

## Ultracold Atoms

### Exercise class n°2: Sagnac interferometer

Aurélien Perrin

We consider an atom interferometer, linked to a reference frame  $\mathcal{R}$  in uniform rotation with respect to a Galilean reference frame  $\mathcal{R}'$ . The rotation vector  $\boldsymbol{\Omega}$  is assumed along the  $Oz$  axis of a reference frame linked to the reference frame  $\mathcal{R}$ , the atoms evolving in the plane  $(xOy)$ . We will evaluate the phase shift of the wave function of a particle between the two arms of the interferometer due to the rotation of  $\mathcal{R}$  with respect to  $\mathcal{R}'$ . For this purpose, we consider a particle of mass  $m$ , free in  $\mathcal{R}'$ . We note  $\mathbf{r}$ ,  $\mathbf{v}$  and  $\mathbf{p}$  (respectively  $\mathbf{r}'$ ,  $\mathbf{v}'$  and  $\mathbf{p}'$ ) its position, velocity and momentum in  $\mathcal{R}$  (respectively  $\mathcal{R}'$ ). We recall:

$$x = x' \cos(\Omega t) + y' \sin(\Omega t) \quad y = -x' \sin(\Omega t) + y' \cos(\Omega t)$$

## 1 Physics in the rotating frame

1. Give the expression linking  $\mathbf{v}'$ ,  $\mathbf{v}$  and  $\mathbf{r}$ , then recall the Lagrangian  $L'(\mathbf{r}', \mathbf{v}')$  of the particle in  $\mathcal{R}'$  (gravity will be neglected).
2. We introduce the Lagrangian  $L(\mathbf{r}, \mathbf{v}) = L'(\mathbf{r}', \mathbf{v}')$ . Write explicitly  $L(\mathbf{r}, \mathbf{v})$  in terms of  $\mathbf{r}$ ,  $\mathbf{v}$  and  $\boldsymbol{\Omega}$ .
3. It should be recalled that in Lagrangian formalism, the momentum is defined by  $p_i = \partial L / \partial \dot{x}_i$ .
  - (a) Determine the momentum  $\mathbf{p}$  as a function of  $\mathbf{r}$  and  $\mathbf{v}$ . Does this result evoke an analogy with another physical situation?
  - (b) Compare  $\mathbf{p}$  and  $\mathbf{p}'$ , and deduce that the Lagrangian  $L$  is a constant of the motion.
4. Show that the hamiltonian  $H(\mathbf{r}, \mathbf{p})$  of the particle (defined by  $H = \mathbf{p} \cdot \mathbf{v} - L$  is written:

$$H = \frac{\mathbf{p}^2}{2m} - \boldsymbol{\Omega} \cdot \mathbf{L} \quad (1)$$

where  $\mathbf{L} = \mathbf{r} \times \mathbf{p}$  is the angular momentum of the particle.

## 2 Wavefunction propagation

We note  $\Psi(\mathbf{r}, t)$  the wave function associated with the particle. It is assumed that at  $t = 0$ , this wave function is that of a plane wave of momentum  $\mathbf{p}_0$  :  $\Psi(\mathbf{r}, 0) = \exp(i\mathbf{r} \cdot \mathbf{p}_0)$ . If the Lagrangian is quadratic in velocity and position, we admit that the wavefunction at a point  $B$  at time  $T$  is given by:

$$\Psi(\mathbf{r}_B, T) = G(T) \Psi(\mathbf{r}_A, 0) \exp\left(\frac{i}{\hbar} S(\mathbf{r}_B, T; \mathbf{r}_A, 0)\right) \quad (2)$$

where the point  $\mathbf{r}_A$  is such that the trajectory passing in  $\mathbf{r}_A$  at the time  $t = 0$  with the momentum  $\mathbf{p}_0$  passes in  $\mathbf{r}_B$  at the time  $T$ . The function  $G(T)$  depends only on  $T$  (not on the path traveled), and  $S$  is the classical action calculated on the path.

1. Calculate  $\dot{\mathbf{r}}$  and  $\dot{\mathbf{p}}$ . We recall the Hamilton-Jacobi equations :

$$\dot{\mathbf{r}} = \frac{\partial H}{\partial \mathbf{p}} \quad \dot{\mathbf{p}} = -\frac{\partial H}{\partial \mathbf{r}}.$$

2. Integrate the equation over  $\mathbf{p}$ , and then the one over  $\mathbf{r}$ , with initial conditions the particle in  $\mathbf{r}_A$  with momentum  $\mathbf{p}_0$ .
3. Show that the classical action  $S$ , in first order in  $\Omega$ , is given by:

$$S = S_0 + m\Omega(x_A y_B - y_A x_B) \quad \text{with} \quad S_0 = \frac{m}{2T} ((x_B - x_A)^2 + (y_B - y_A)^2). \quad (3)$$

4. Show that at the zeroth order, we have:

$$x_A y_B - y_A x_B = \mathbf{u}_z \cdot \int_0^T \mathbf{r}(t) \times \mathbf{v}(t) dt$$

where the integral is taken along the  $AB$  trajectory.

### 3 Phase shift

We can then show that the wavefunction in  $\mathbf{r}_B$  at  $T$  is written:

$$\Psi(\mathbf{r}_B, T) = \Psi^{(0)}(\mathbf{r}_B, T) \exp\left(\frac{i}{\hbar} m\Omega \cdot \int_0^T \mathbf{r}(t) \times \mathbf{v}(t) dt\right) \quad (4)$$

where  $\Psi^{(0)}(\mathbf{r}_B, T)$  represents the wavefunction for  $\Omega = 0$ .

*This point is more difficult to show than it seems, and requires careful inspection of the trajectories.*

1. The interferometer is assumed to be in the shape of an  $ACBD$  quadrilateral. A particle from point  $A$  can therefore take two paths ( $ACB$  or  $ADB$ ). Deduce from the previous questions the phase shift of the particle wavefunction between the two arms of the interferometer according to  $\Omega$  and the area of the interferometer.
2. In a 1996 experiment in California, Mark Kasevich and his group demonstrated this effect on cesium atoms (mass  $m = 2.2 \times 10^{-25}$  kg). The velocity of the atomic beam used is 300 m/s. At point  $A$ , the beam is divided into two components. One ( $AD$ ) is transmitted without modification; the other ( $AC$ ) is deflected by interaction with a laser beam. More precisely, the transverse velocity of the deflected atoms is modified by two recoil velocities, using a stimulated Raman pulse. The same phenomenon is used at points  $C$  and  $D$  to deflect the atoms again, and at  $B$  to close the interferometer.
  - (a) The lasers wavelength is  $\lambda = 852$  nm. Calculate the recoil velocity.
  - (b) We give  $AD = CB = 1$  m. Calculate the area of the interferometer.
  - (c) Calculate the phase shift related to the rotation of the earth and compare it to the result found by Mr. Kasevich: 4.2 radian.