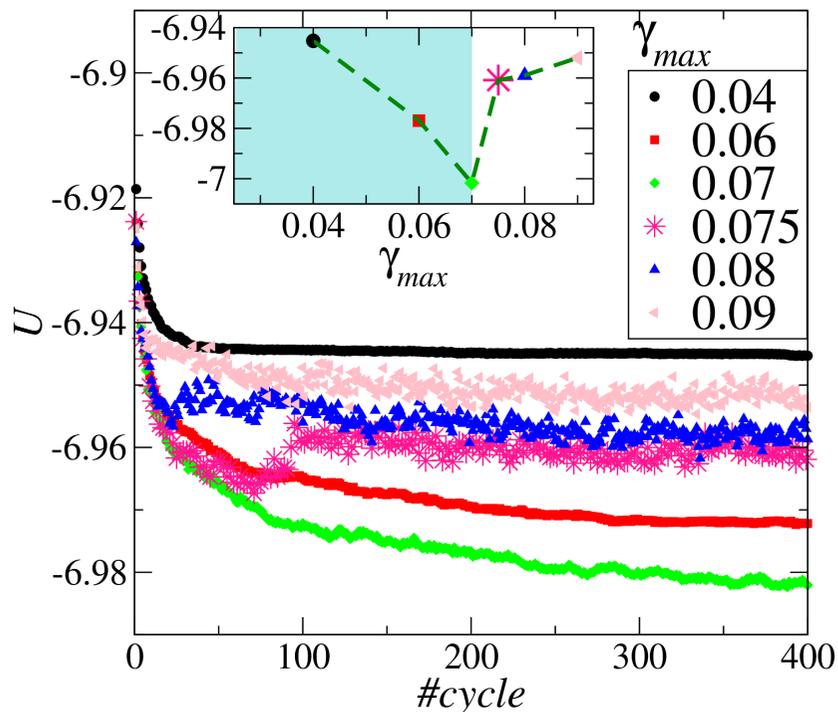
What happens if we go beyond the AQS limit?

We perform shear deformation at finite strain rates and temperatures.

The system is sheared (*via* SLLOD) at range of temperatures 0.01, 0.1, at strain rate  $10^{-5}$  and varying amplitudes ( $\gamma_o$ ).

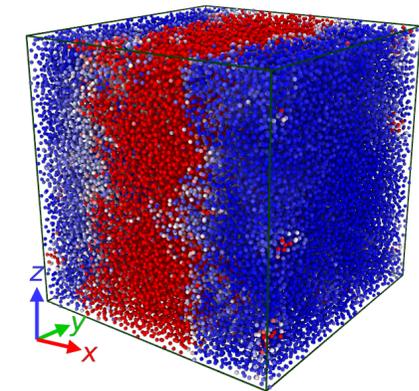
$$\gamma(t) = \gamma_o \sin(\omega T)$$

The picture does not change.



$\gamma = 0.07$

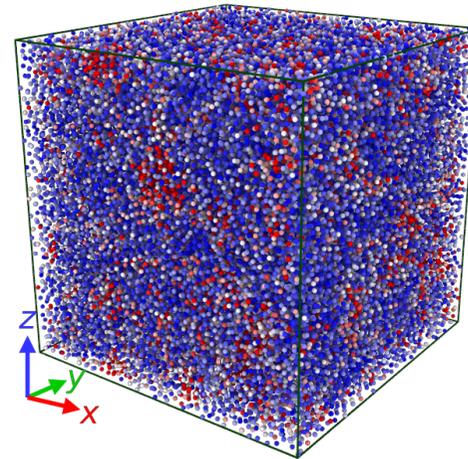
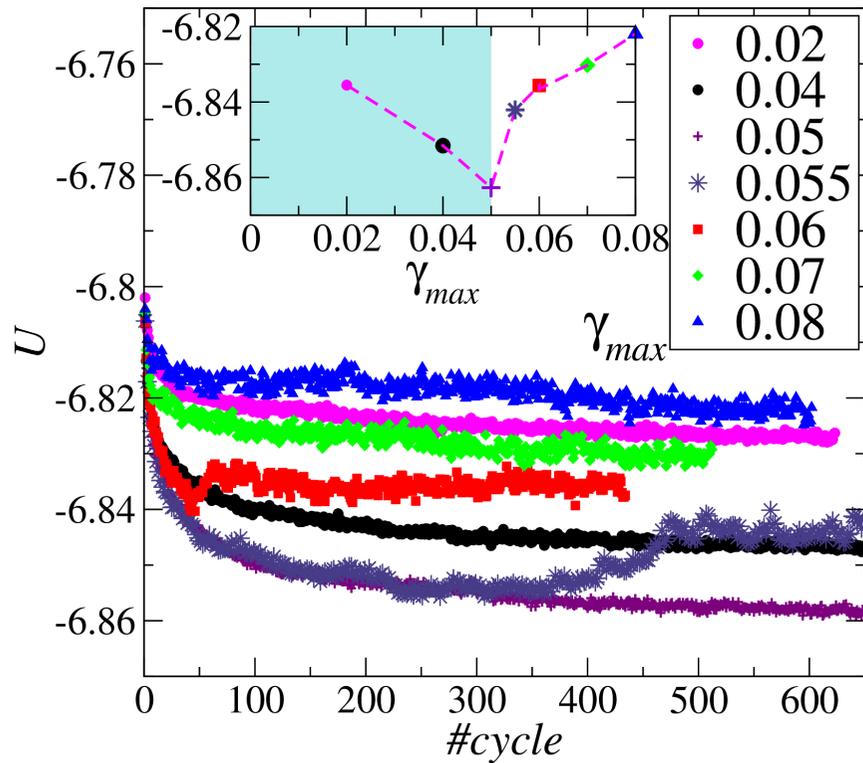
$\gamma = 0.075$



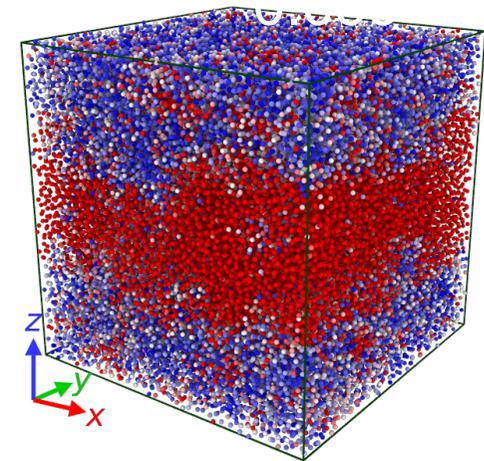
$T = 0.001$

# Finite Temperature and Shear Rates

$T = 0.1$



$\gamma = 0.055$



At a higher temperature, shear banding still present, but more fluctuations.

Lower yield strain amplitude.

# Entropic characterization of the yielding

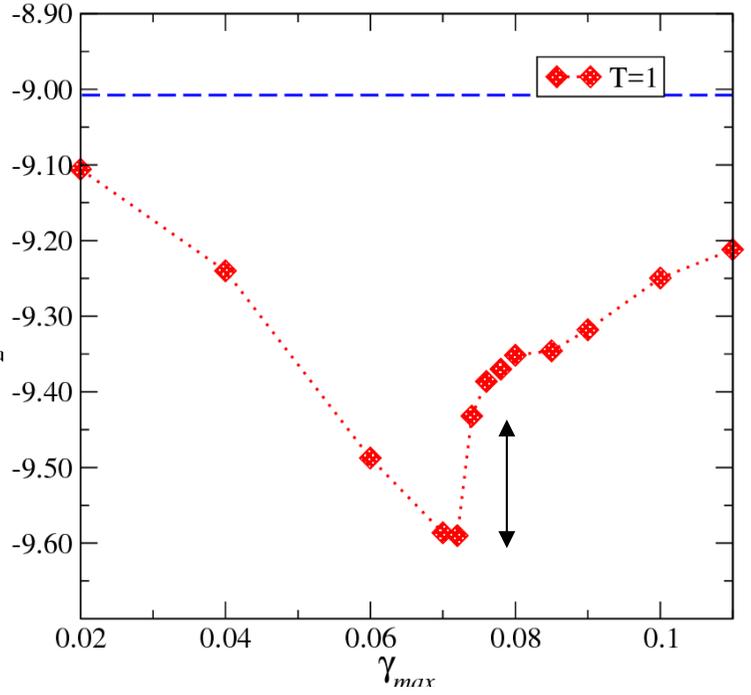
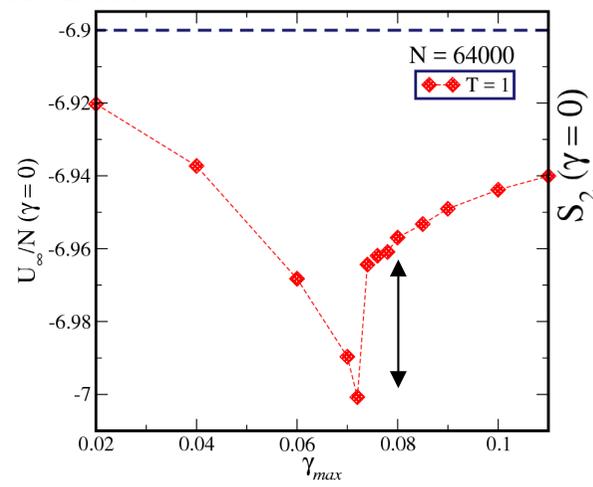
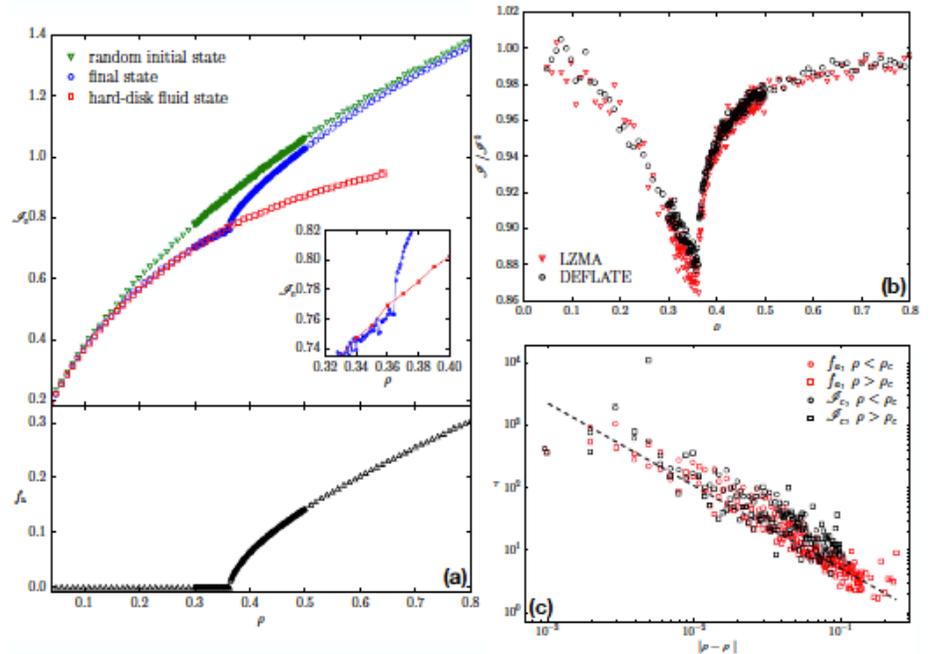
Recent interest in information measures  
In driven systems (CID Martiniani et al 2017)

**How about conventional structure based entropy? Pair excess entropy ( $S_2$ ) for a binary system:**

$$\frac{S_2}{k_B} = -2\pi\rho \sum_{\alpha,\beta} x_\alpha x_\beta \int_0^\infty \{g_{\alpha\beta}(r) \ln g_{\alpha\beta}(r) - [g_{\alpha\beta}(r) - 1]\} r^2 dr,$$

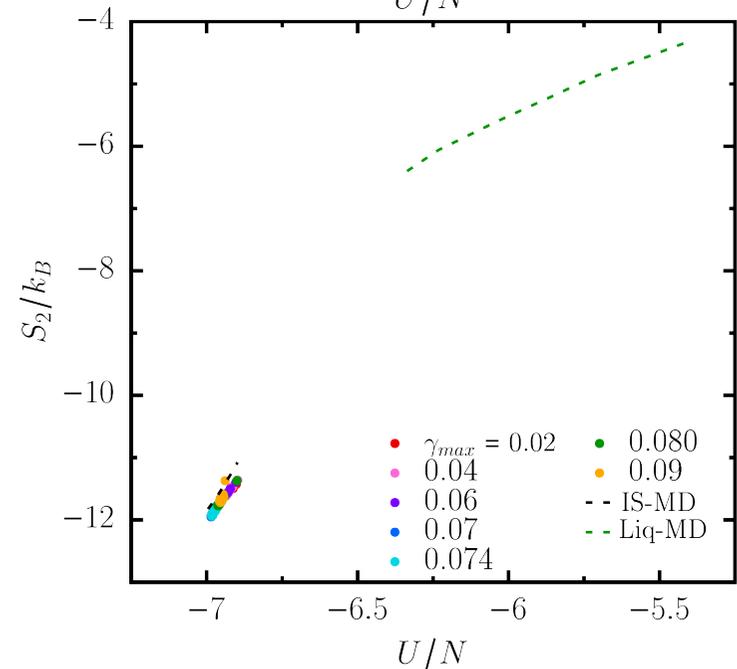
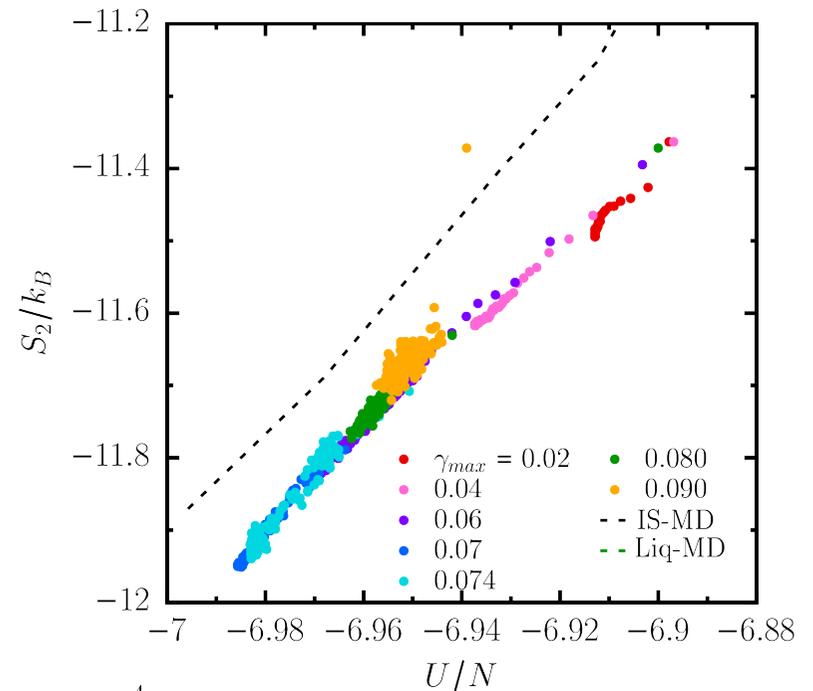
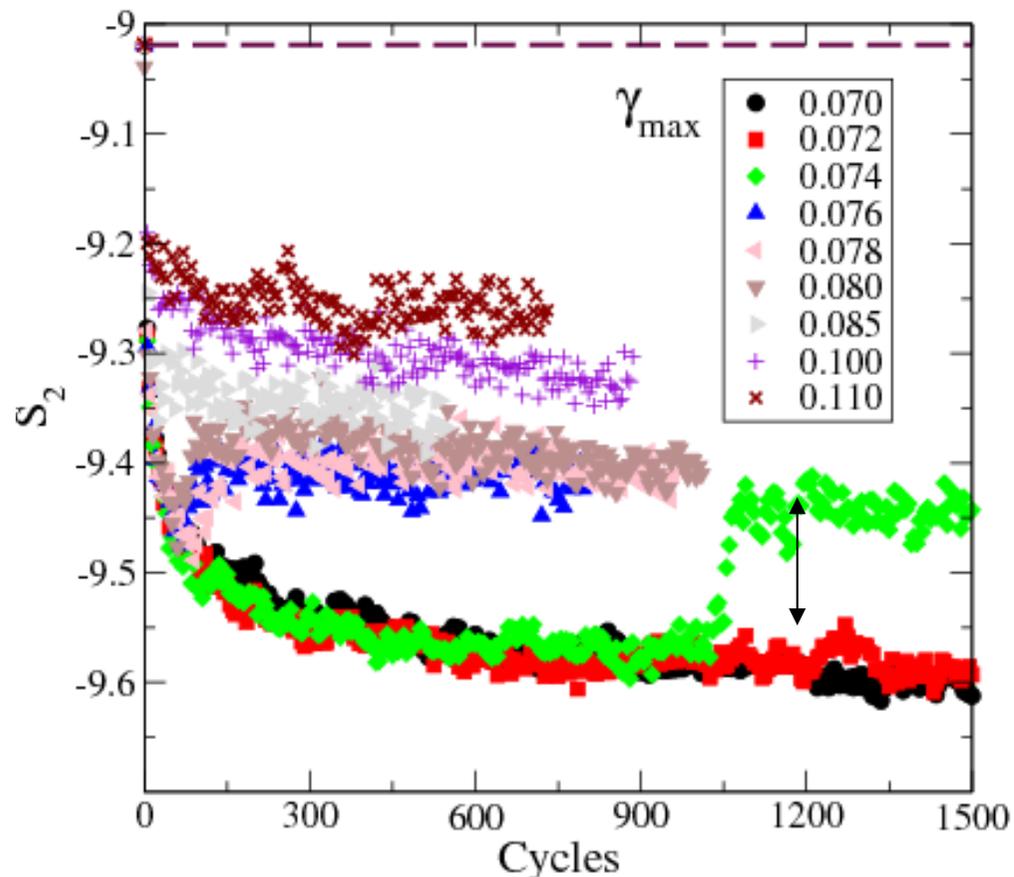
$g_{\alpha\beta}(r)$  is the pair correlation function for various components

**The  $S_2$  and total energy jumps across the yielding amplitude.**



# Entropic characterization of the yielding

- $S_2$  tracks the energy of the inherent structures.
- Close to relation found in  $T$  dependence for the liquid, but not the same.



## Annealing Effects are Important

Many past studies argue for a role of annealing (variously, relaxation, aging etc.. ) in the phenomenon of shear banding:

Pulled front/STZ: Alix-Williams and Falk 2018

SGR model: Fielding et al 2016

Elasto-plastic model: Martens et al, 2010

Analytic theory for avalanches: Dahmen et al 2009

In many cases what is discussed is shear banding in flow.

Transient vs permanent shear bands.

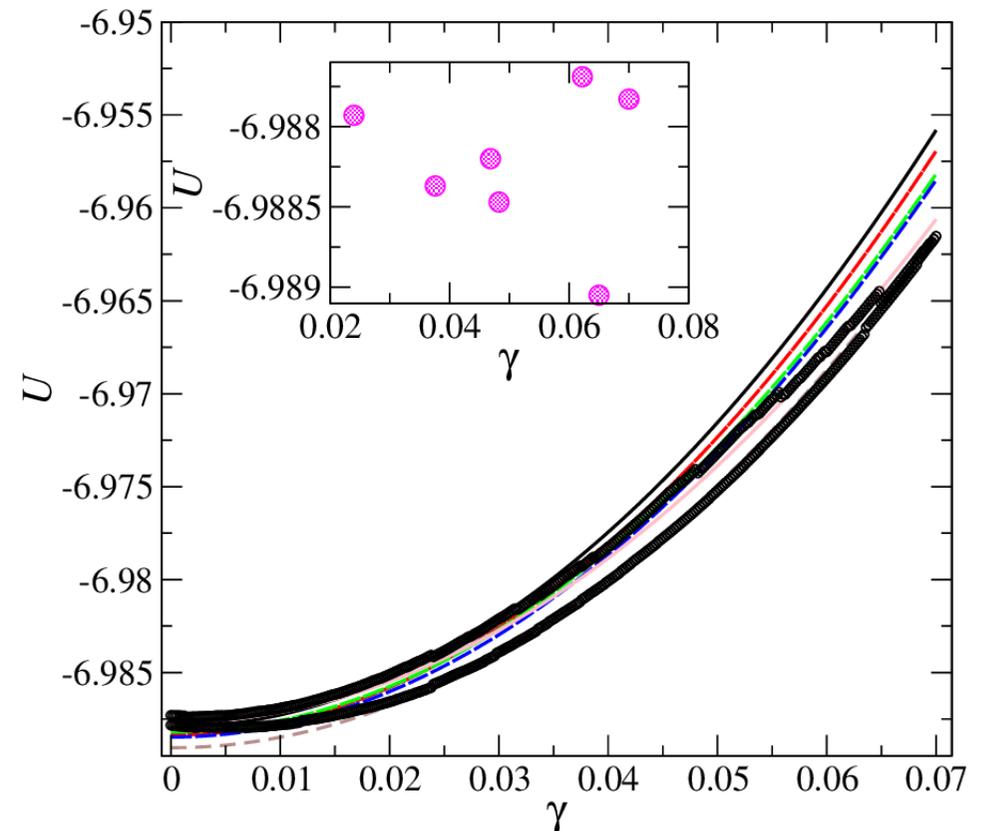
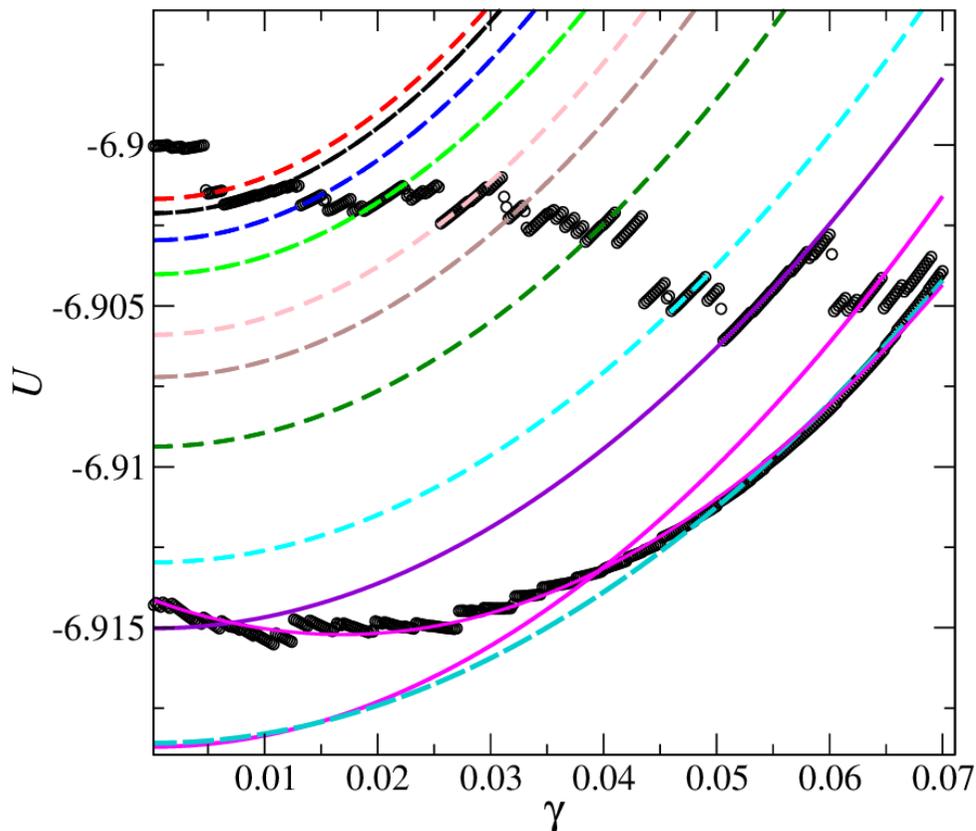
Incorporation of annealing effects in cyclic deformation (e.g. SGR model) leads to shear banding in the sense that is discussed here.

# Annealing Effects under uniform shear

Do annealing effects play a role in yielding under uniform shear?

For poorly annealed glasses, they clearly do. Energy decreases as the system is sheared.

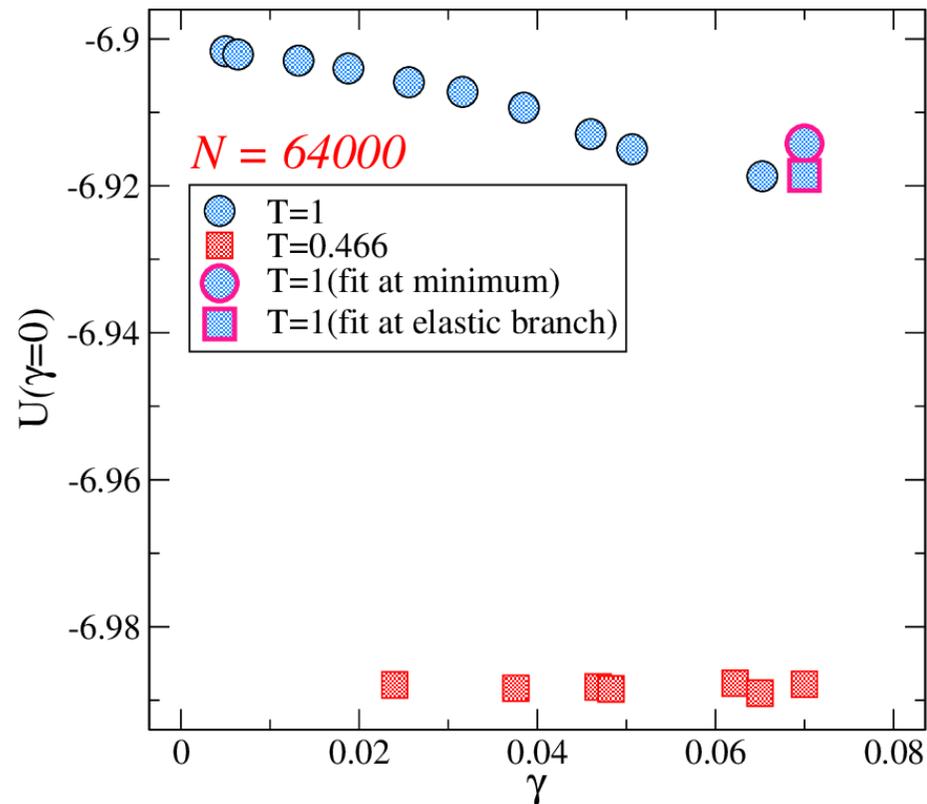
Even when it doesn't obviously, one must consider energy when the system is unloaded.



# Annealing Effects under uniform shear

A harmonic extrapolation back to zero strain reveals that a uniformly sheared glass anneals as a function of strain.

Preliminary analysis.



Ref: Dubey et al 2016

# Summary

The yielding transition appears as a sharp transition associated with an abrupt divergence of the sizes of avalanches, at the yielding transition, for cyclic deformation.

At variance with descriptions predicting divergences upon approach to the transition.

Strong and significant annealing effects before and after the yielding transition.

Annealing under cyclic deformation an interesting phenomenon to understand better.

The role of annealing not prominently treated in approaches to understanding yielding, but various threads in the literature which acknowledge its role.

Observation of shear banding above the yielding point. How do we understand them? Role of annealing. Connection to shear banding in flow.

Description of yielding as a mechanical instability. Integrating with other aspects of glass physics?