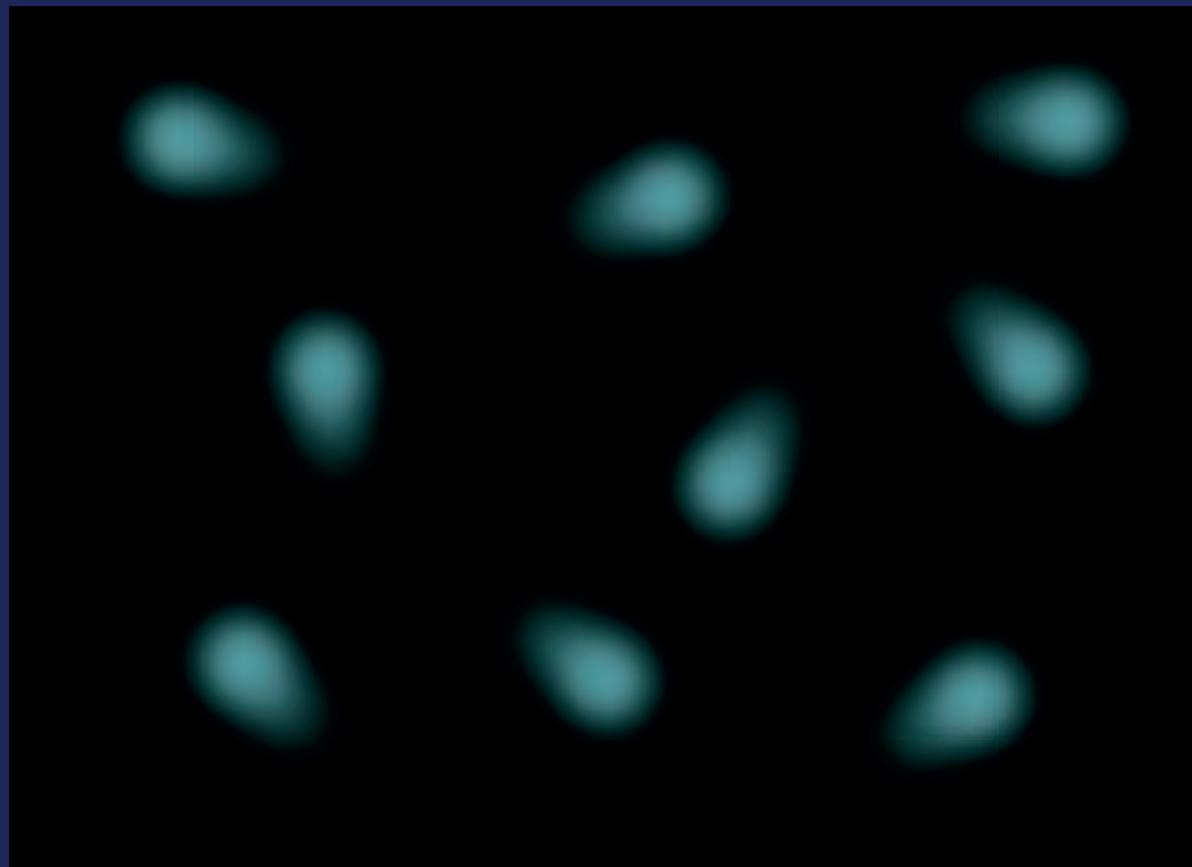
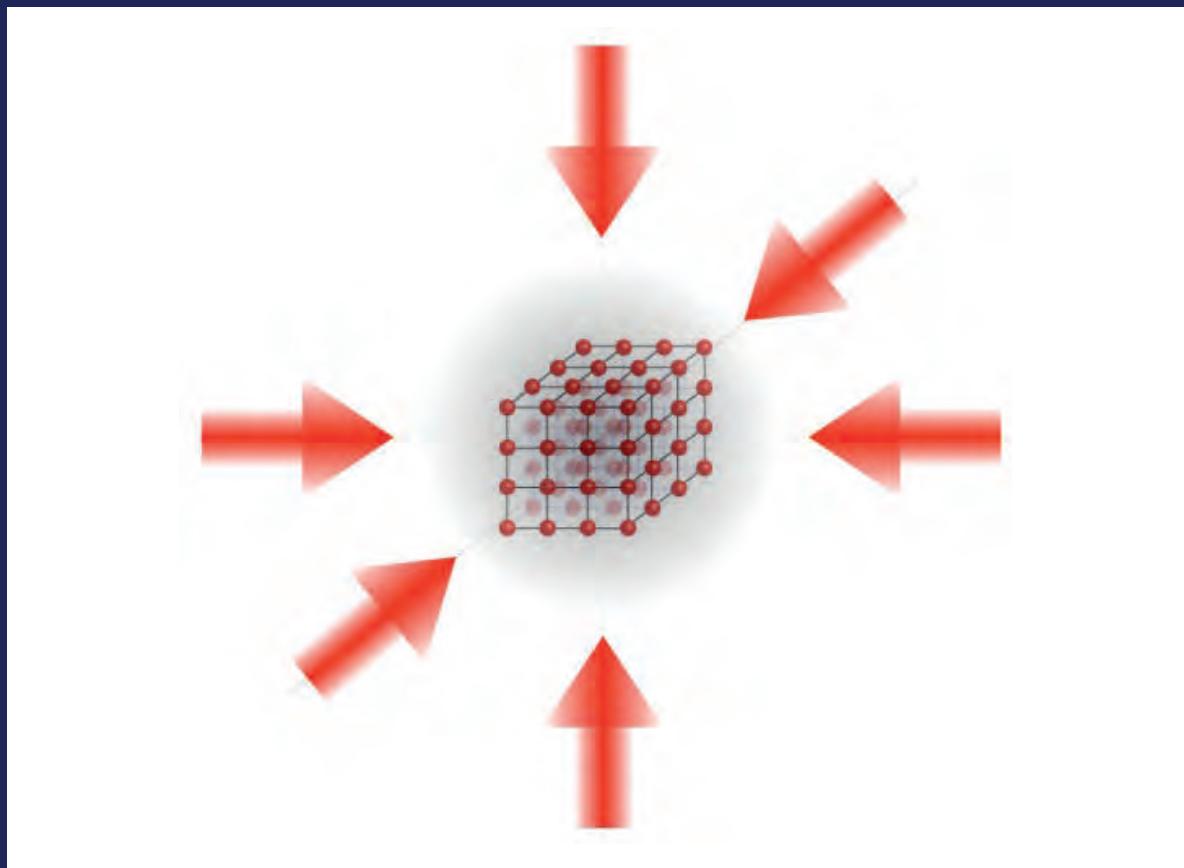


# Building quantum systems from scratch

Tilman Esslinger ETH Zürich

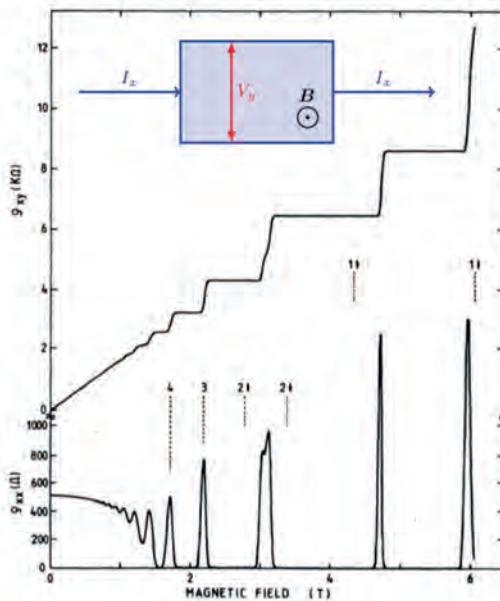
Funding: ETH, EU (ERCadv TransQ, SIQS, TherMiQ, QUIC, ColOpt),  
NCCR QSIT, SNF  
[www.quantumoptics.ethz.ch](http://www.quantumoptics.ethz.ch)





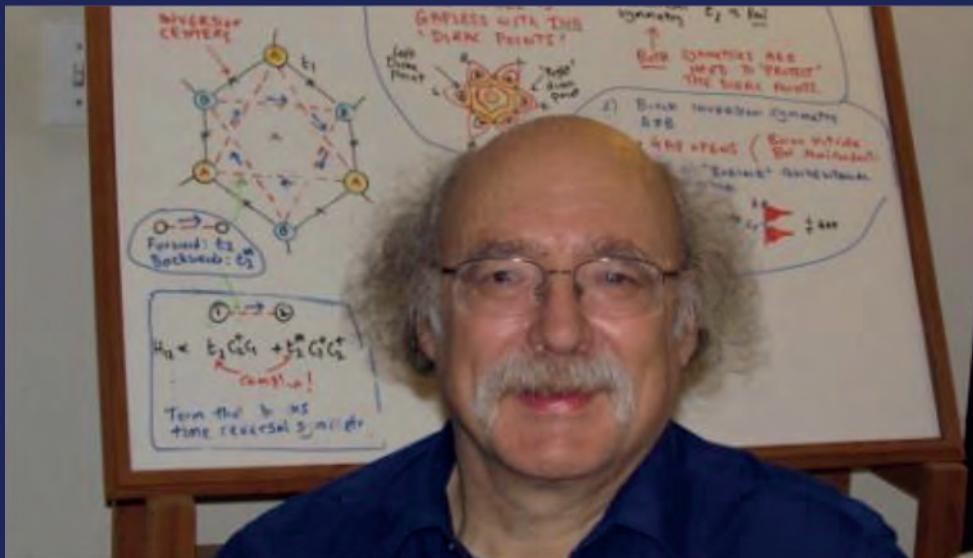


**FIGURE 1.** Classification topologique obtenue en comptant le nombre de "poignées" d'un objet. Cette classification est robuste, puisqu'un objet ne change pas de classe quand on le déforme légèrement, et elle ne dépend pas d'éventuelles symétries spatiales géométriques de l'objet.

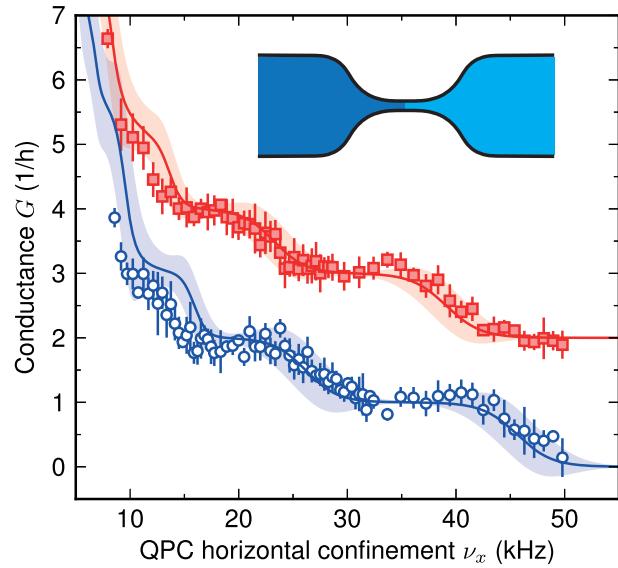


**FIGURE 2.** Exemple de manifestation d'un ordre topologique dans un gaz 2D d'électrons plongé dans un grand champ magnétique : la résistivité de Hall  $\rho_{xy}$  mesurée dans une expérience d'effet Hall quantique présente des plateaux quantifiés quand on varie le champ magnétique ; dans cette mesure, un courant  $I_x$  circule dans un échantillon rectangulaire  $L_x \times L_y$  et on mesure une tension  $V_y$  aux bornes de l'échantillon. Cette quantification est robuste : elle subsiste en présence d'un désordre (modéré) dans l'échantillon. Figure extraite de VON KLITZING (1986).

## Haldane model



## Quantized conductance

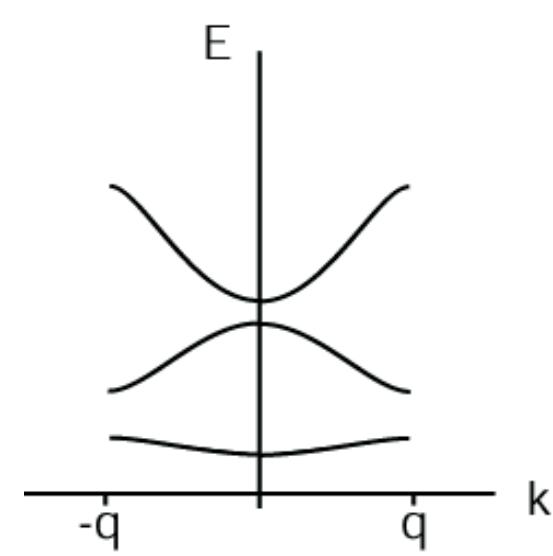


# Where is the physics?

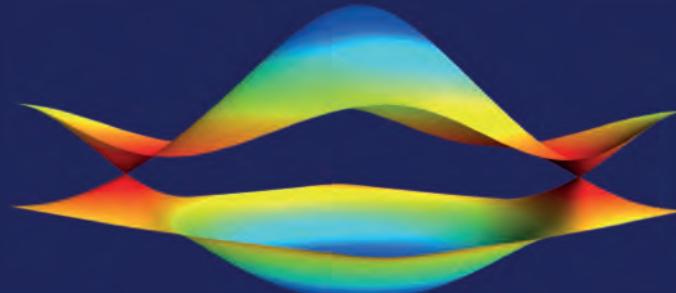
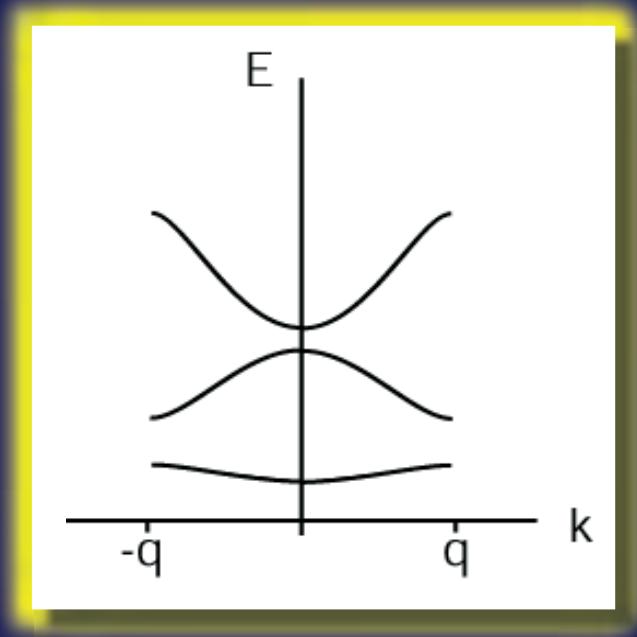
$$H = T + U + V_{\text{trap}}$$

$$H = T + U + V_{\text{trap}}$$

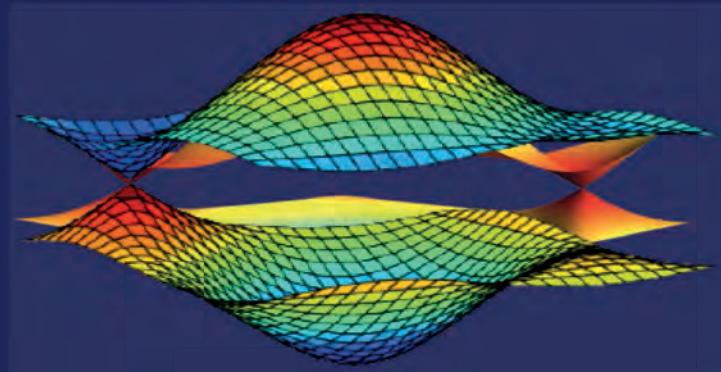
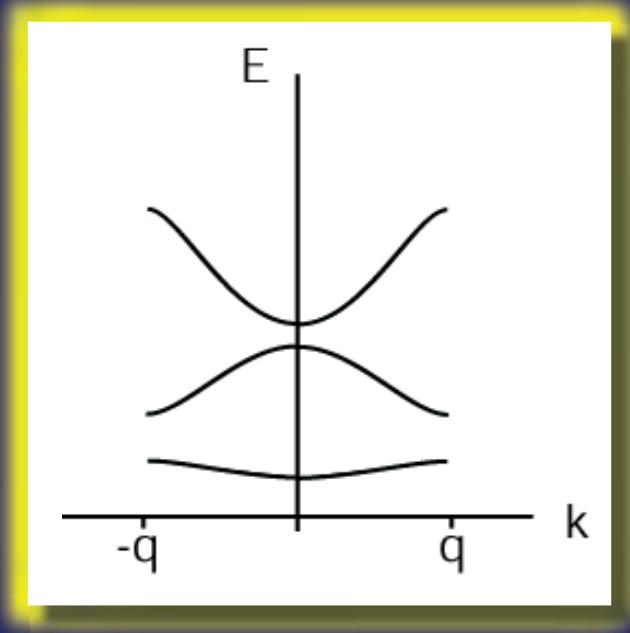
$$H = T + U + V_{\text{trap}}$$



$$H = T + U + V_{\text{trap}}$$

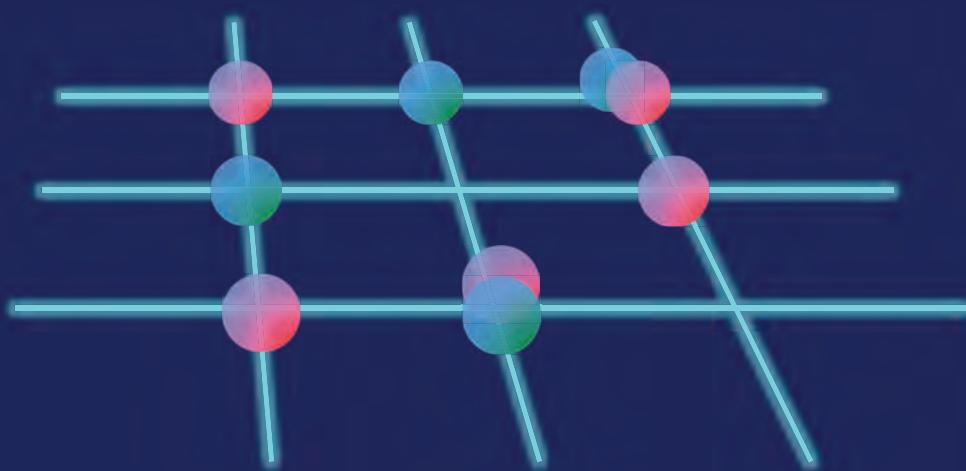


$$H = T + U + V_{\text{trap}}$$

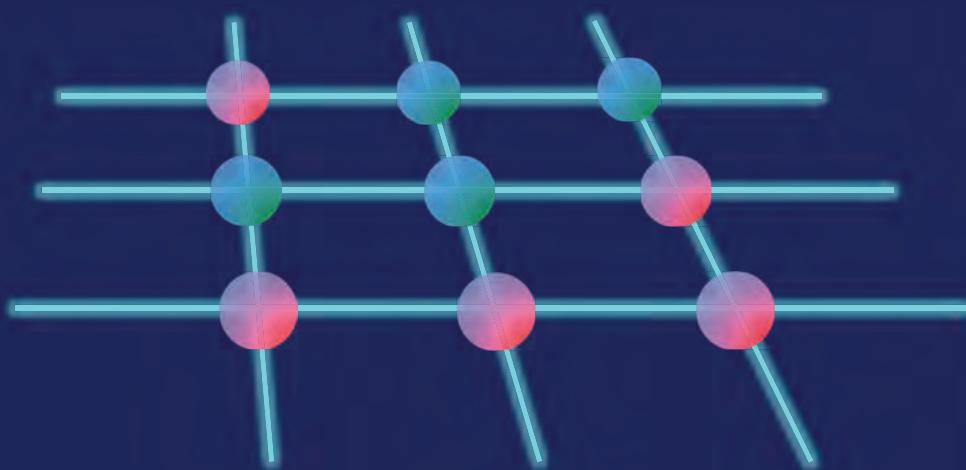


$$H = T + U + V_{\text{trap}}$$

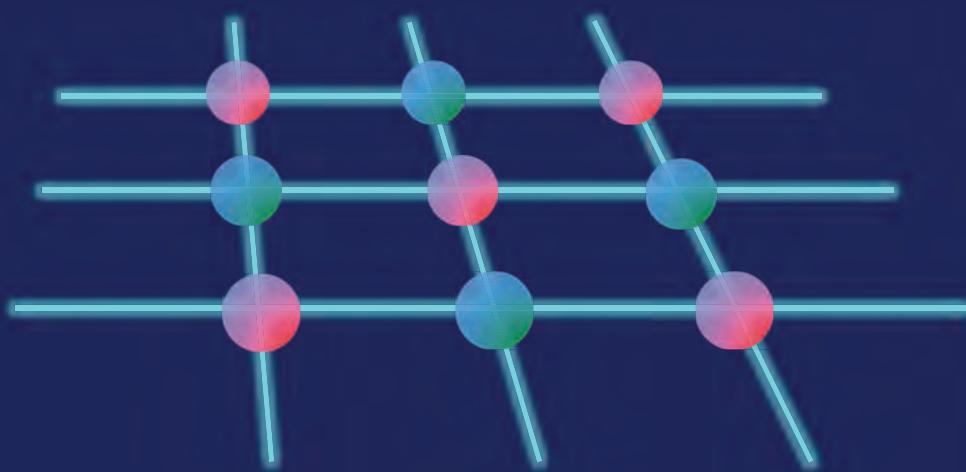
$$H = T + U + V_{\text{trap}}$$



$$H = T + U + V_{\text{trap}}$$

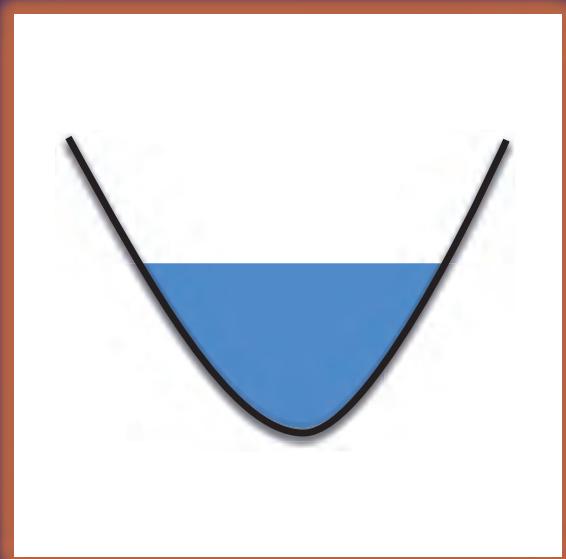


$$H = T + U + V_{\text{trap}}$$

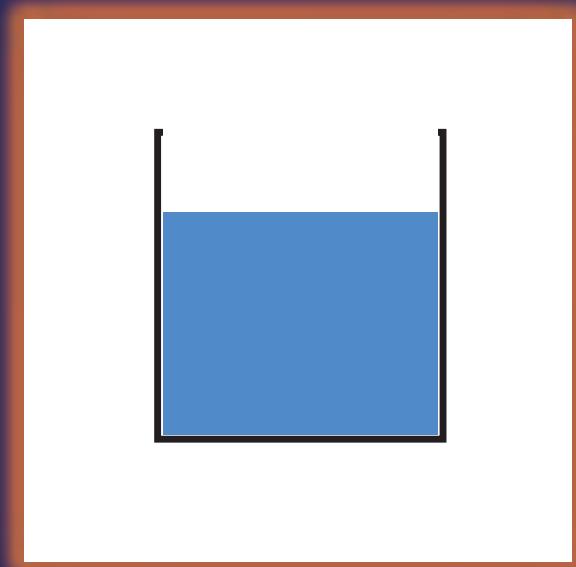


$$H = T + U + V_{\text{trap}}$$

$$H = T + U + V_{\text{trap}}$$

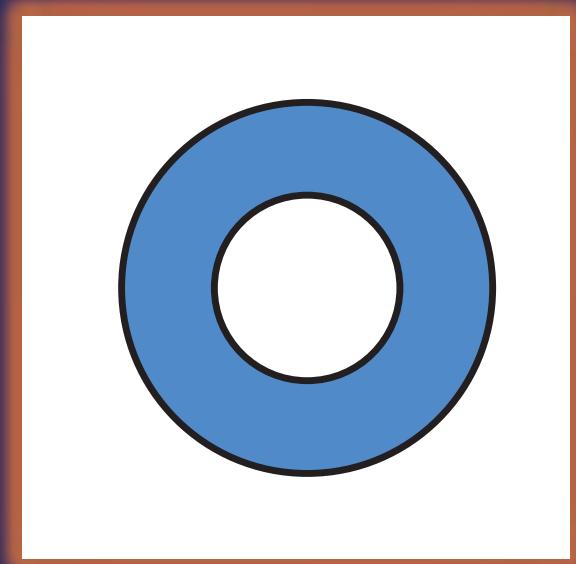


$$H = T + U + V_{\text{trap}}$$



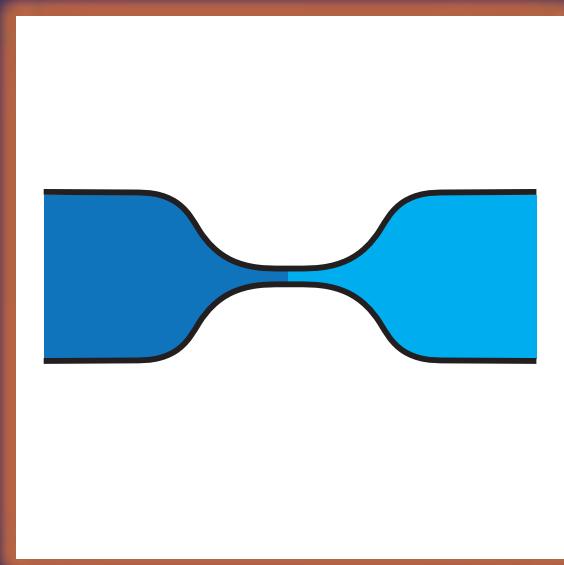
e.g. Cambridge, ENS

$$H = T + U + V_{\text{trap}}$$



e.g. NIST (Campbell), ENS

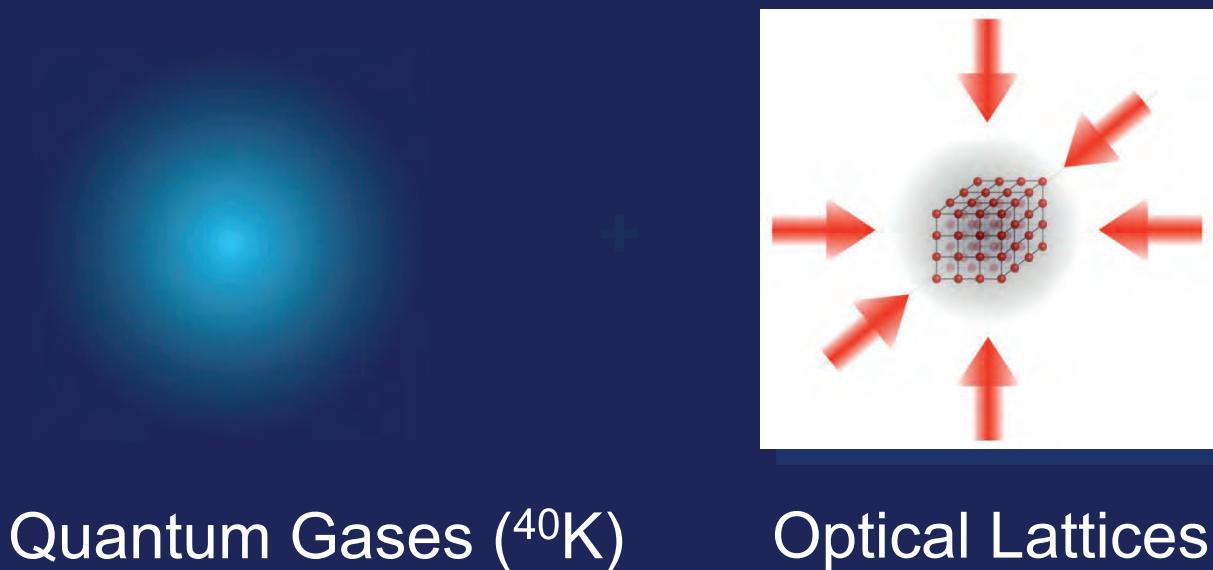
$$H = T + U + V_{\text{trap}}$$



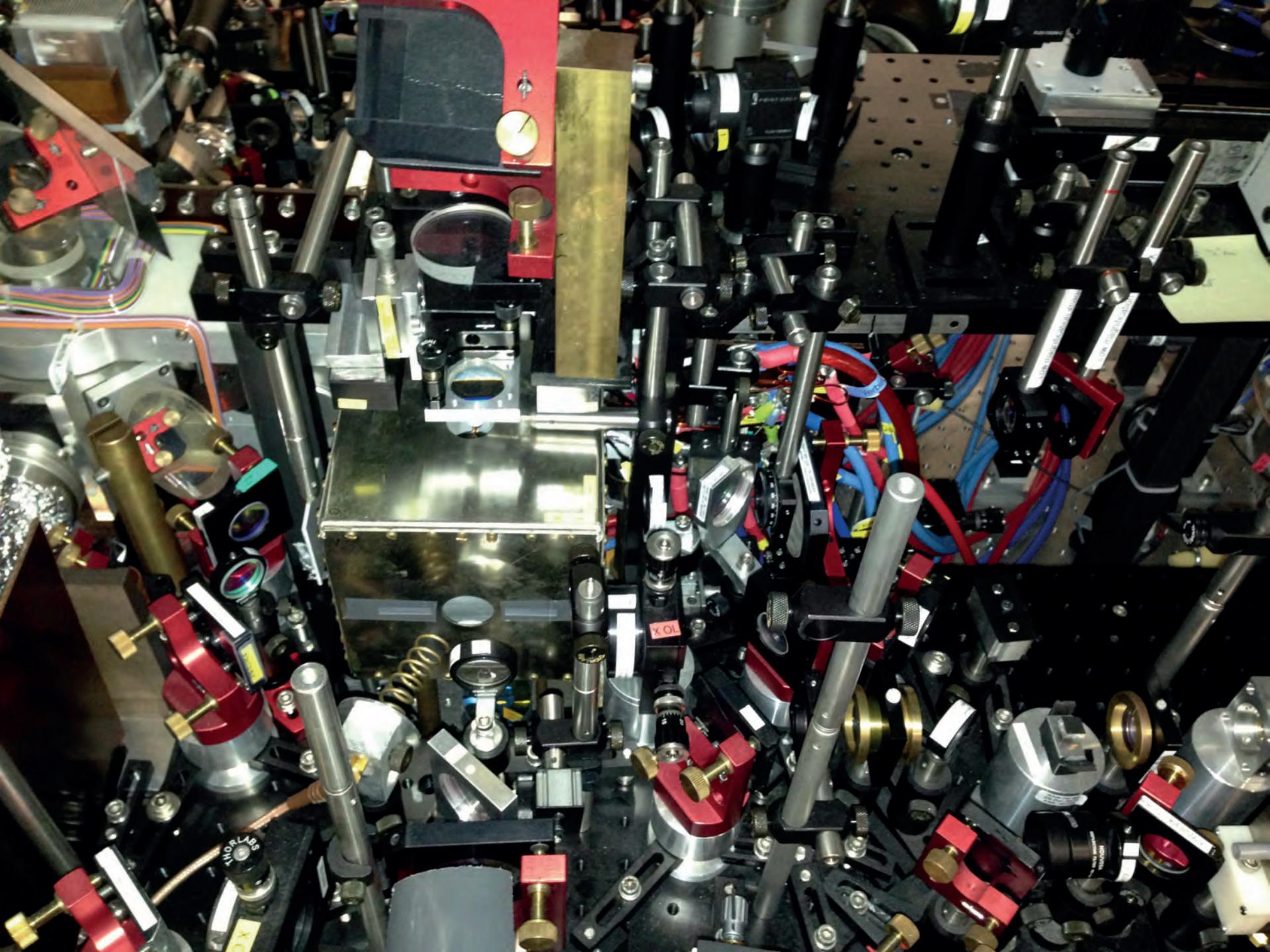
Transport: ETH, JILA,  
Josephson: Heidelberg, Firenze,...



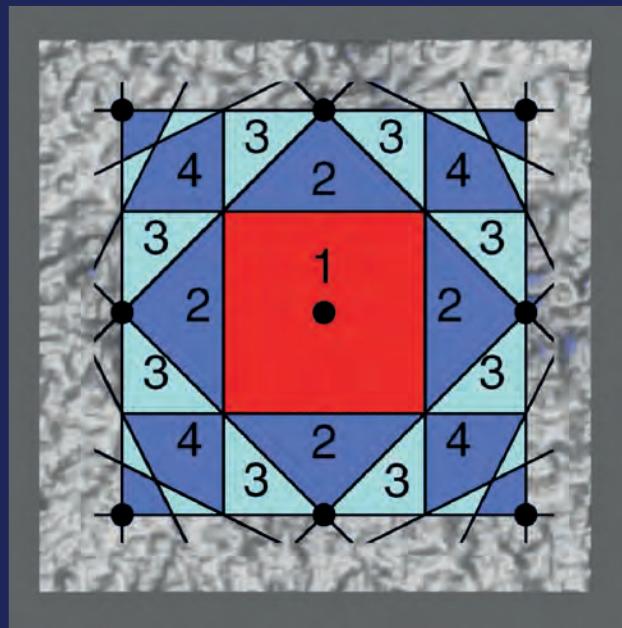
# Building the Hamiltonian



See also: Mainz/Munich, Hamburg, MIT, Illinois, Rice,...

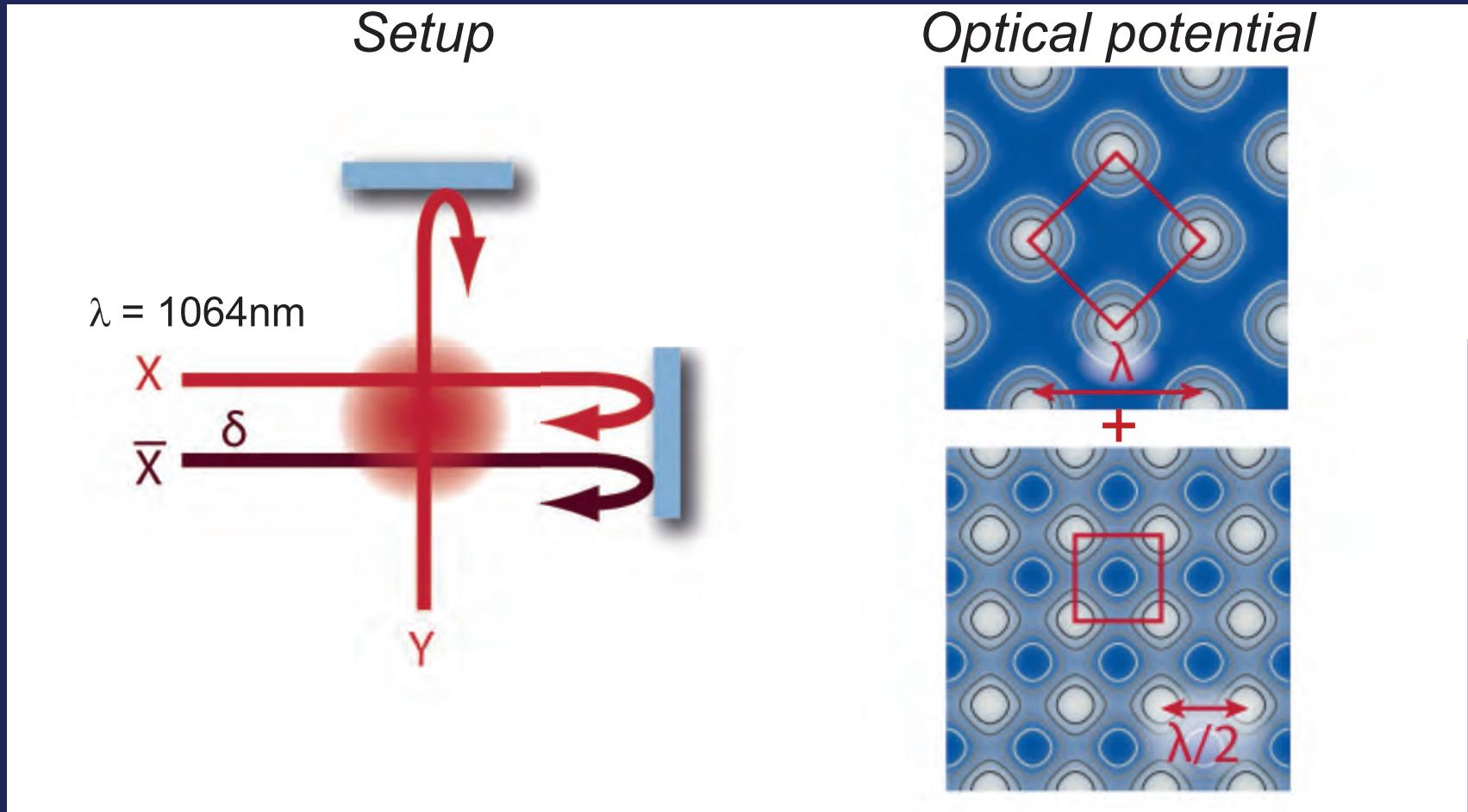


# Simple Measurement...



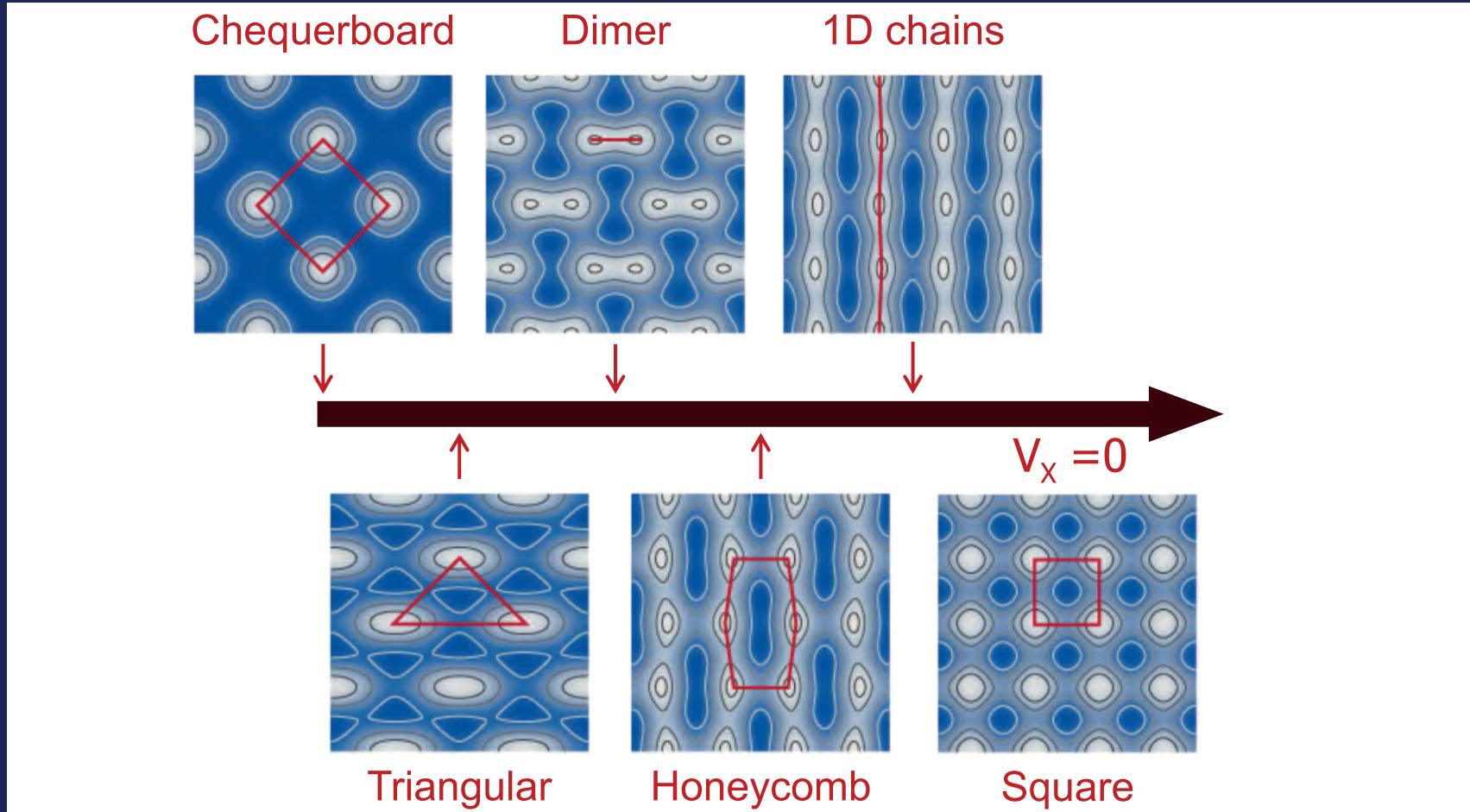
Simple structure...

# Tunable Geometry Optical Lattice

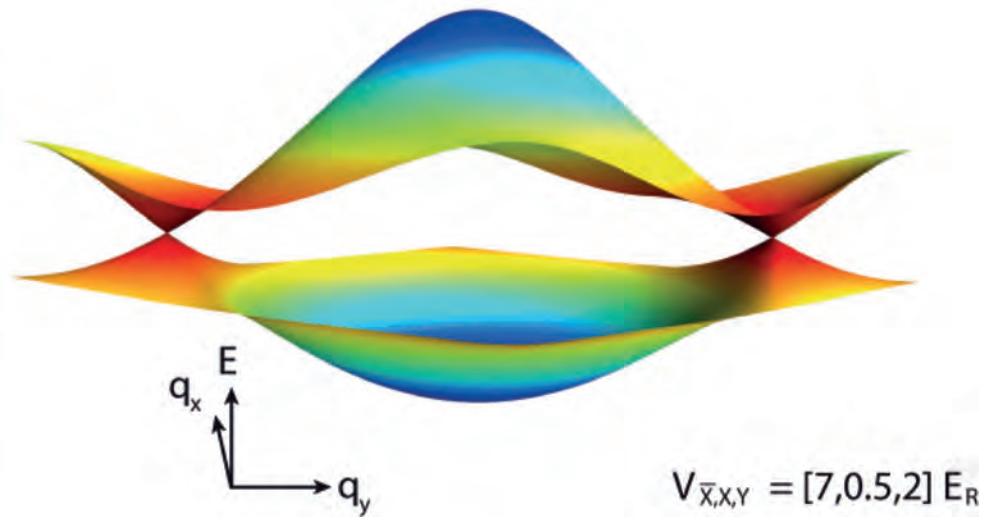
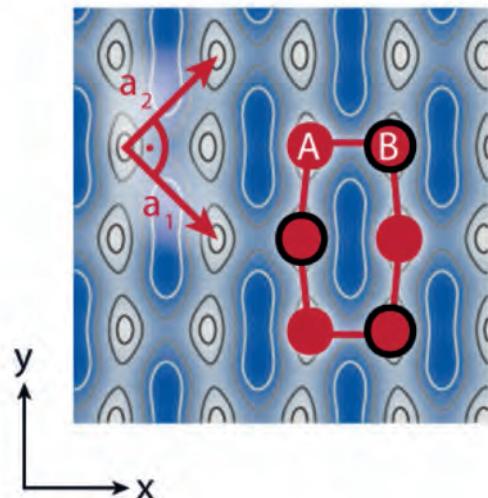


Other complex lattices: NIST, Munich, Hamburg, Berkeley, ...

# Tunable Geometry Optical Lattice



# Honeycomb Lattice



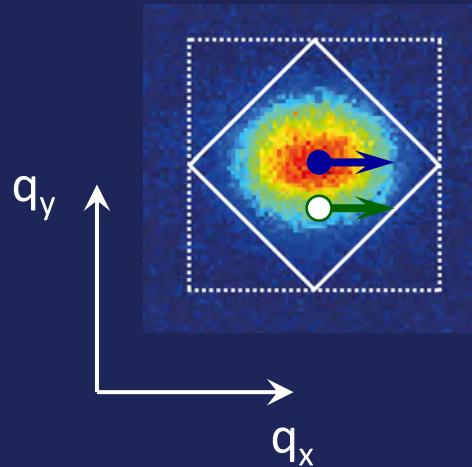
# Probing the Dirac points

vanishing density of states

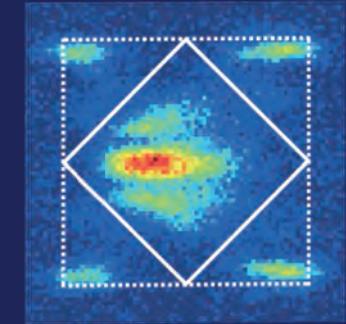
small energy scales

# Bloch oscillation and interband transitions

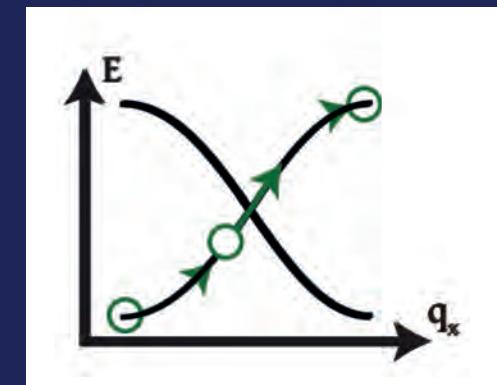
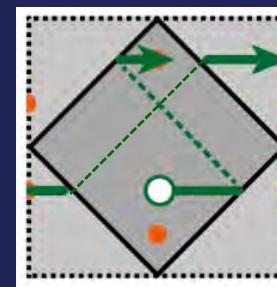
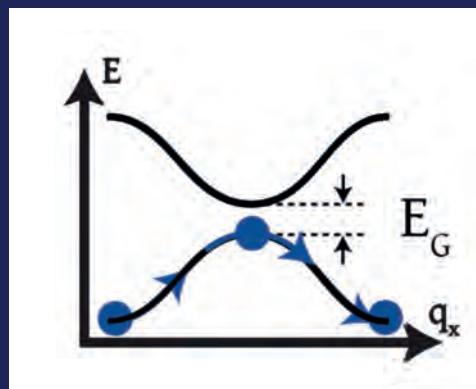
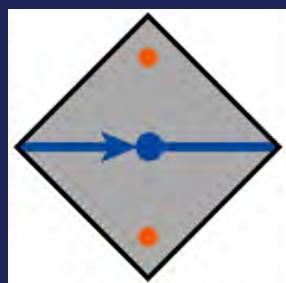
Starting point



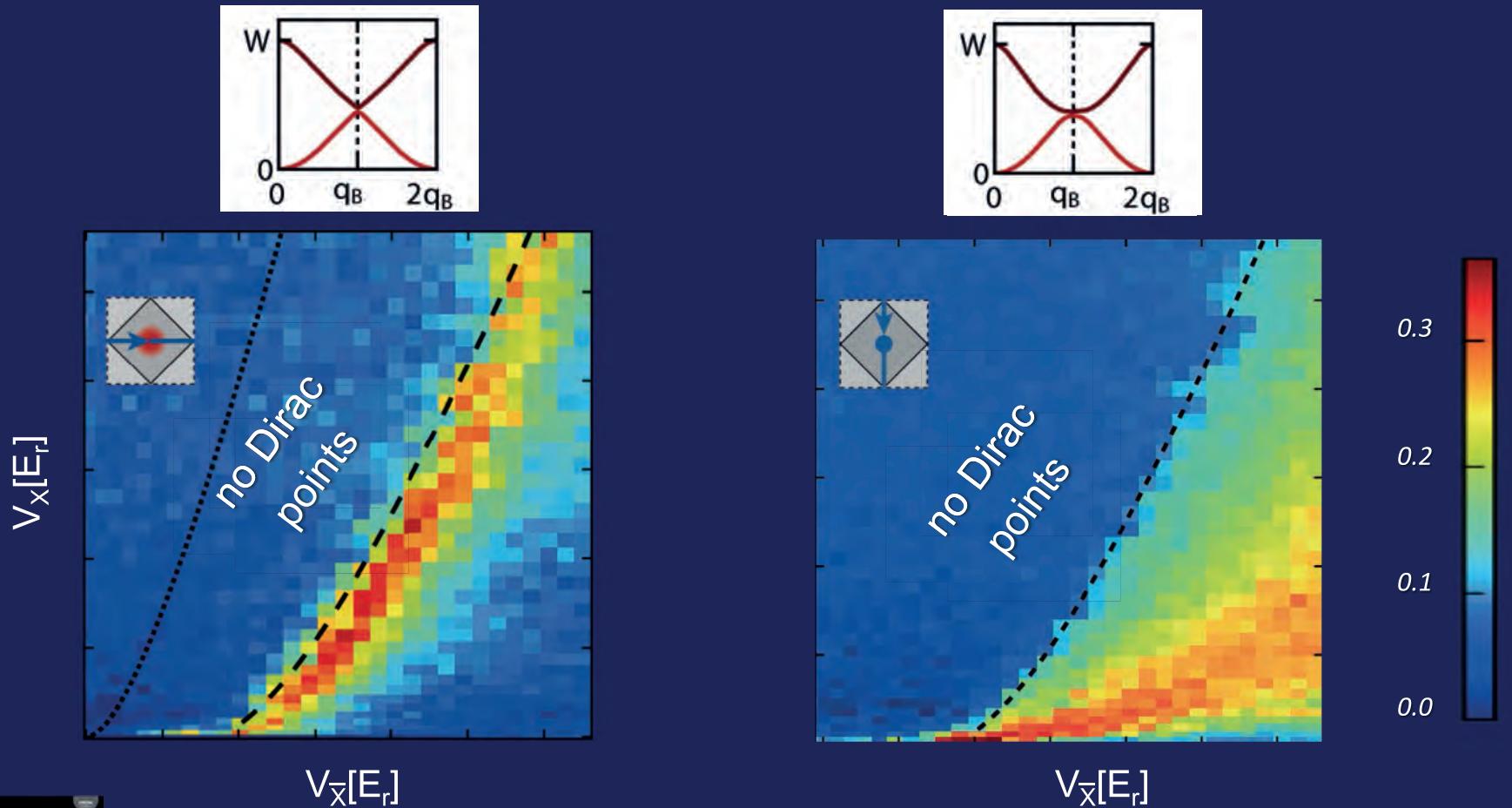
+ magnetic gradient



Transfer to 2nd band



# Touching Dirac points

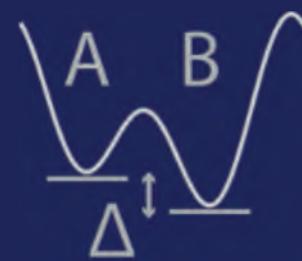


L. Tarruell, D. Greif, T. Uehlinger, G. Jotzu, and T. Esslinger, Nature 483, 302–305 (2012).

Theory, see also: L.-K. Lim, J.-N. Fuchs, G. Montambaux, PRL 108, 175303 (2012)

# Breaking Inversion Symmetry

# Berry curvature

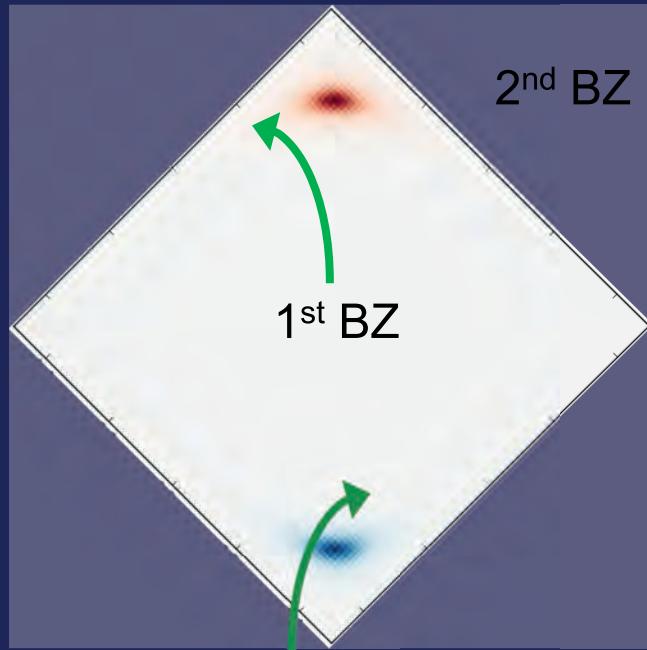


See also: L. Duca, Science 347, 288 (2015)

Review: N. Goldman, G. Juzeliunas, P. Ohberg, I. Spielman, Rep. Prog. Phys. 77, 126401, (2014)

Goldman, N., Cooper, N., & Dalibard, J. (2017). Preparing and Probing Chern Bands with Cold Atoms. doi:10.1017/9781316084366.016

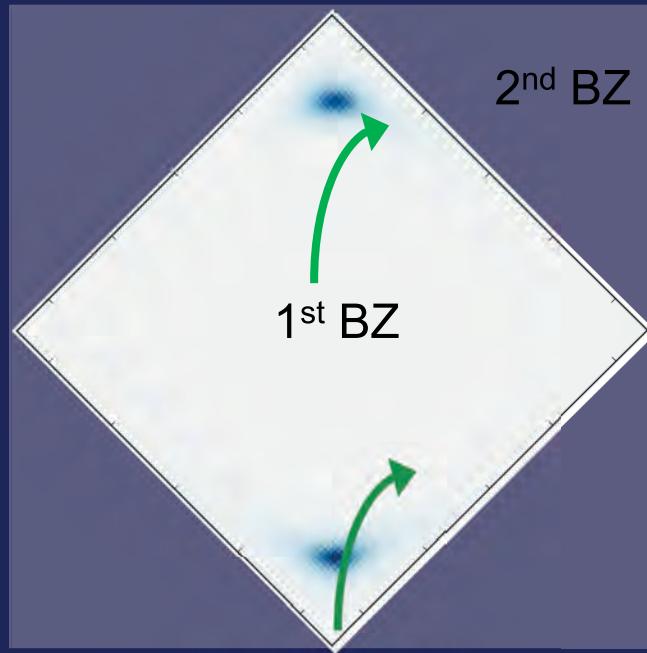
# Berry Curvature and Transverse Drift



$$\dot{\mathbf{r}} = \frac{1}{\hbar} \partial_{\mathbf{k}} \epsilon(\mathbf{k}) - \dot{\mathbf{k}} \times \boldsymbol{\Omega}(\mathbf{k})$$
$$\hbar \dot{\mathbf{k}} = \mathbf{F}(\mathbf{r})$$

Chang and Niu, PRL 75, 1348 (1995)  
Price and Cooper, PRA 85, 033620 (2012)

# Berry Curvature and Transverse Drift

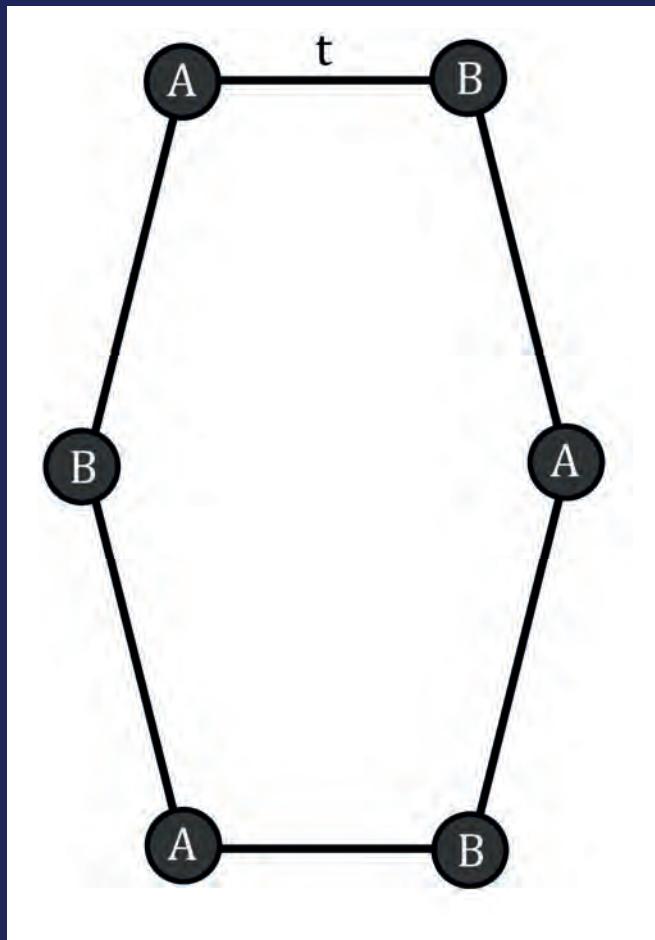


Like a Hall current

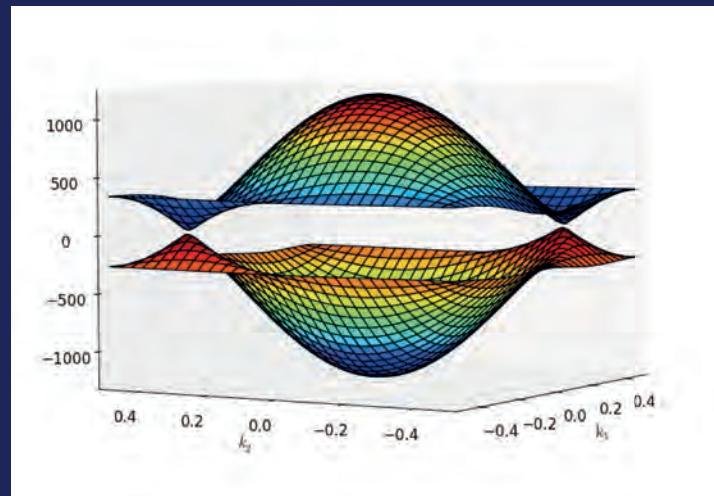
# Topological Haldane model

Proposal for Quantum Hall Effect *without* magnetic field!  
Haldane, PRL **61**,2015-2018 (1988)

# Topological Haldane model



Start from a honeycomb lattice

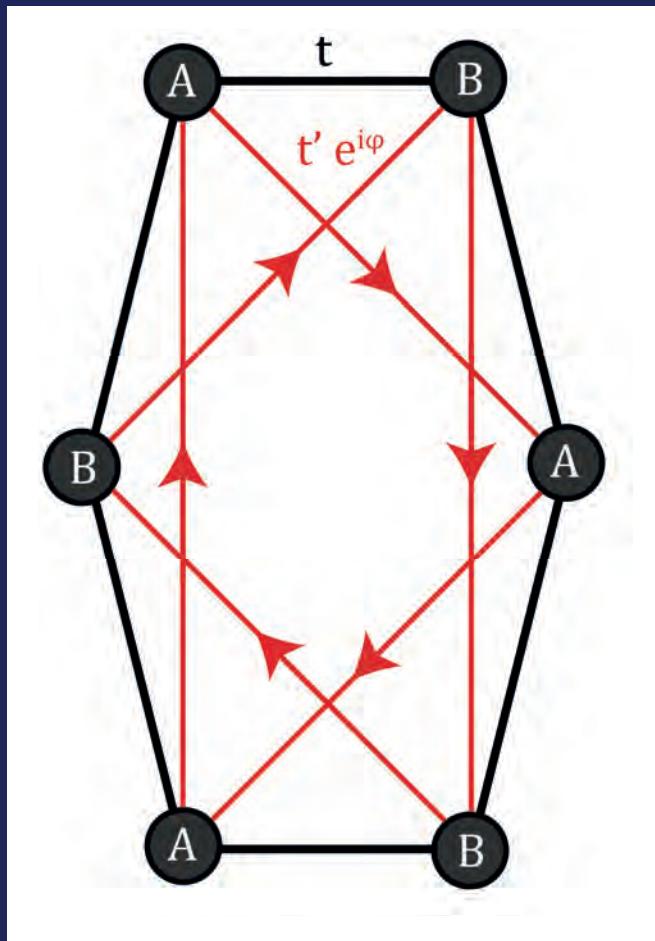


inversion and time-reversal symmetry

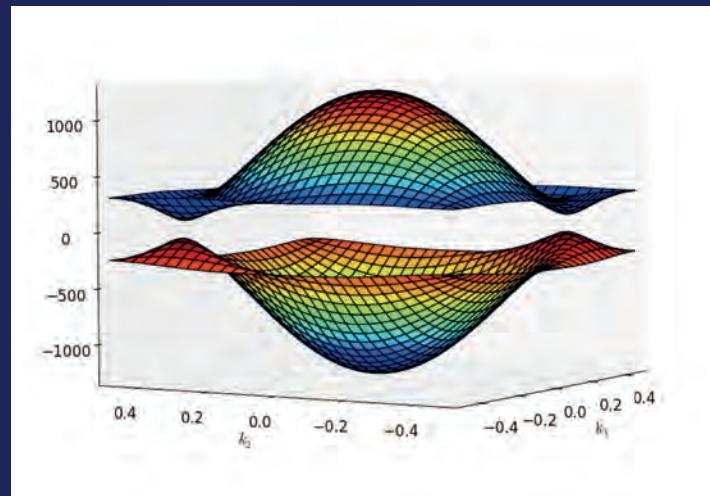
An aerial photograph of a large, open urban square. In the center is a prominent, multi-tiered circular fountain or monument. The square is paved with light-colored stone and has several green trees and small shrubs planted along its edges. Surrounding the square are numerous multi-story buildings, mostly residential houses with red-tiled roofs. The overall scene is a mix of historical architecture and modern urban planning.

**Topological Haldane model  
break time-reversal symmetry**

# Topological Haldane model



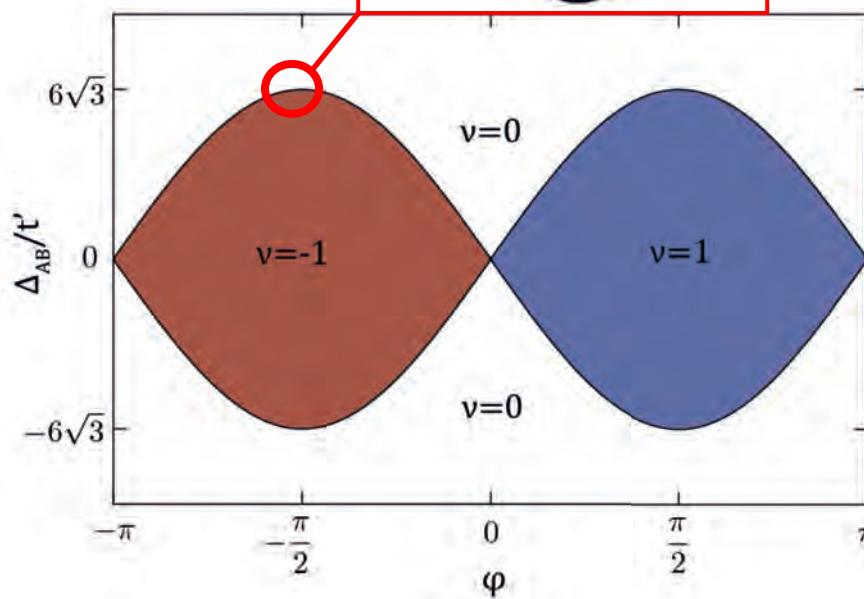
break time-reversal symmetry with  
complex next-nearest neighbour tunnellings



→ Topological Chern insulator, with non-zero Hall conductance

# Topological Haldane model

Inversion symmetry



Time-reversal symmetry

Haldane, PRL 61, 2015-2018 (1988)

How?



VOLUME 61, NUMBER 18

PHYSICAL RE

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geometrical constant of order unity, and  $g$  is the Landé  $g$  factor for the electrons.

While the particular model presented here is unlikely to be directly physically realizable, it indicates that, at least in principle, the QHE can be placed in the wider context of phenomena associated with broken time-reversal invariance, and does not necessarily require external magnetic fields, but could occur as a consequence of magnetic ordering in a quasi-two-dimensional system.

Haldane, PRL 61,2015-2018 (1988)

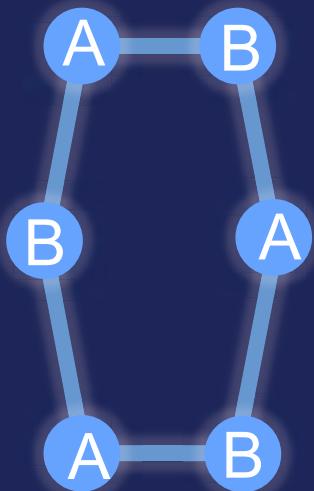
# Breaking time-reversal symmetry

Proposal for Photovoltaic Hall effect in graphene

T. Oka und H. Aoki, PRL **79**, 081406 (2009)



# Breaking time-reversal symmetry



Other proposals to realize topological Hamiltonians:

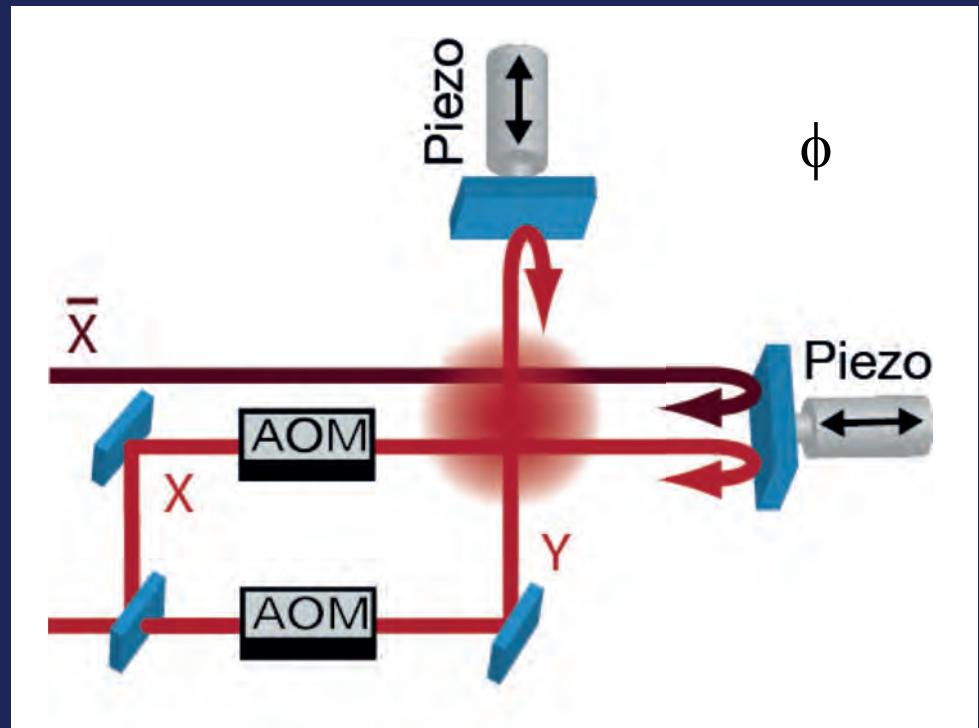
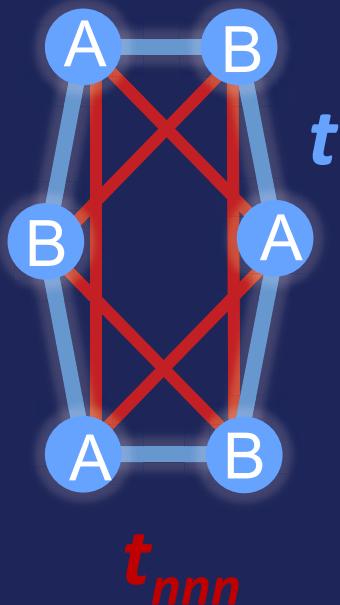
T. Kitagawa et al., Phys. Rev. B 82, 235114 (2010)

P. Hauke et al., Phys. Rev. Lett 109, 145301 (2012)

Realisation in photonic system: Rechtsman et. al Nature 496, 196–200 (2013)

# Breaking time-reversal symmetry

## Lattice Shaking



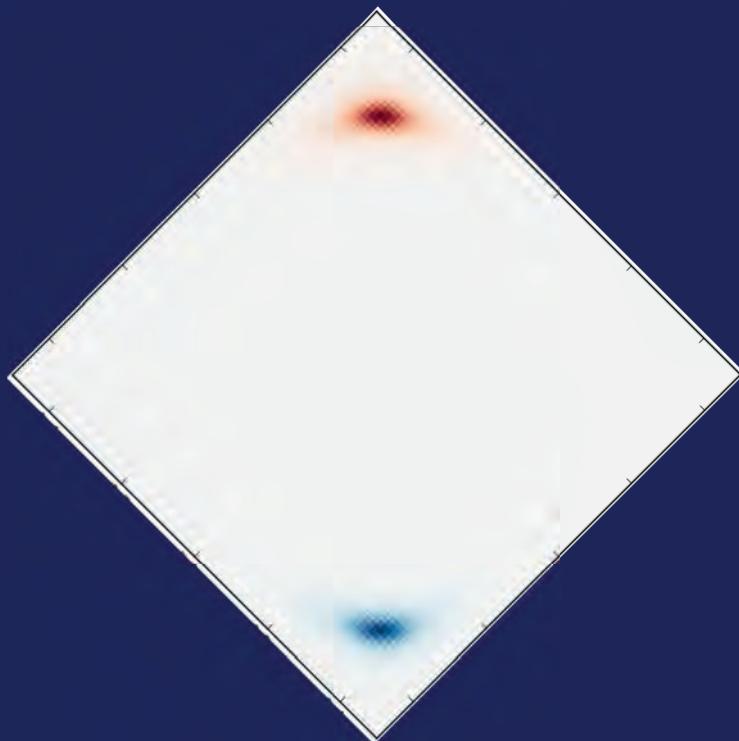
Lattice shaking: Pisa — Lignier, PRL 99, 220403 (2007)

Hamburg/Barcelona — Struck, Science 333, 996-9 (2011), PRL 108, 225304 (2012)

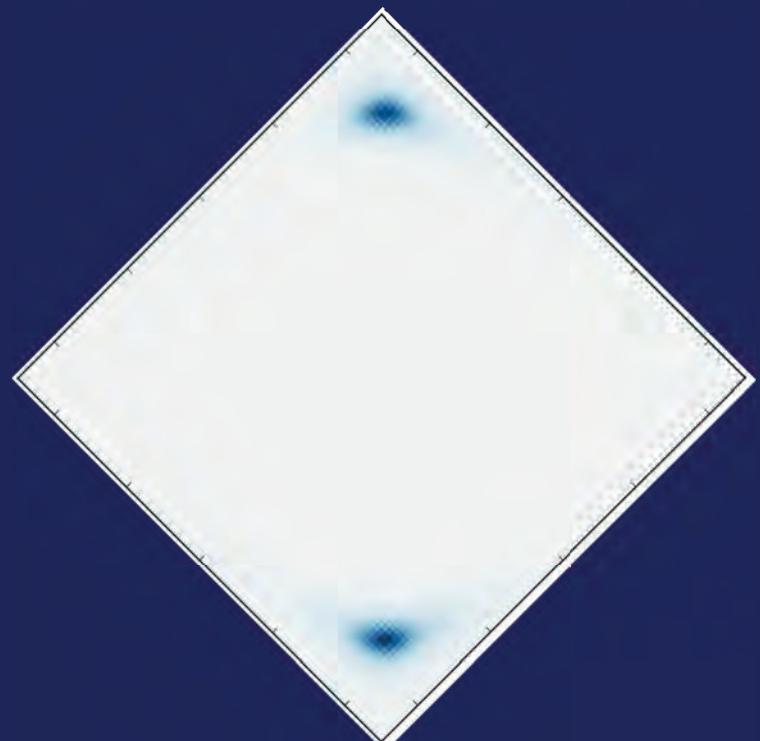
Chicago — Parker, Nat. Phys. 9, 769-774 (2013)

# Berry Curvature

Trivial band insulator



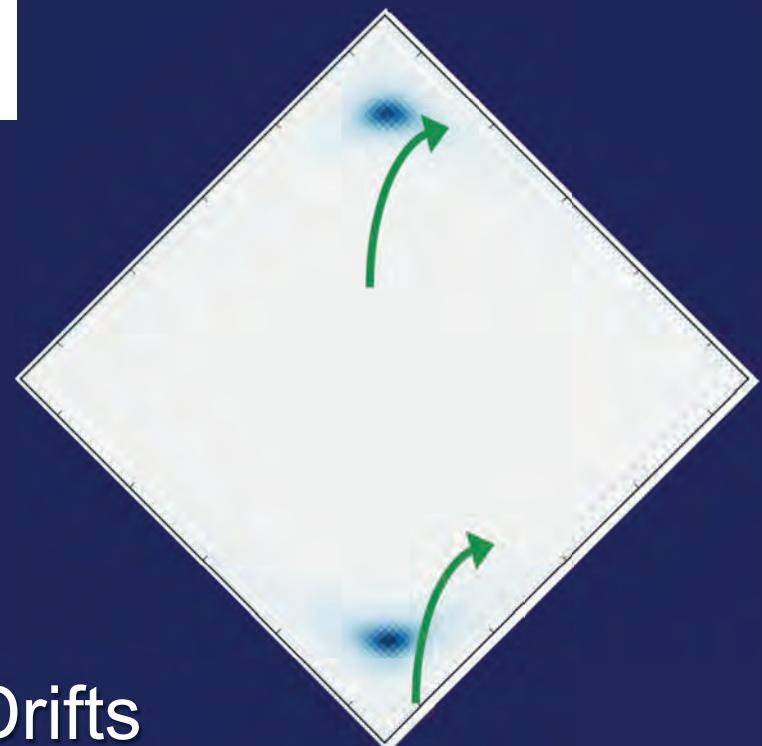
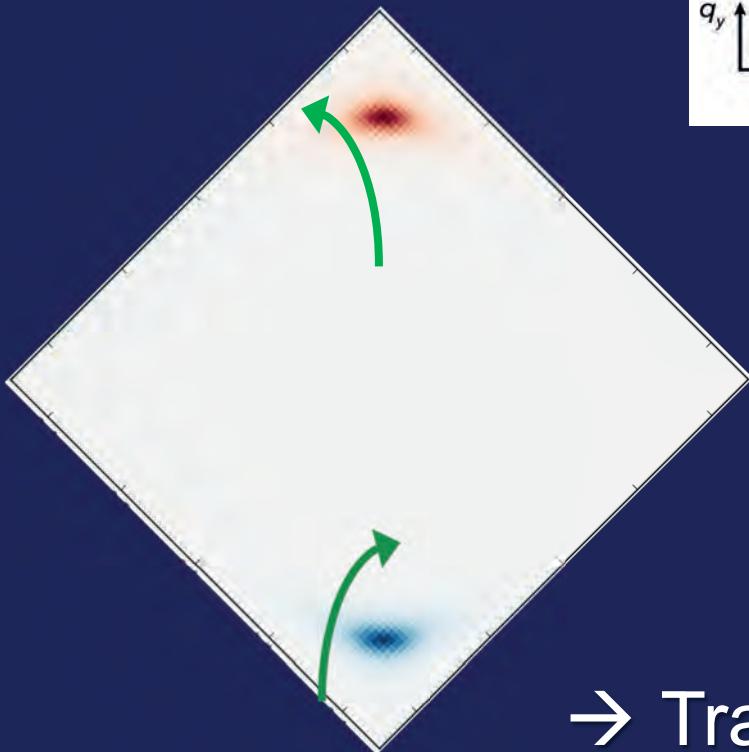
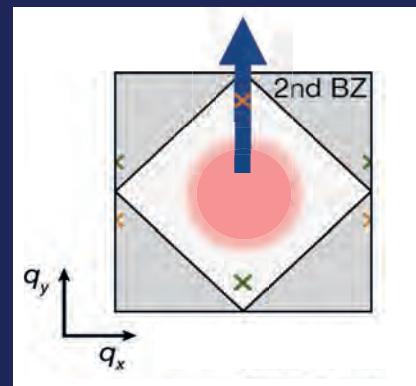
Chern insulator



Chern number 0

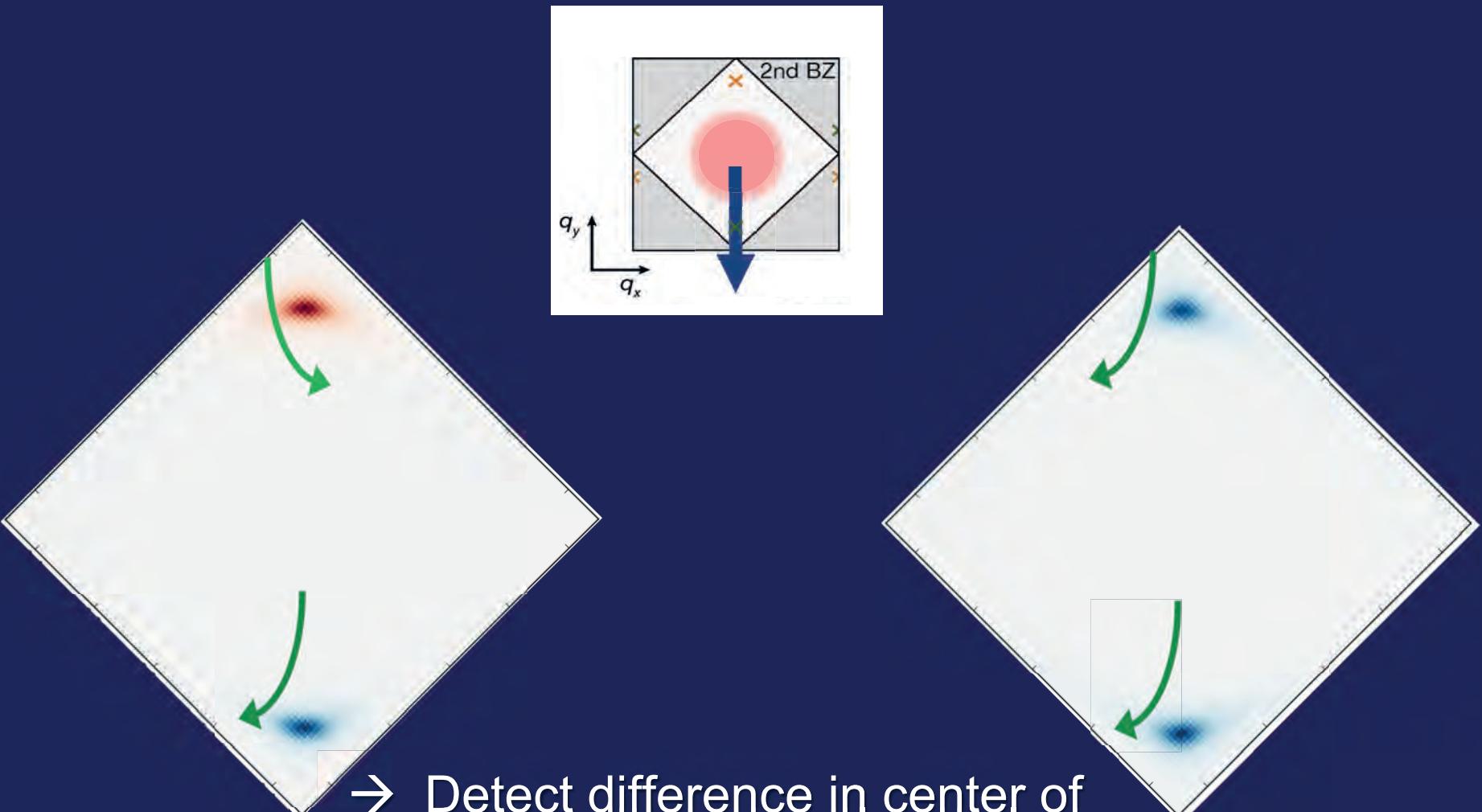
Chern number -1

# Berry Curvature - Measurement



→ Transverse Drifts

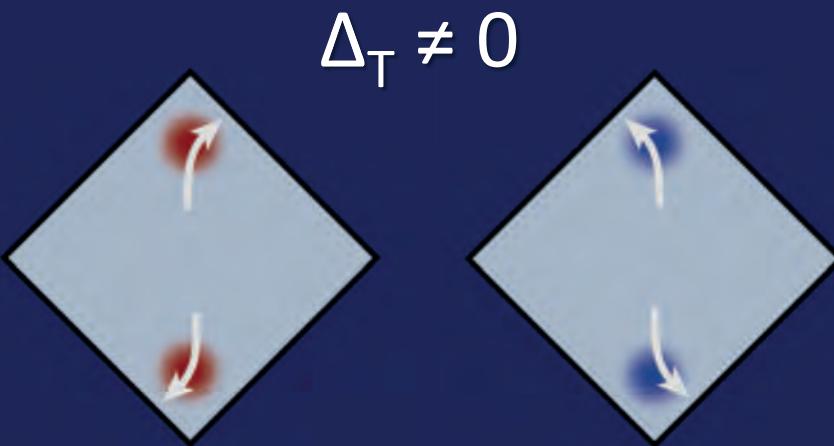
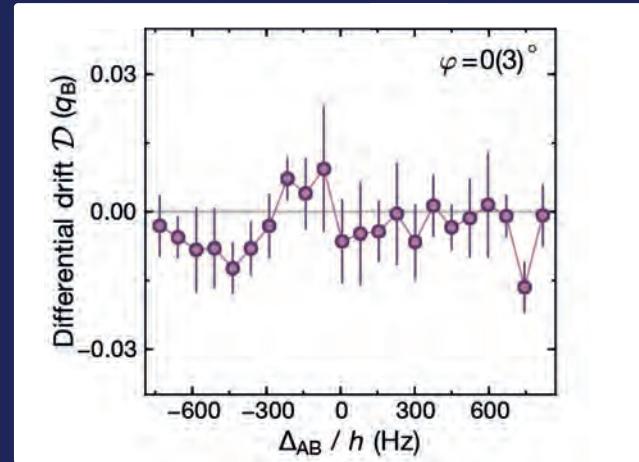
# Berry Curvature - Measurement



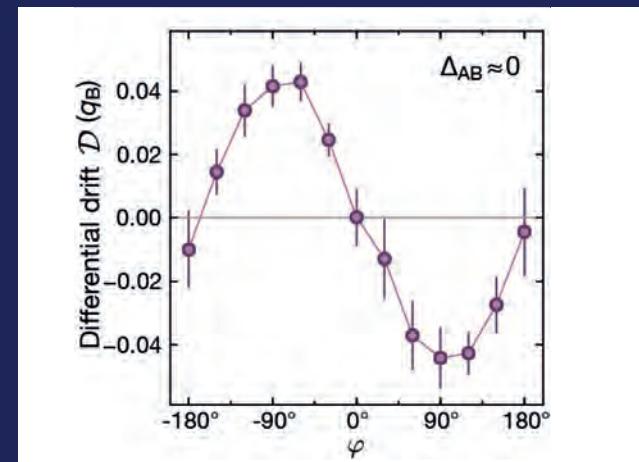
→ Detect difference in center of mass position after full Bloch cycle

# Topological features of the system

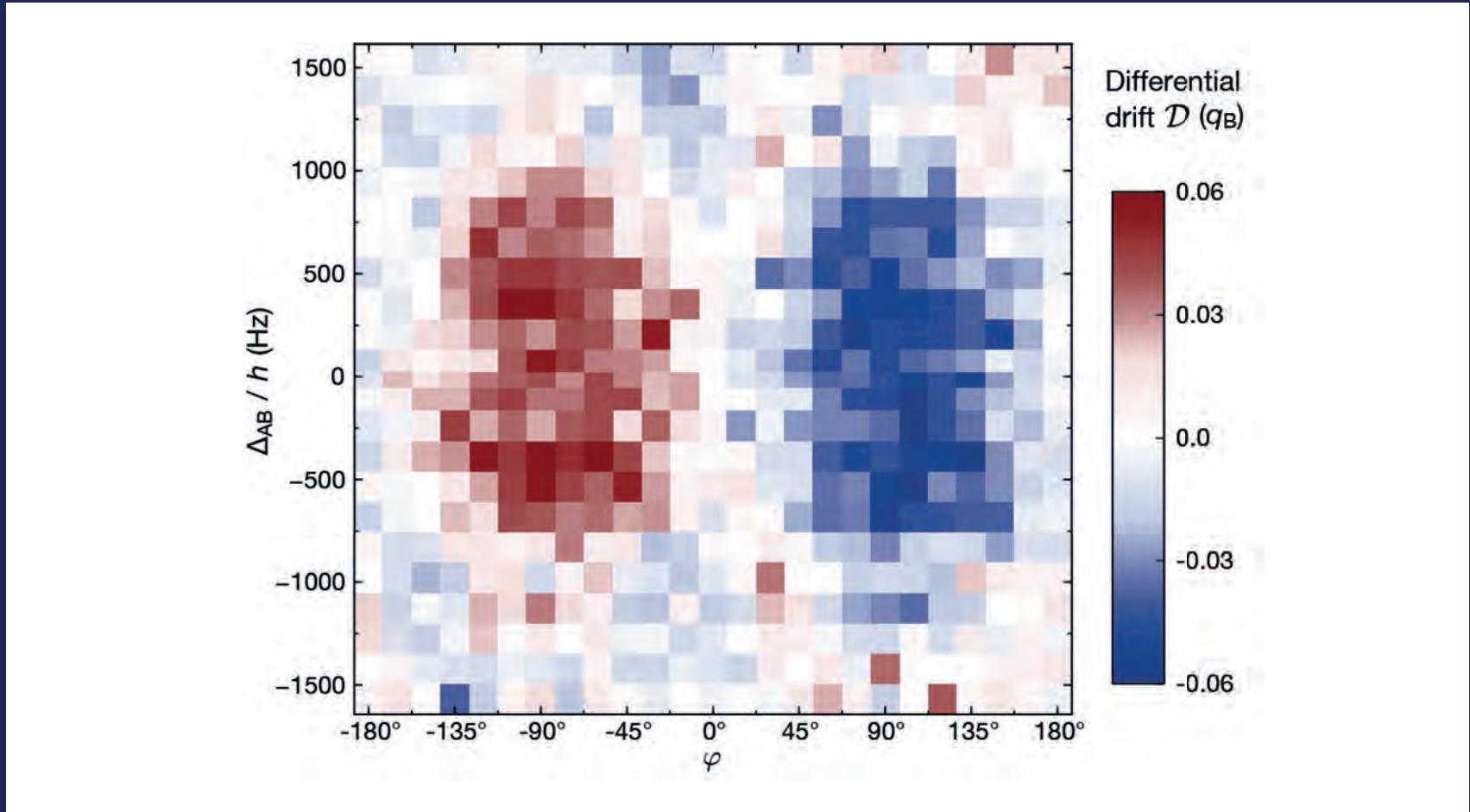
## topologically trivial



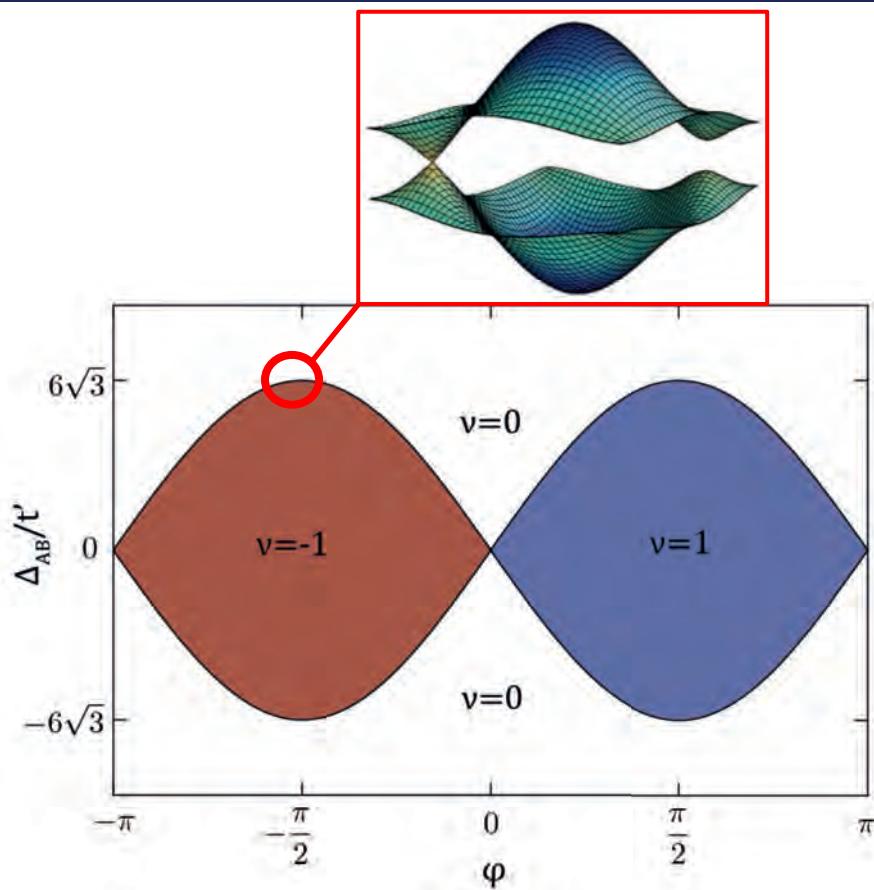
nonzero Chern number



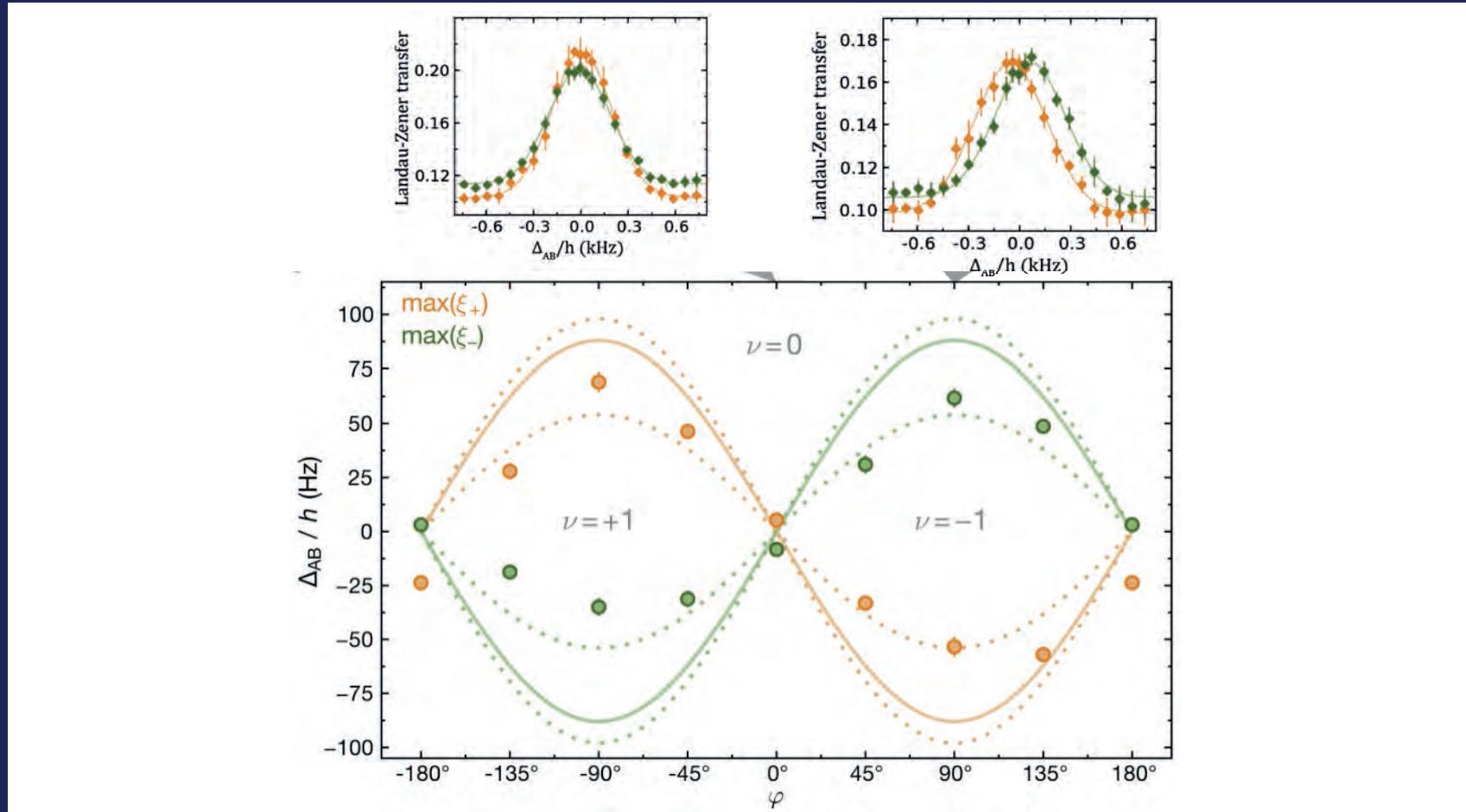
# Observing Transverse Drifts



# Mapping out the transition line

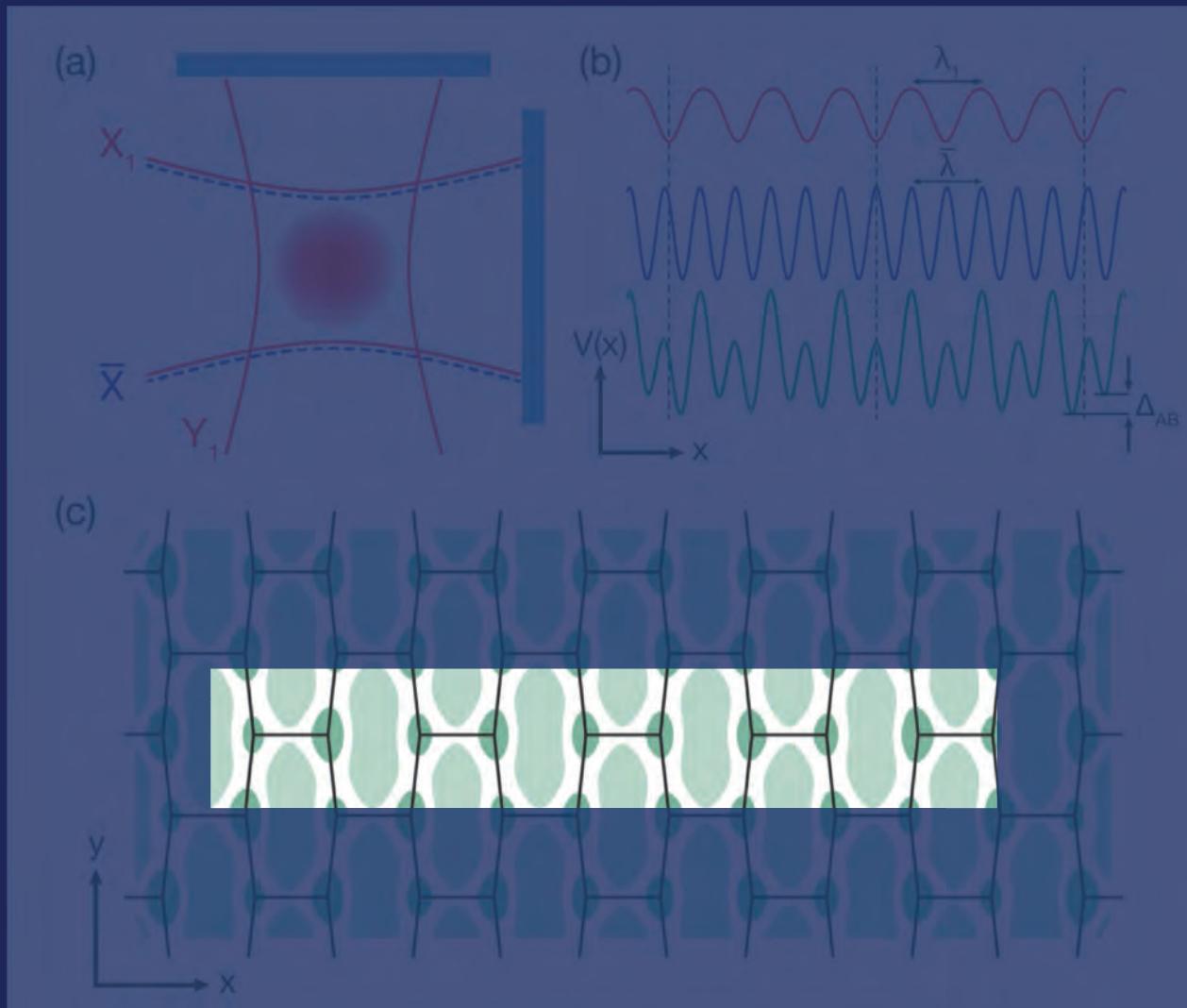


# Mapping out the transition line

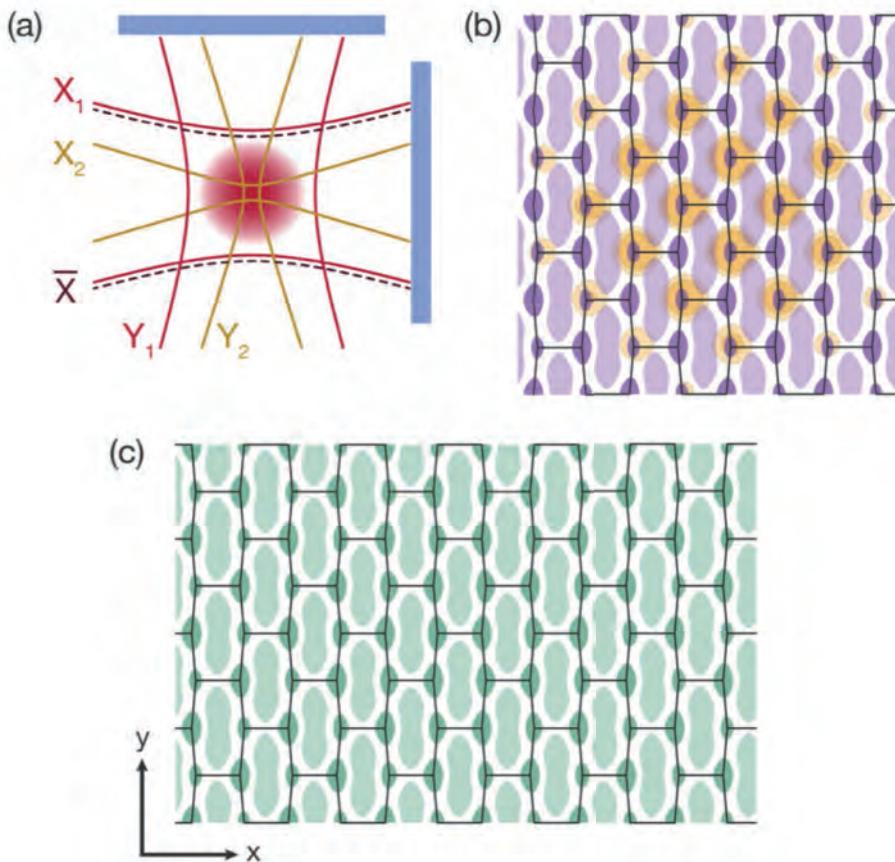


# Edge states

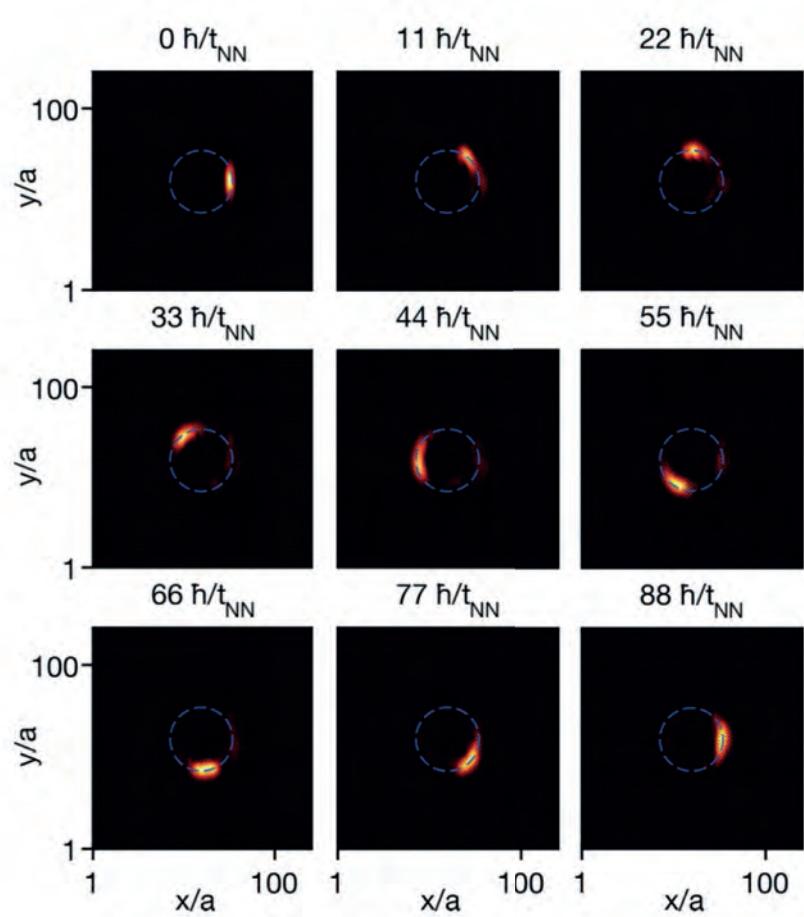
# Edge states



# Edge states



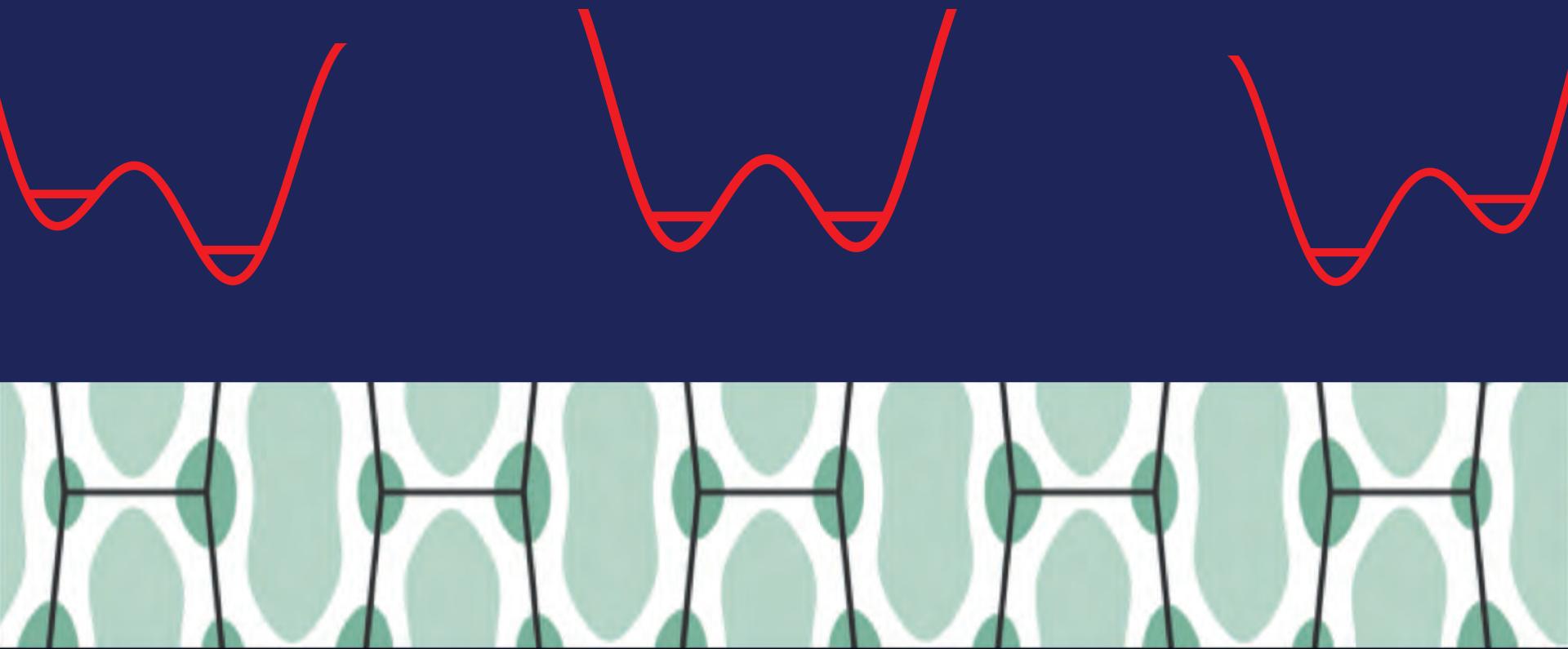
# Edge states

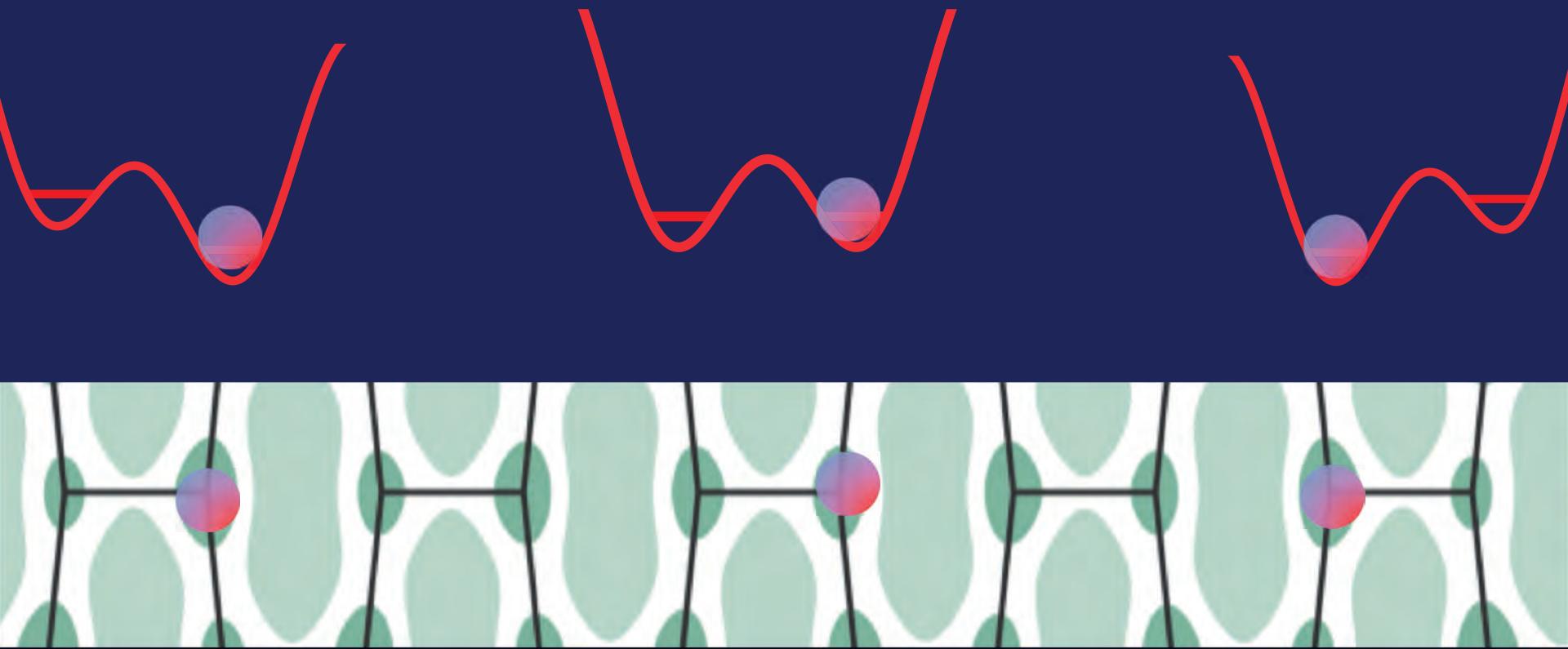


# Interactions?

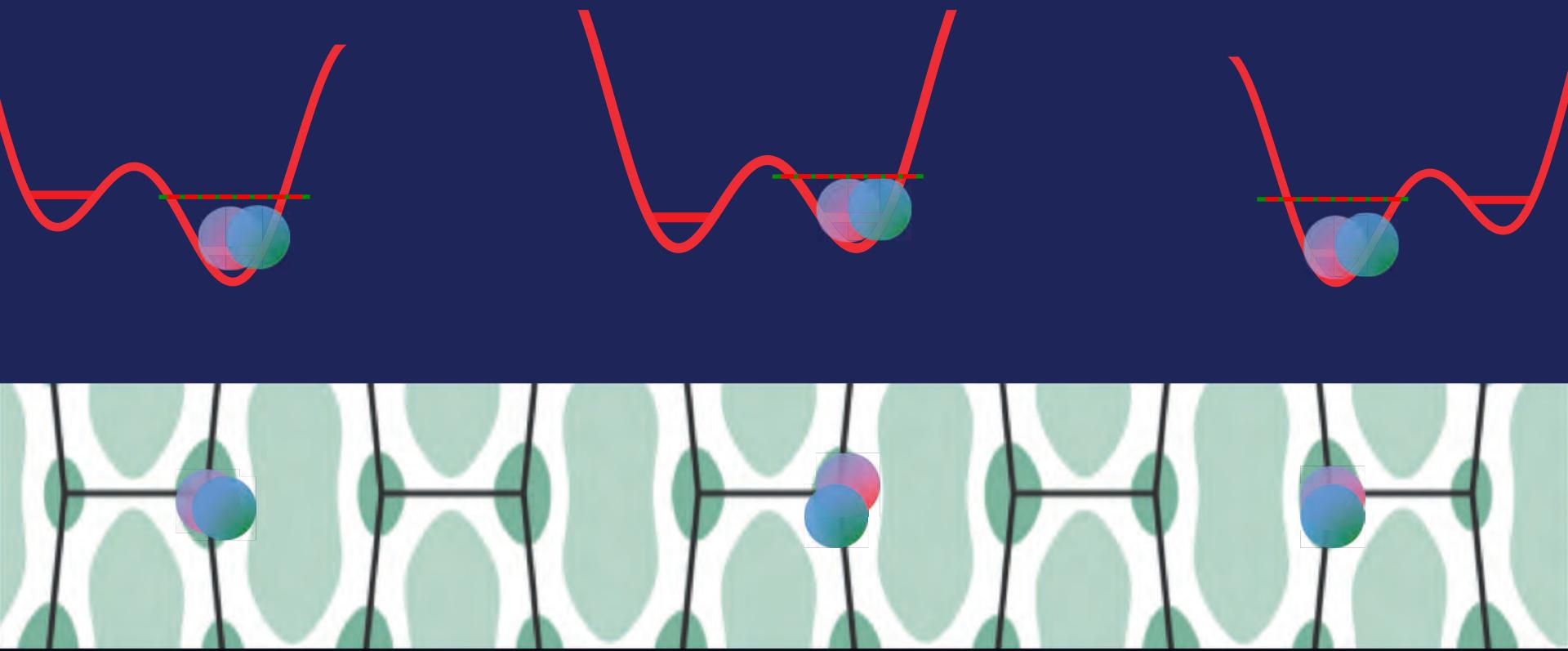
Theoretical challenge

Experimental challenge

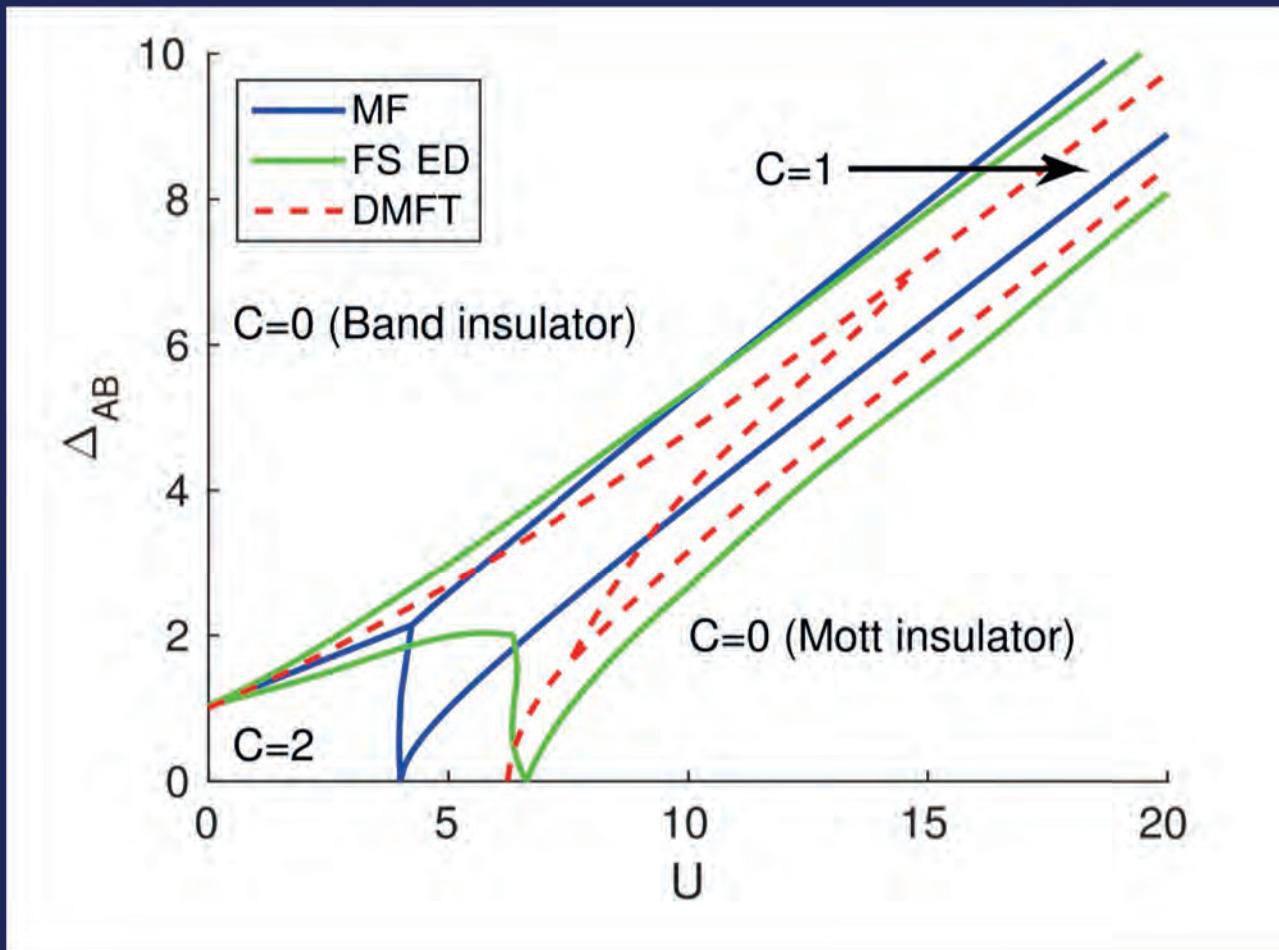




# Interactions

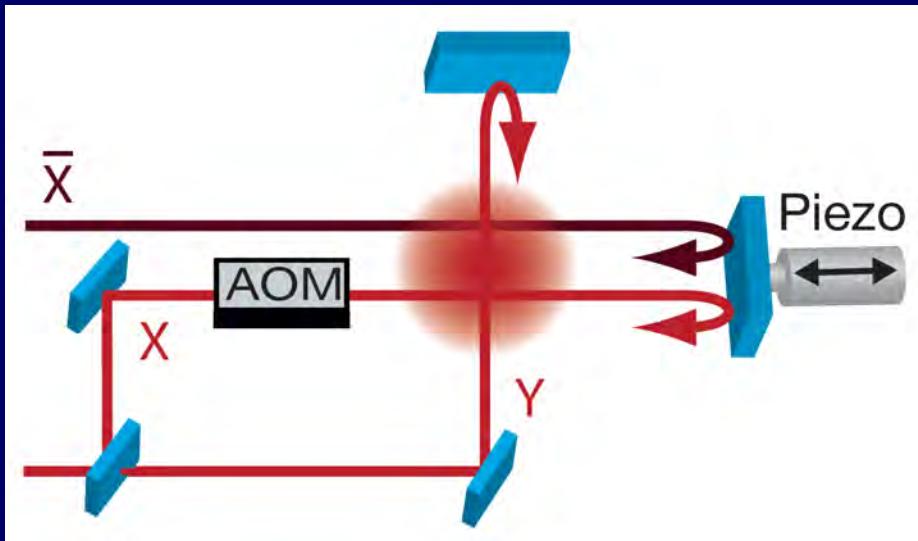


# Interactions

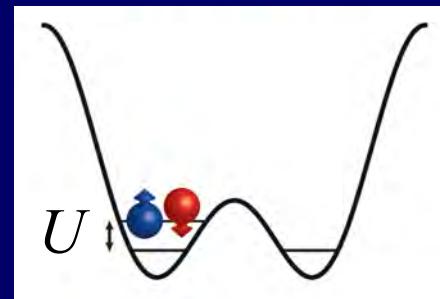
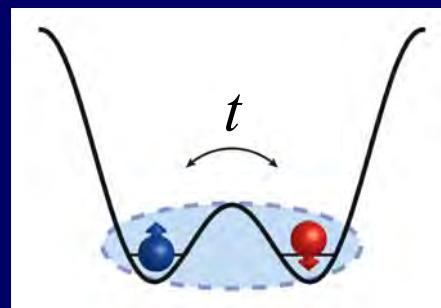
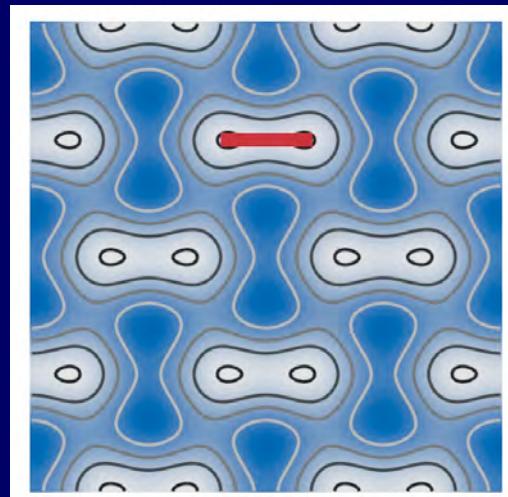


Topological Phase Transitions in the Repulsively Interacting Haldane-Hubbard Model  
T. I. Vanhala, T. Siro, L. Liang, M. Troyer, A. Harju, and Päivi Törmä, PRL 116, 225305 (2016)

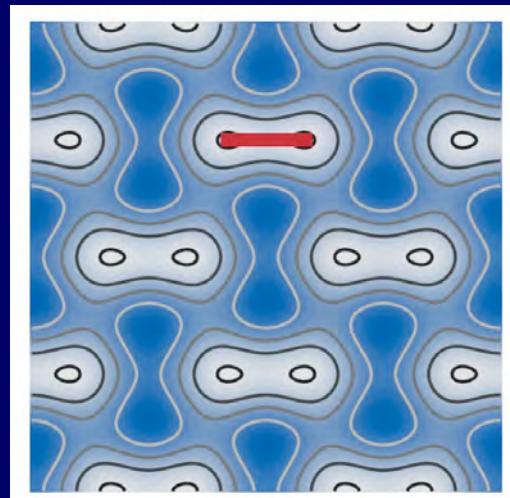
# Shaking and interactions



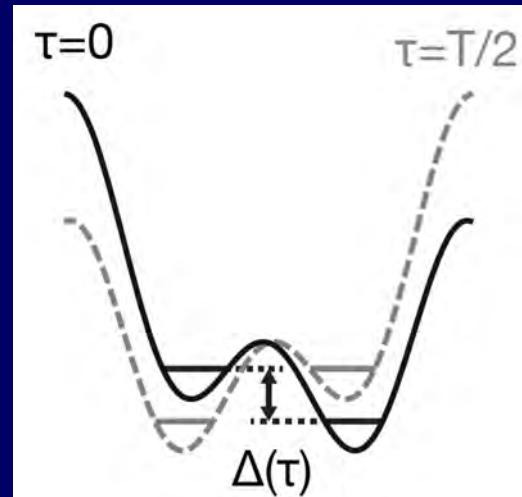
# Double well



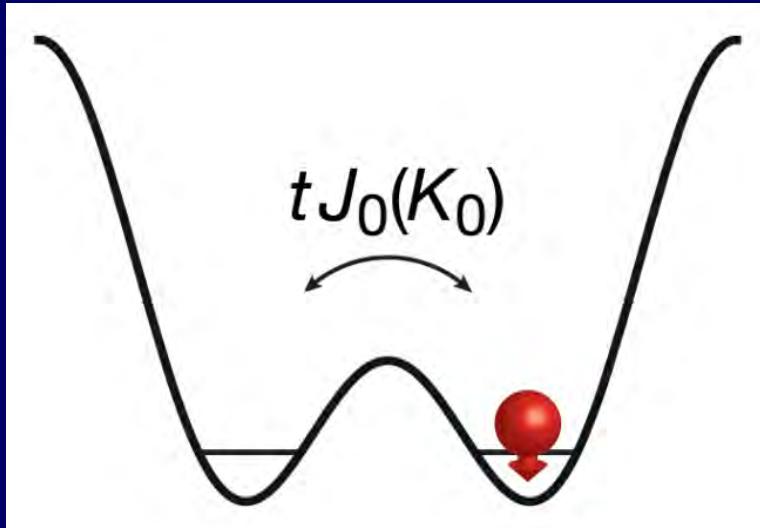
# Driven double well



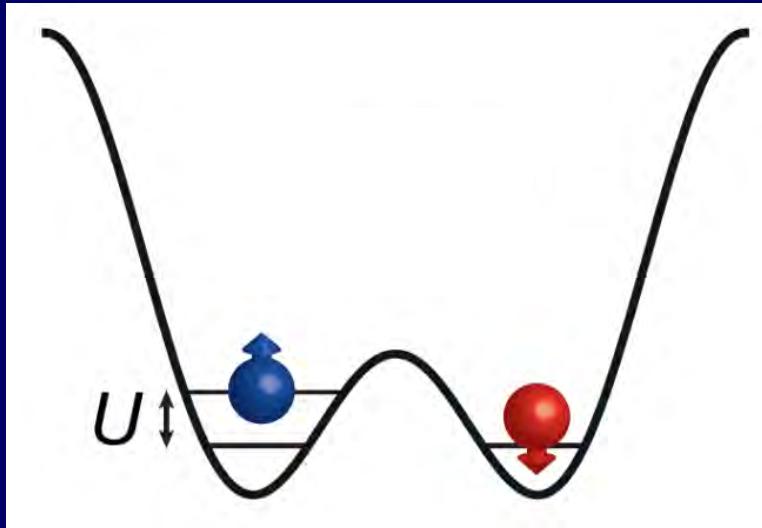
$$t \rightarrow tJ_0(K_0)$$



# Driven double well

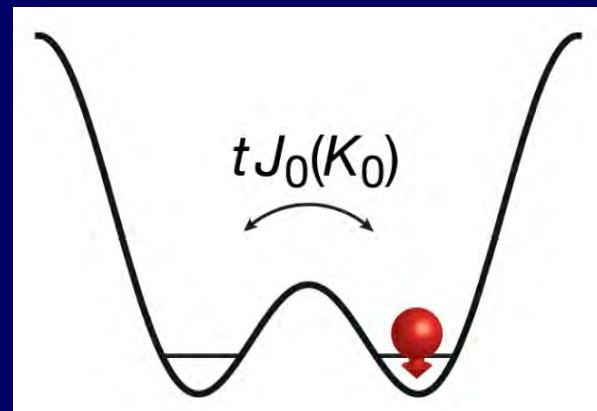
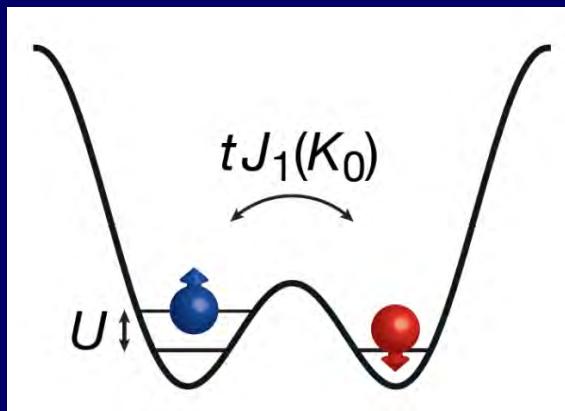


# Driven double well



$$U \approx \omega$$

# Driven double well

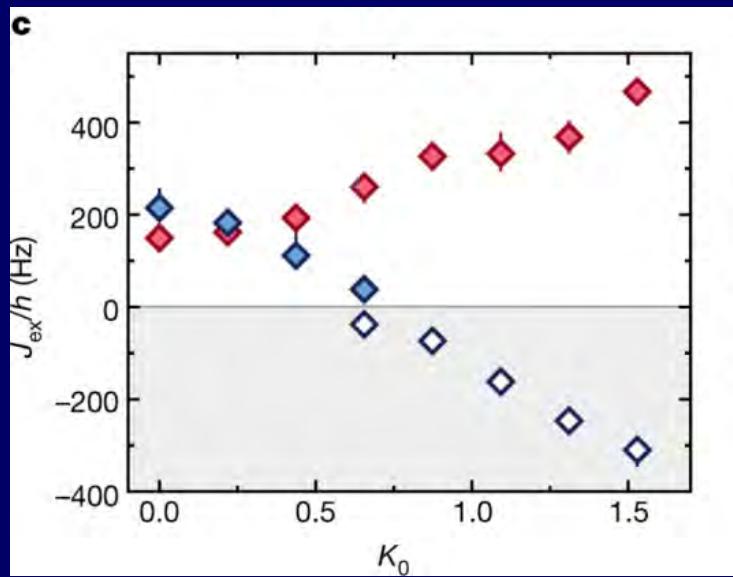
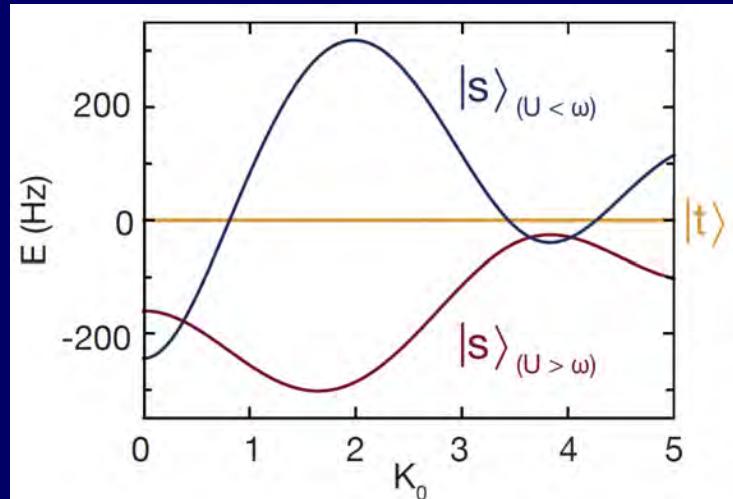


**Exchange**

$$\hbar\omega \gg (\hbar\omega - U) > t$$

$$J_{\text{ex}} = -\frac{4t^2 J_0^2(K_0)}{U} - \frac{4t^2 J_1^2(K_0)}{U - \hbar\omega}$$

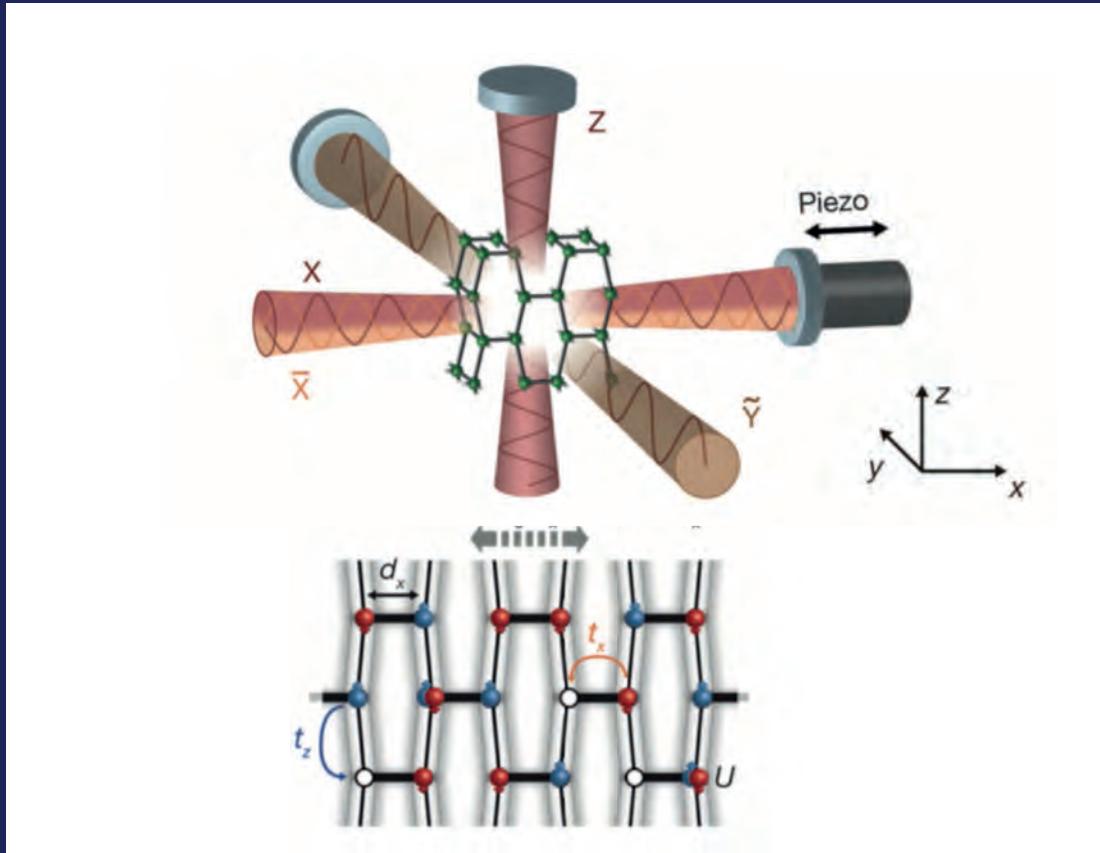
# Driven double well



$U=9.1$  kHz  $> \omega$

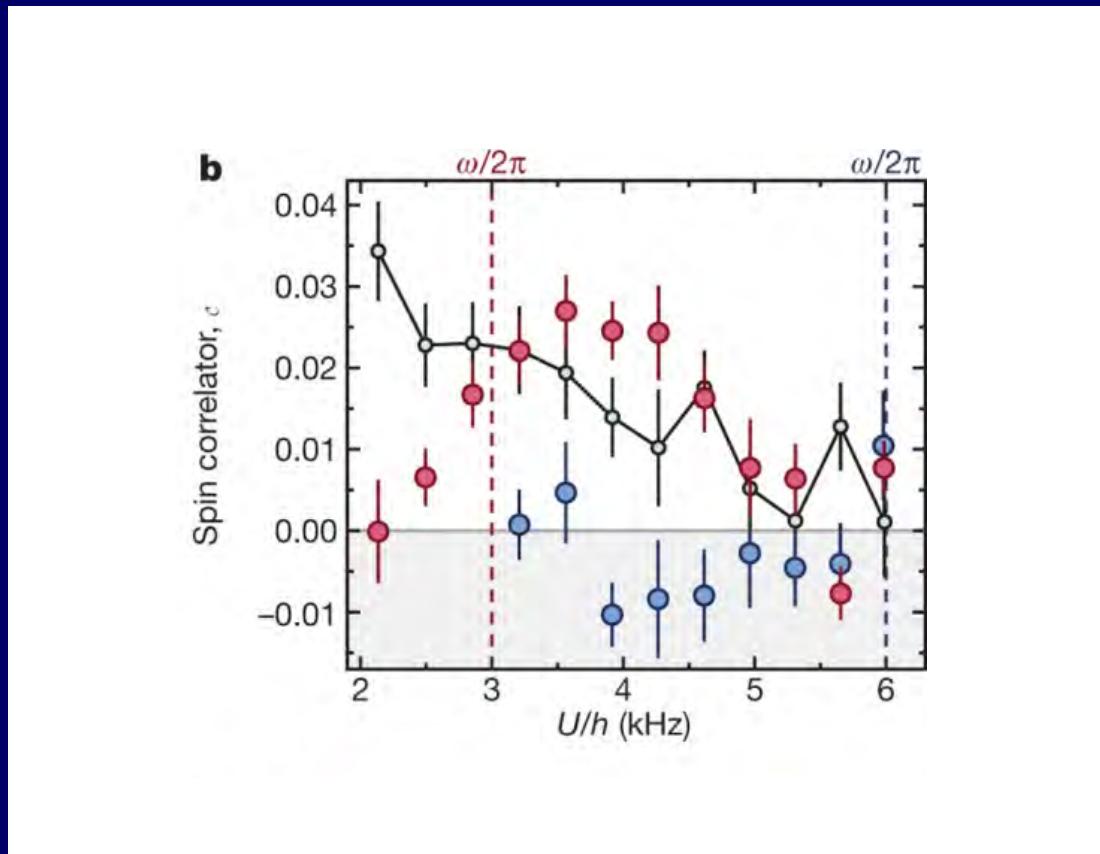
$U=6.5$  kHz  $< \omega$

# Magnetic correlations in driven lattice



F. Görg, M. Messer, K. Sandholzer, G. Jotzu, R. Desbuquois, T.E., Nature 553, 481-485 (2018)  
Also: J. Coulthard, S. R. Clark, S. Al-Assam, A. Cavalleri, D. Jaksch, Phys. Rev. B 96, 085104 (2017)

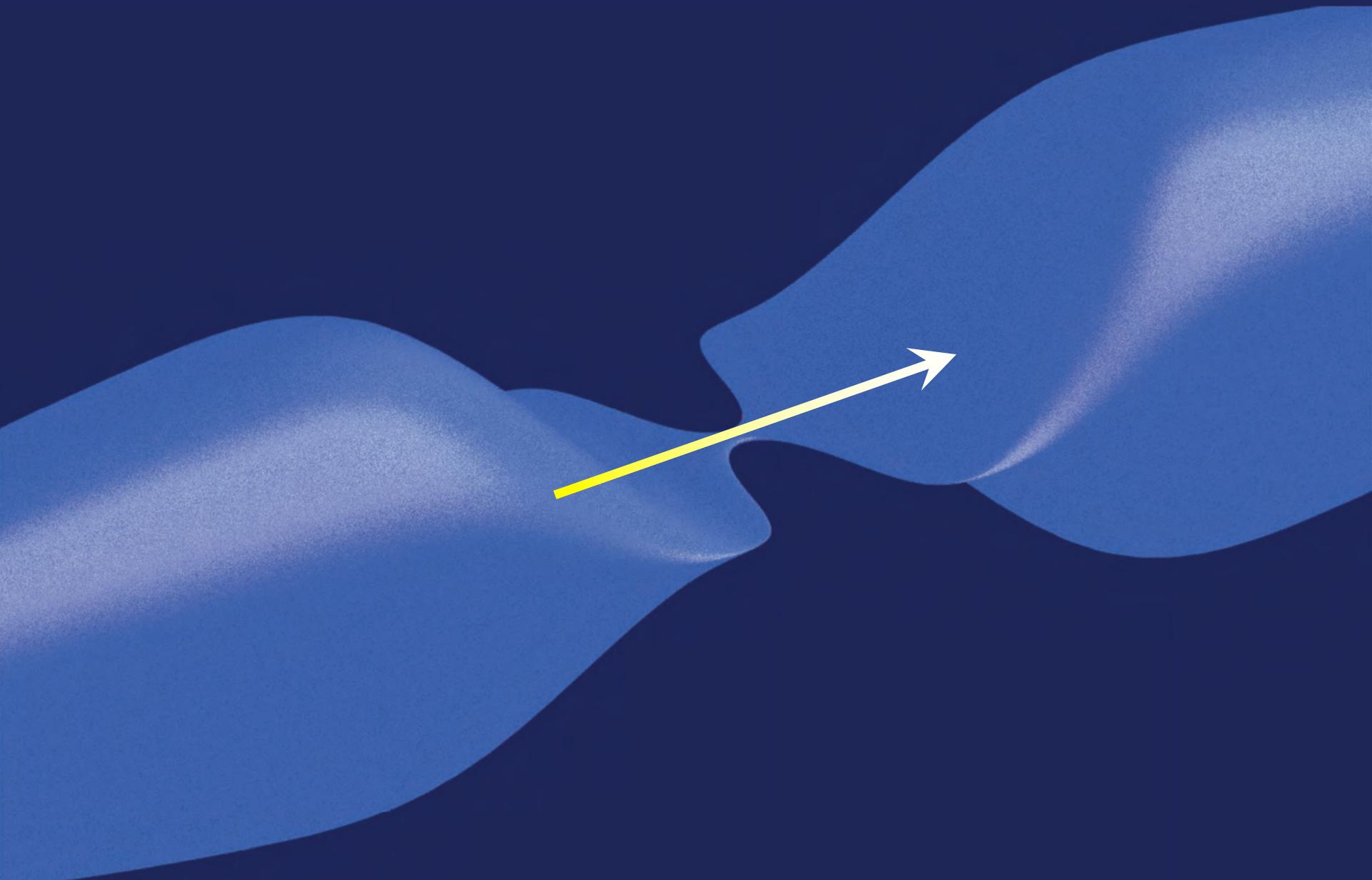
# Enhancement of magnetic correlations in driven lattice



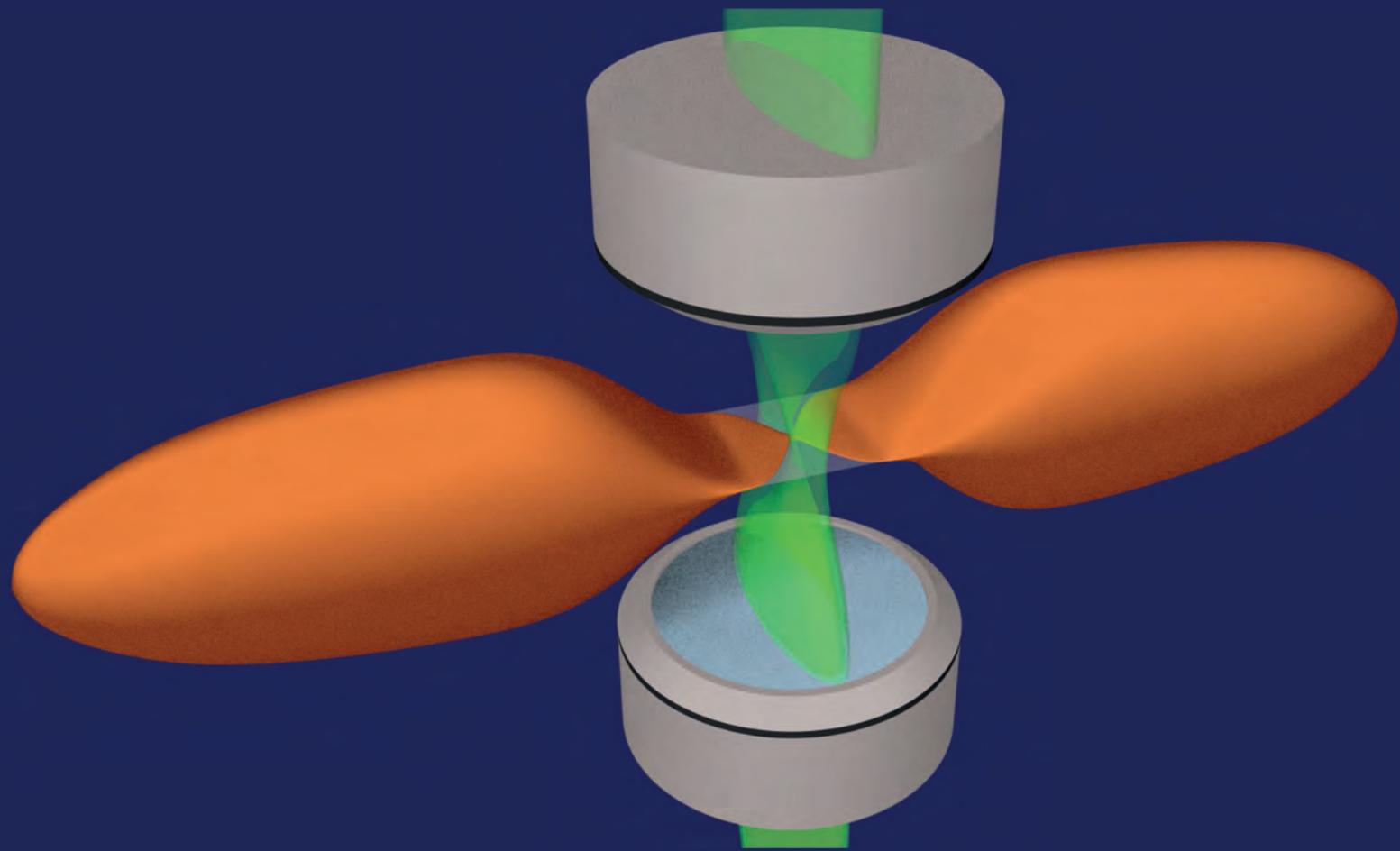
F. Görg, M. Messer, K. Sandholzer, G. Jotzu, R. Desbuquois, T.E., Nature 553, 481-485 (2018)  
Also: J. Coulthard, S. R. Clark, S. Al-Assam, A. Cavalleri, D. Jaksch, Phys. Rev. B 96, 085104 (2017)



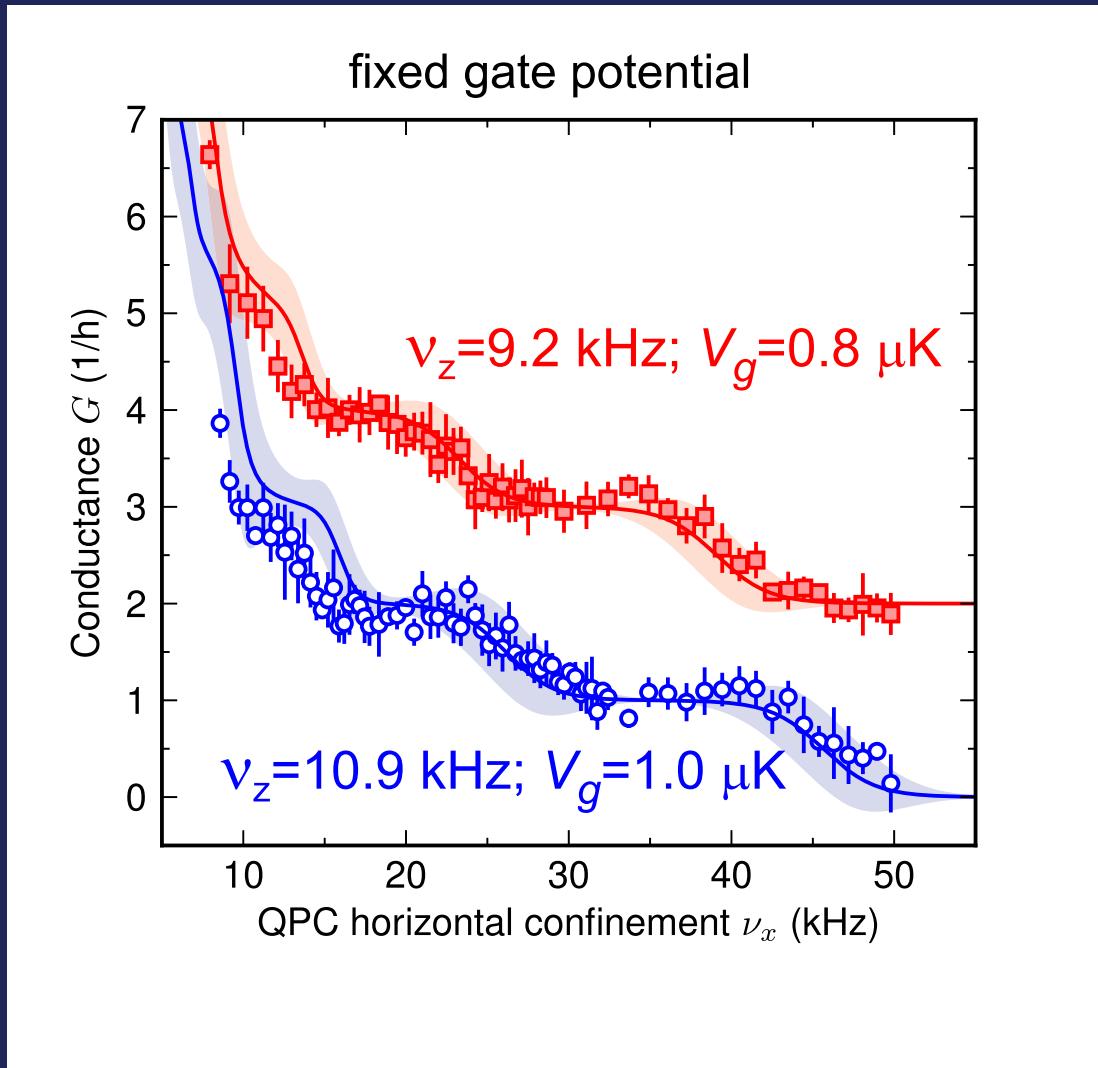
# Quantum Point Contact for Fermions



# Quantum Point Contact for Fermions



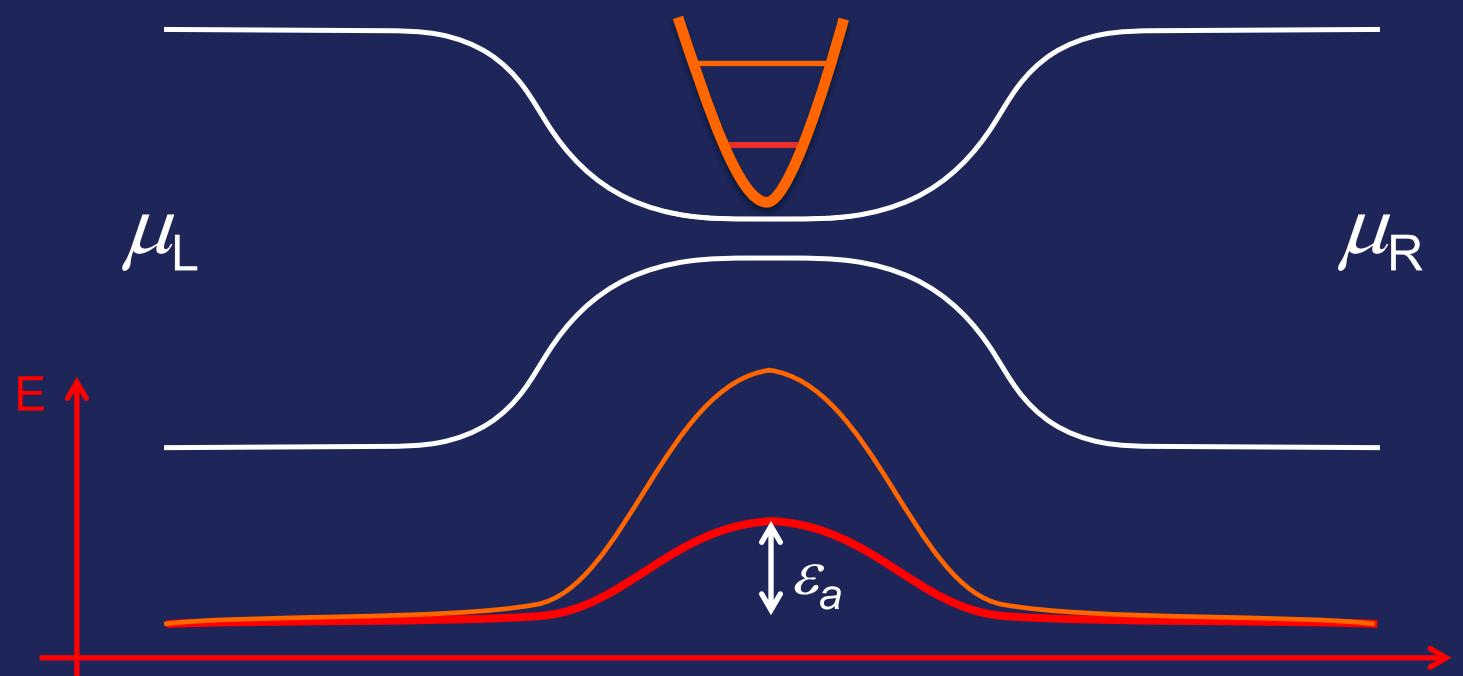
# Quantum Point Contact

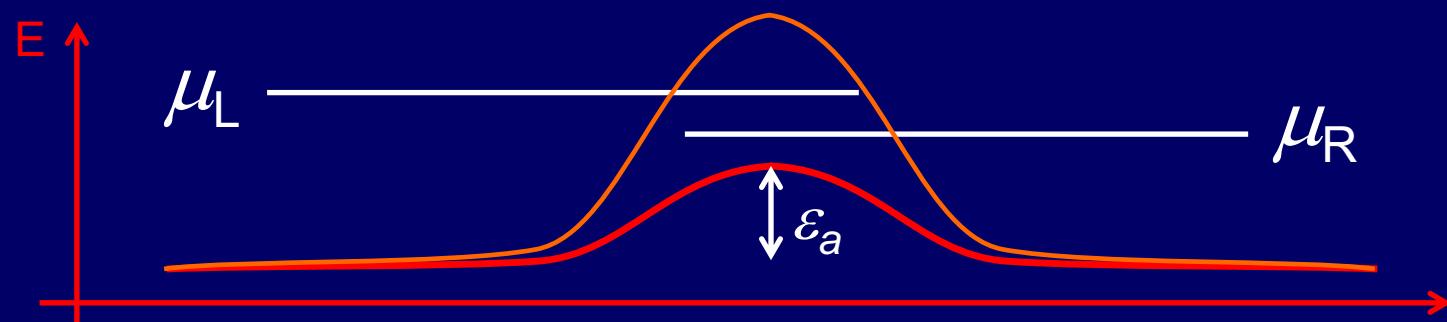


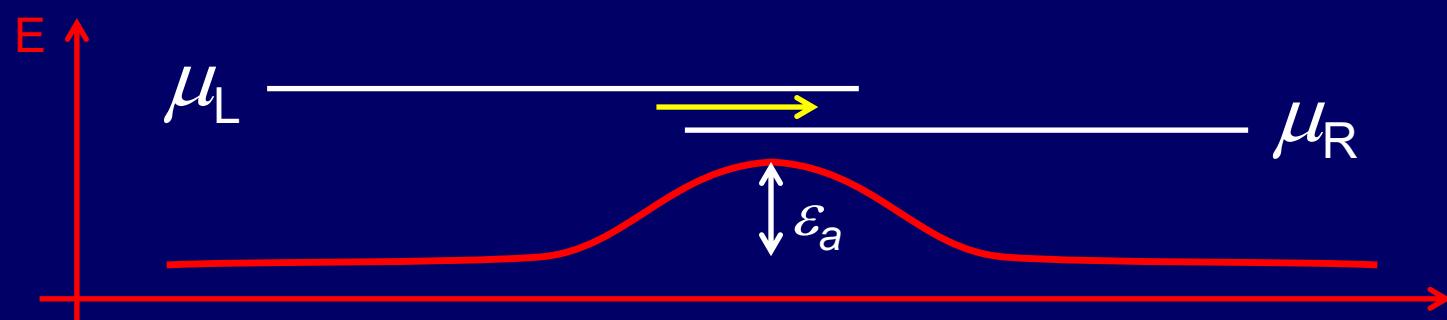
# Left Reservoir

# Constriction

# Right Reservoir

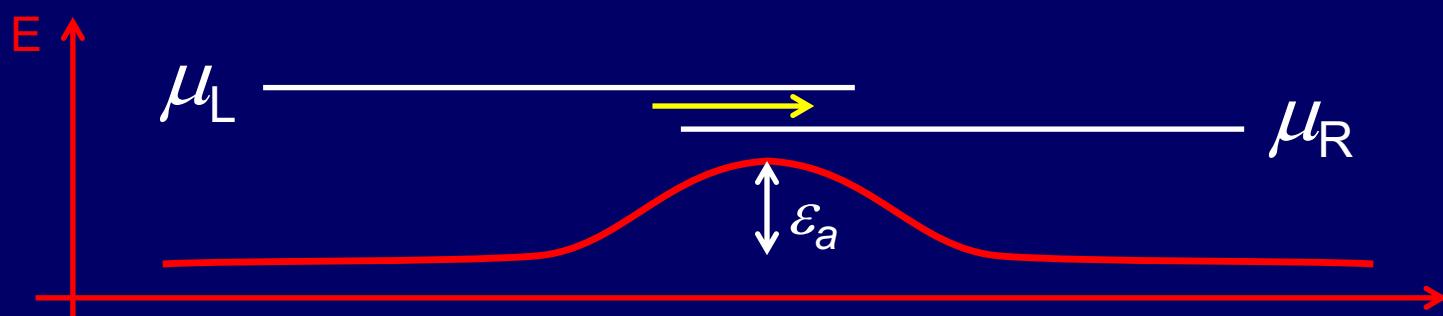






Current for  $T=0$ ,  $T_a=1$ :

$$I_a = \int_{\varepsilon_F}^{\varepsilon_F + \Delta} d\varepsilon g_a(\varepsilon) V_a(\varepsilon) T_a(\varepsilon)$$



Current for  $T=0$ ,  $T_a=1$ :

$$I_a = \int_{\varepsilon_F}^{\varepsilon_F + \Delta} d\varepsilon \ g_a(\varepsilon) \ V_a(\varepsilon) \ T_a(\varepsilon)$$

# Current for $T=0$ , $T_a=1$ :

$$I_a = \int_{\varepsilon_F}^{\varepsilon_F + \Delta} d\varepsilon g_a(\varepsilon) v_a(\varepsilon) T_a(\varepsilon) = \frac{\Delta}{h}$$

velocity:

$$v_a(\varepsilon) = \frac{\hbar k_a}{m} = \sqrt{2(\varepsilon - \varepsilon_a)/m}$$

density of states:  
(right movers)

$$g_a(\varepsilon) = \frac{1}{2\pi} \frac{dk_a}{d\varepsilon} = \frac{1}{2\pi \hbar v_a(\varepsilon)}$$

Current for  $T=0$ ,  $T_a=1$ :

$$I_a = \int_{\varepsilon_F}^{\varepsilon_F + \Delta} d\varepsilon g_a(\varepsilon) V_a(\varepsilon) T_a(\varepsilon) = \frac{\Delta}{h}$$

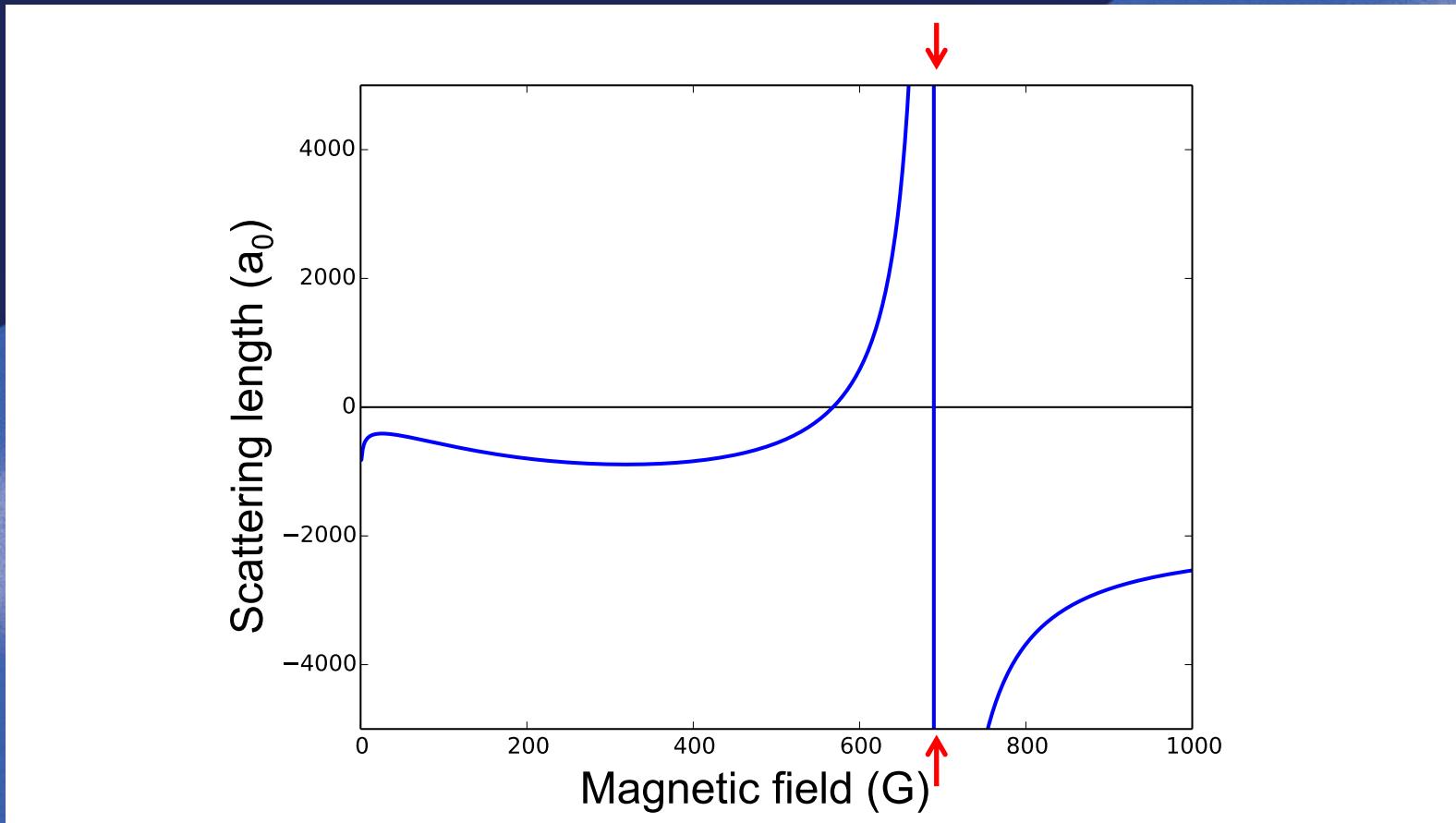
Conductance  $G = \frac{1}{h}$

Consequence of Heisenberg + Pauli's principle

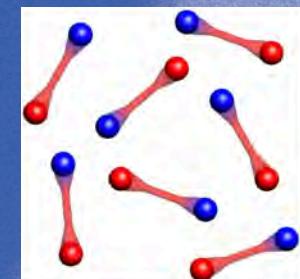
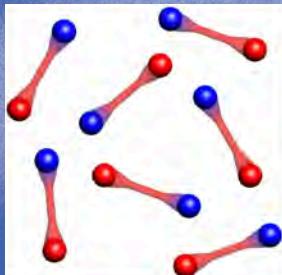
Landauer, Büttiker, Imry

# Connecting two strongly correlated superfluids with quantum point contact

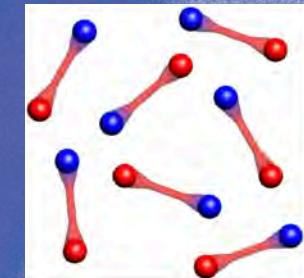
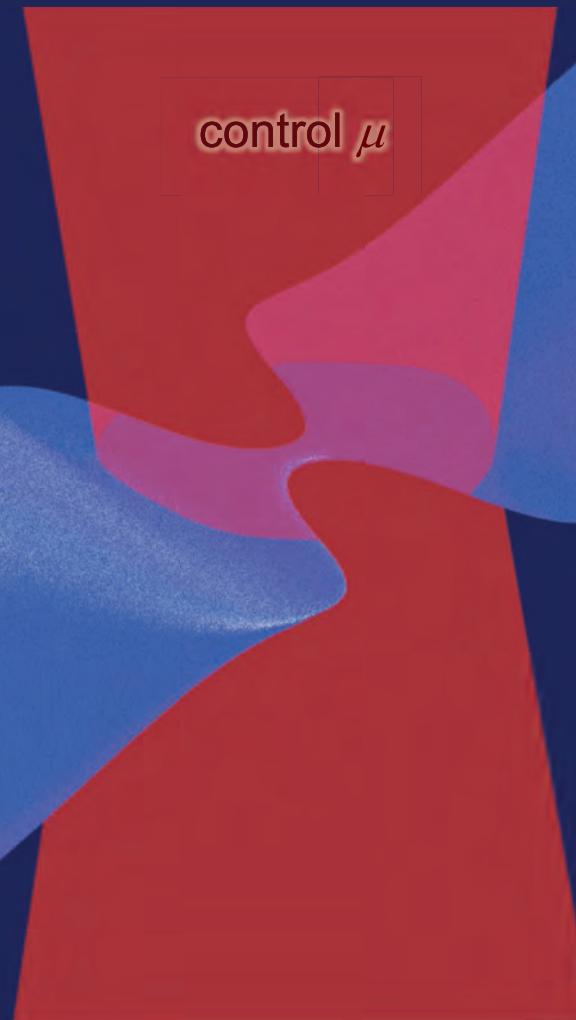
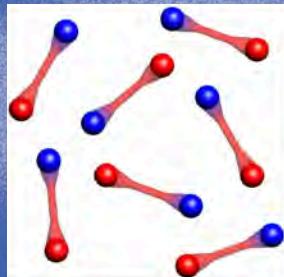
# Connecting two strongly correlated superfluids with quantum point contact



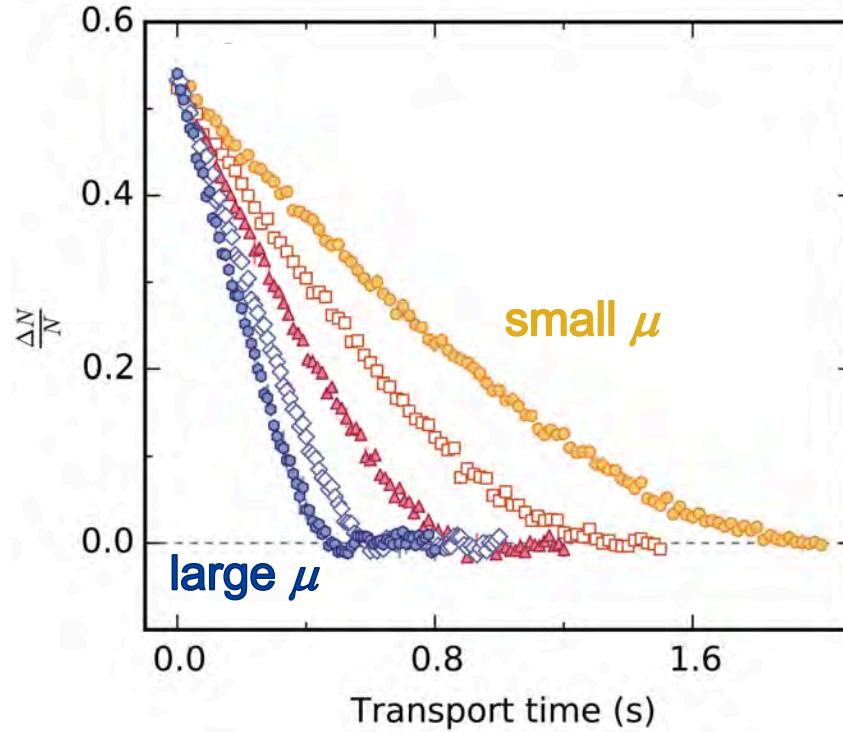
# Connecting two strongly correlated superfluids with quantum point contact



# Connecting two strongly correlated superfluids with quantum point contact



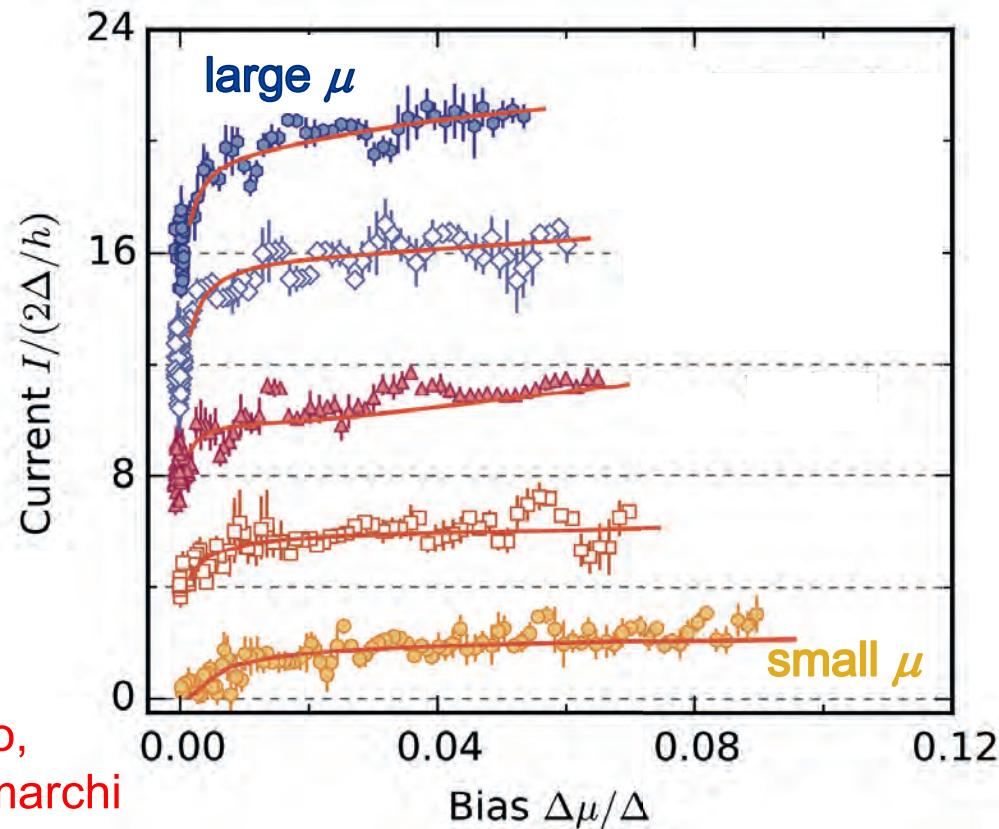
# Connecting two strongly correlated superfluids with quantum point contact



# Connecting two strongly correlated superfluids with quantum point contact

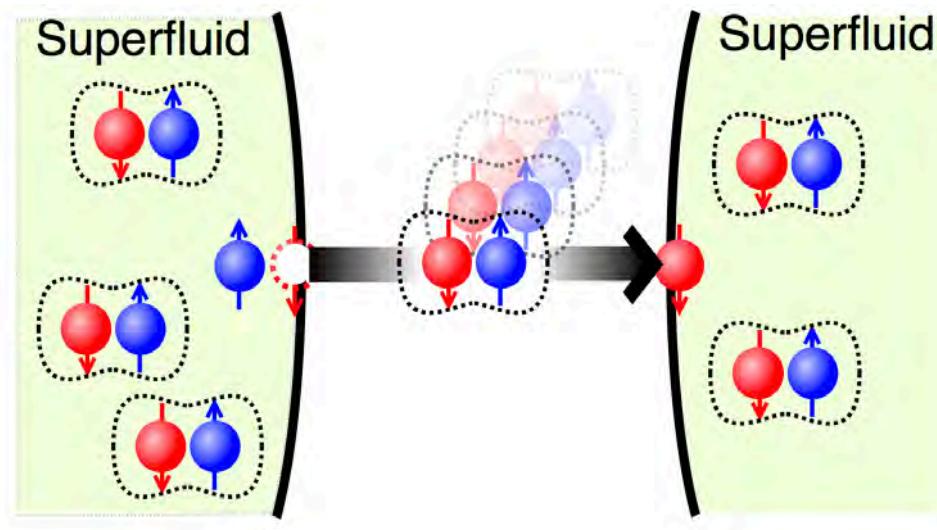


Theory:  
Shun Uchino,  
Thierry Giamarchi



Dominik Husmann, Shun Uchino, Sebastian Krinner, Martin Lebrat, Thierry Giamarchi, Tilman Esslinger, Jean-Philippe Brantut, Science 350, 1498-1501 (2015)

# Connecting two strongly correlated superfluids with quantum point contact



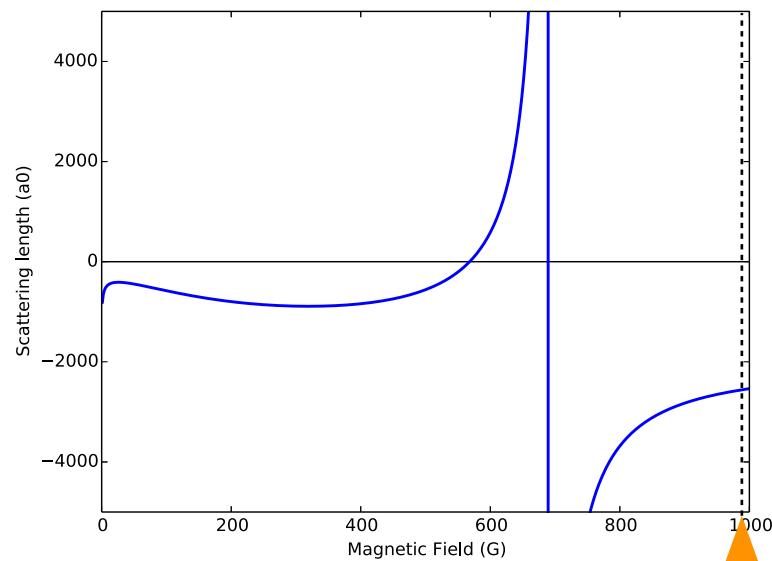
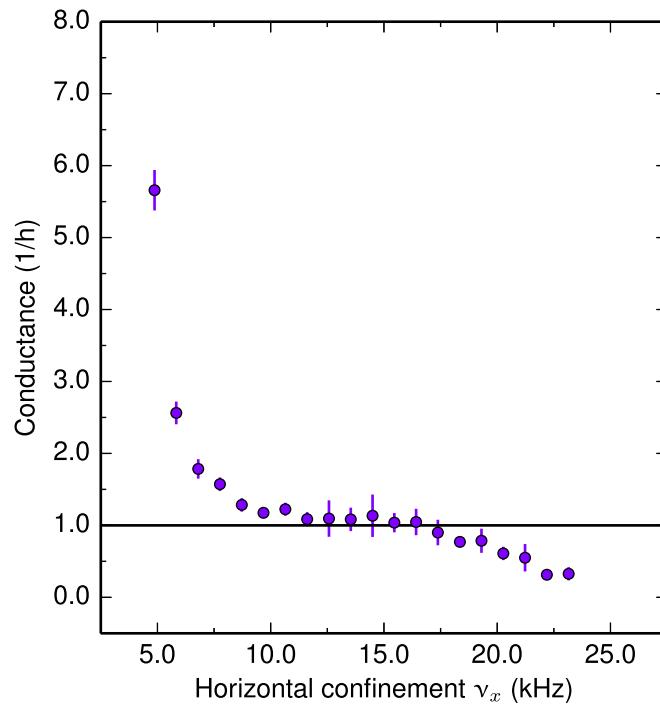
**Theory:**  
Shun Uchino,  
Thierry Giamarchi

Gap for single particle transfer bridged by  
coherent transfer of  $n$  pairs  
(multiple Andreev reflection)

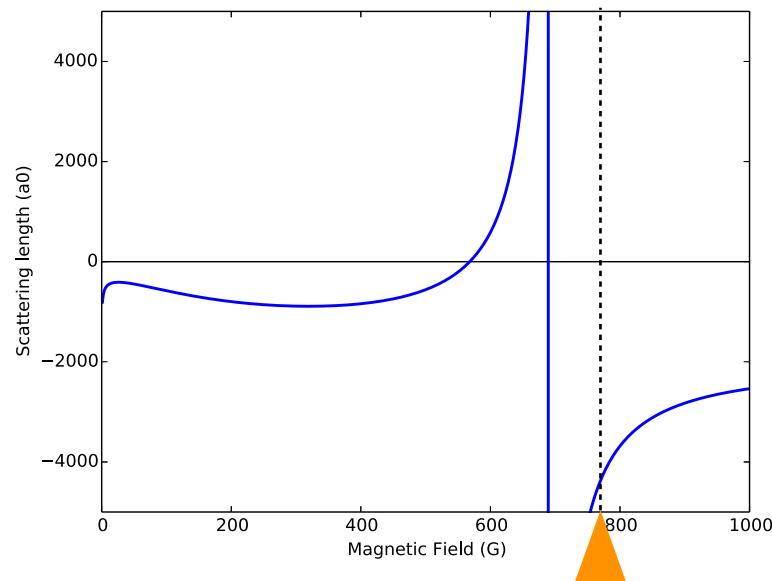
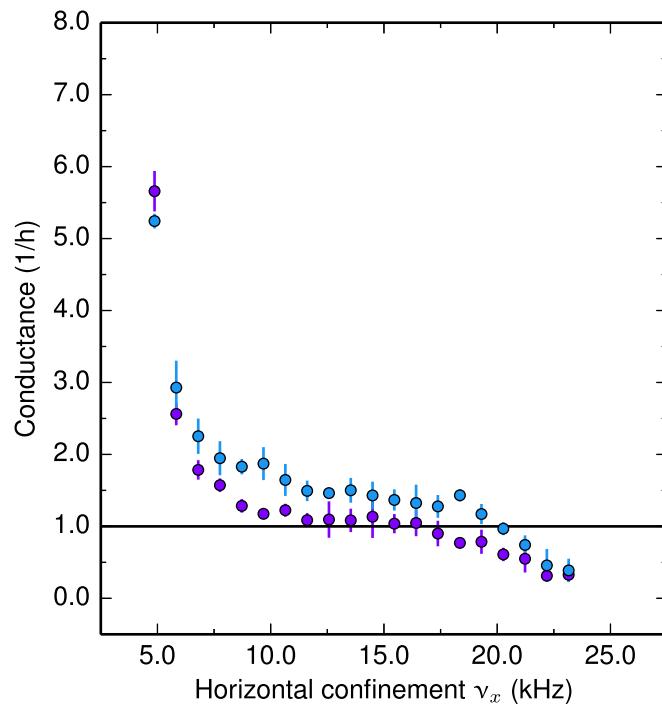
Dominik Husmann, Shun Uchino, Sebastian Krinner, Martin Lebrat, Thierry Giamarchi,  
Tilman Esslinger, Jean-Philippe Brantut, Science 350, 1498-1501 (2015)

Josephson Effect: G. Valtolina, , A. Burchianti, A. Amico, E. Neri, K. Xhani, J. A. Seman, A. Trombettoni, A. Smerzi, M. Zaccanti, M. Inguscio, G. Roati, Science 3501505-1508 (2015).

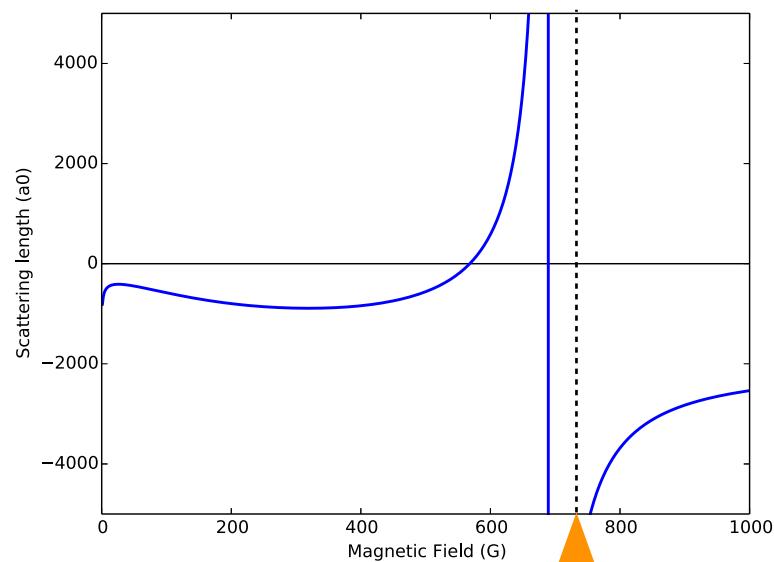
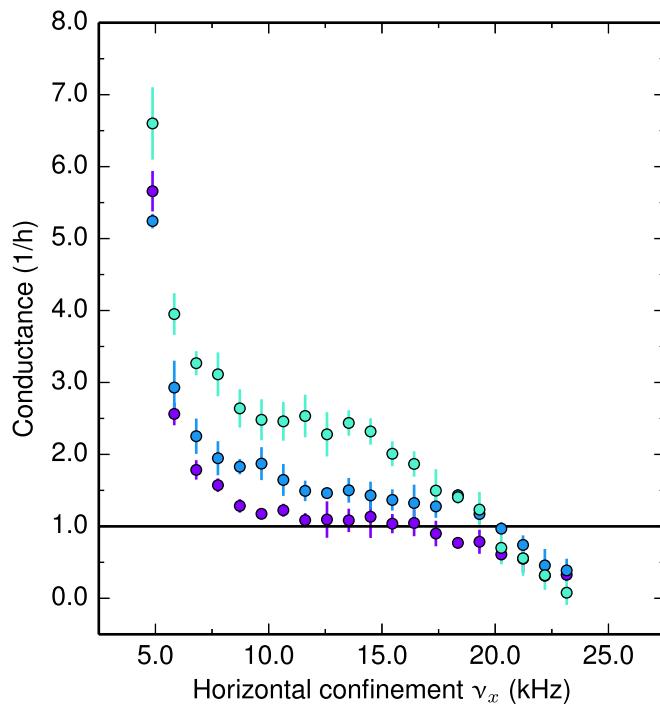
# Changing Interactions



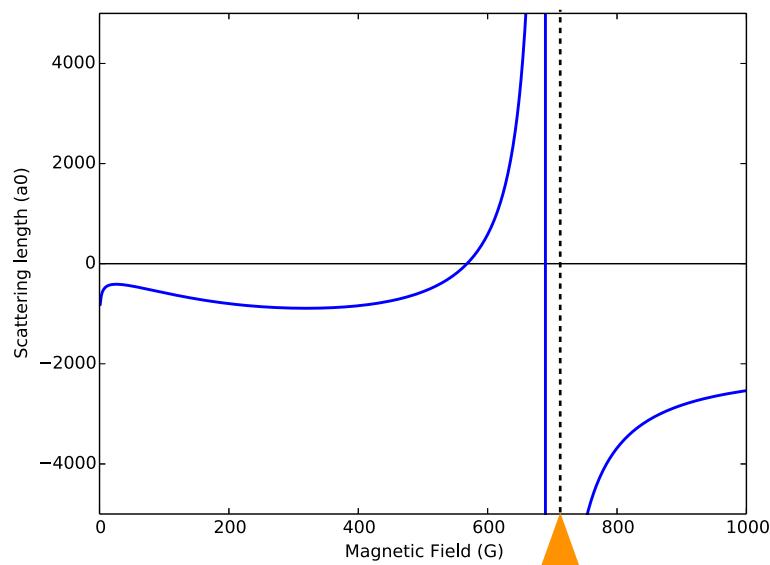
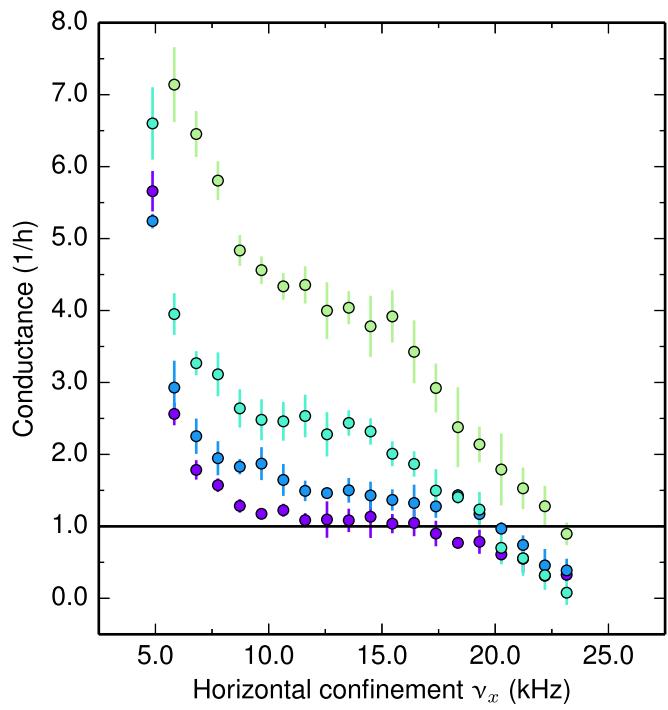
# Changing Interactions



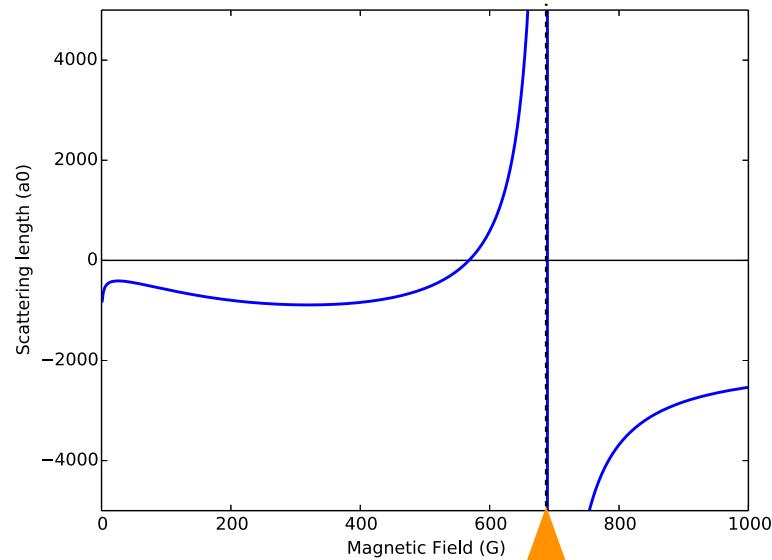
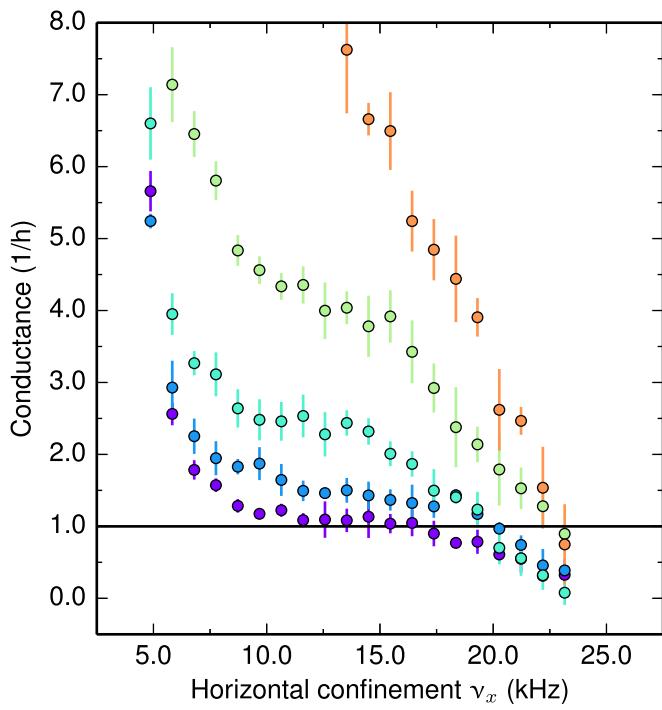
# Changing Interactions



# Changing Interactions



# Changing Interactions



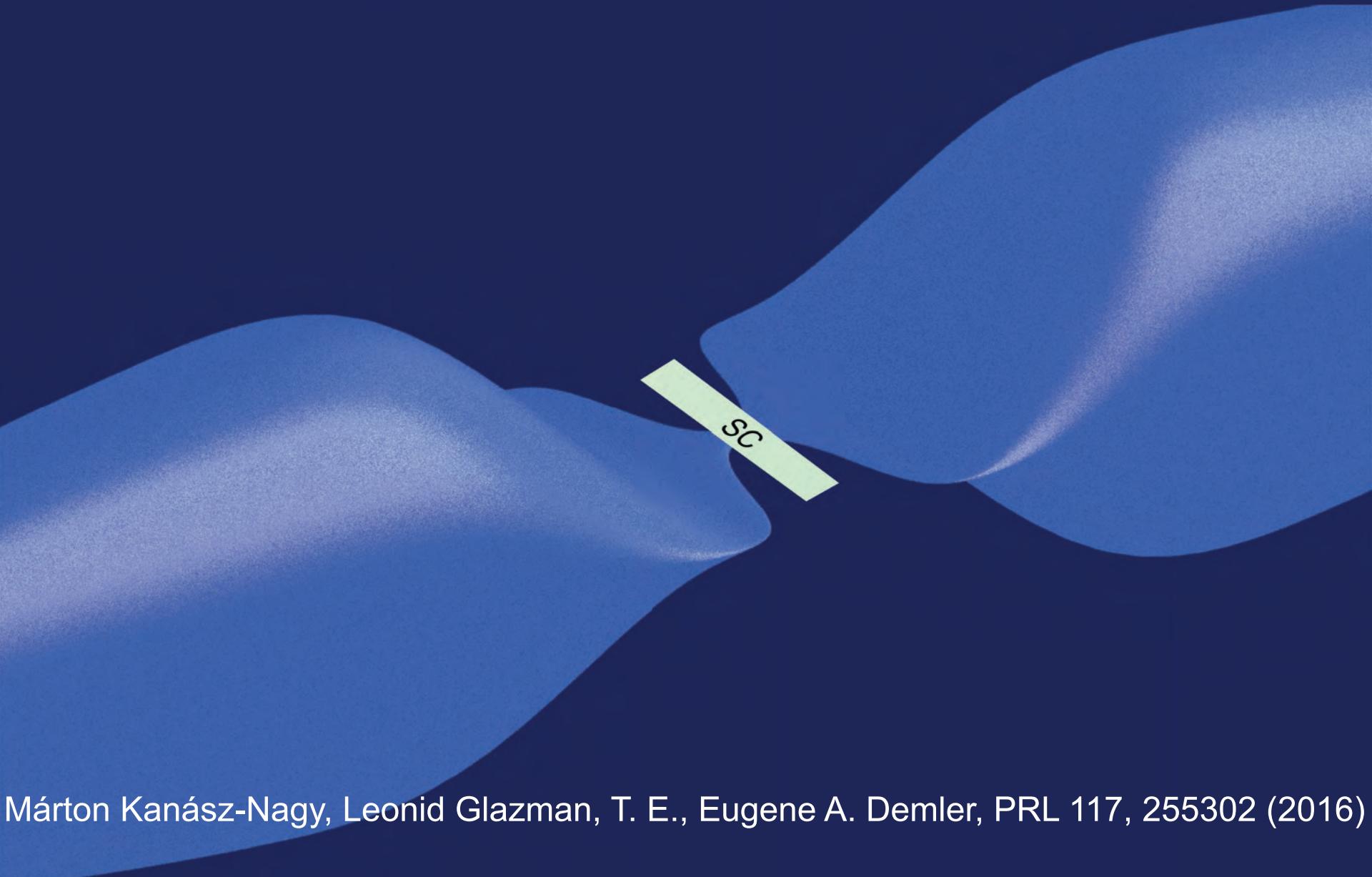
Sebastian Krinner, Martin Lebrat, Dominik Husmann, Charles Grenier, Jean-Philippe Brantut, T.E. PNAS 113, 8144 (2016).

???

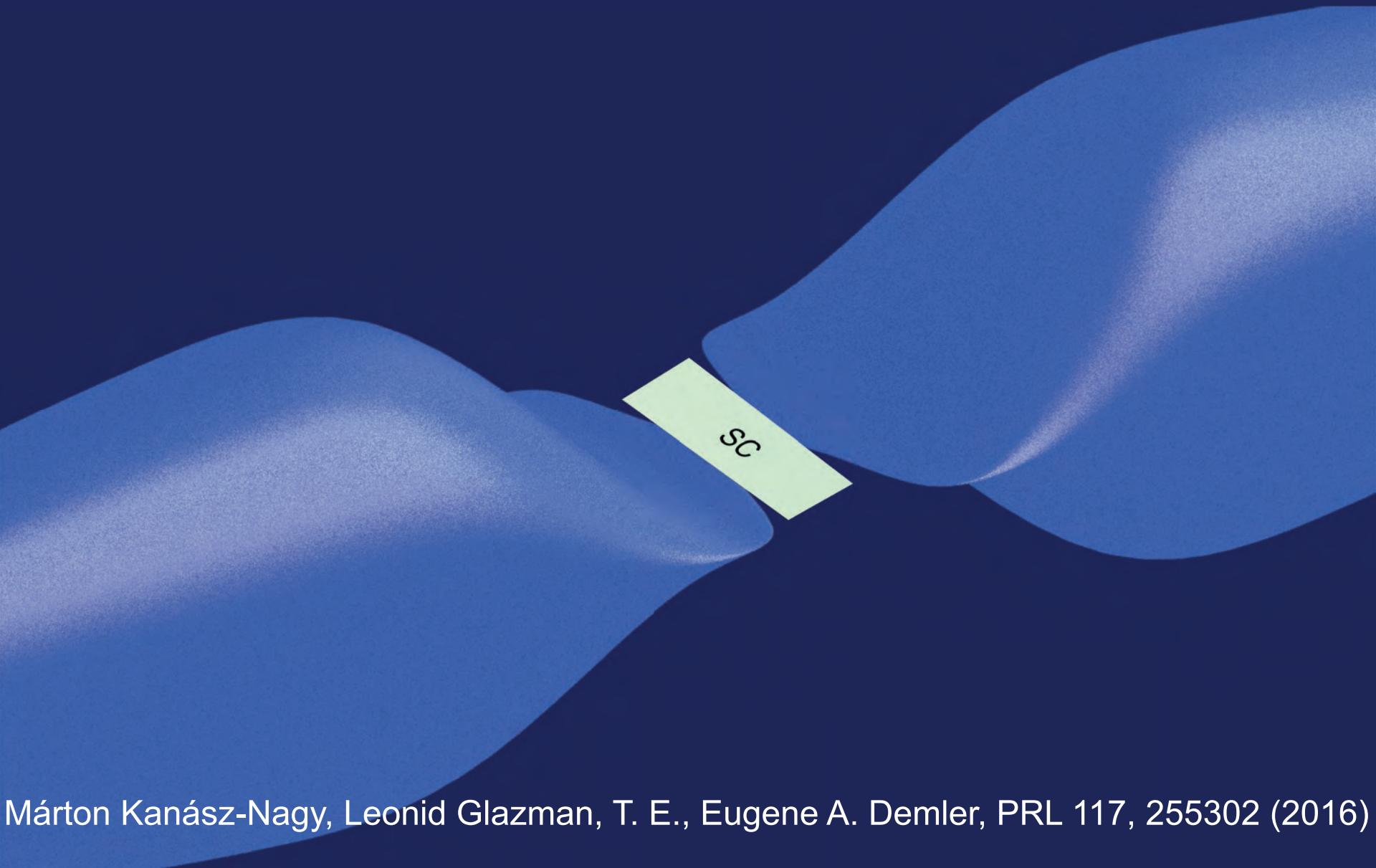
Shun Uchino and Masahito Ueda, Phys. Rev. Lett. **118**, 105303 (2017).

Boyang Liu, Hui Zhai, and Shizhong Zhang, Phys. Rev. A **95**, 013623 (2017).

Márton Kanász-Nagy, Leonid Glazman, T. E., Eugene A. Demler, PRL 117, 255302 (2016)

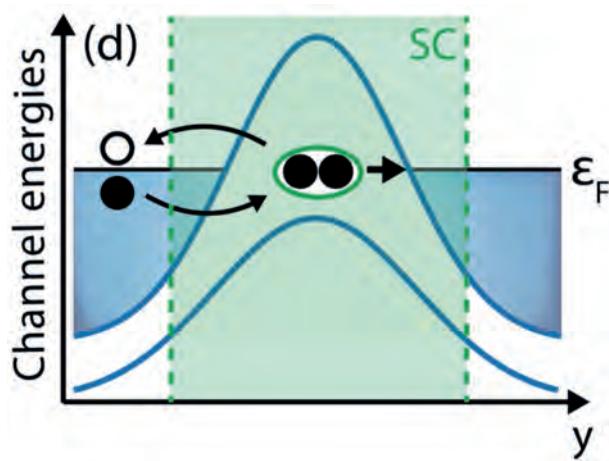


Márton Kanász-Nagy, Leonid Glazman, T. E., Eugene A. Demler, PRL 117, 255302 (2016)



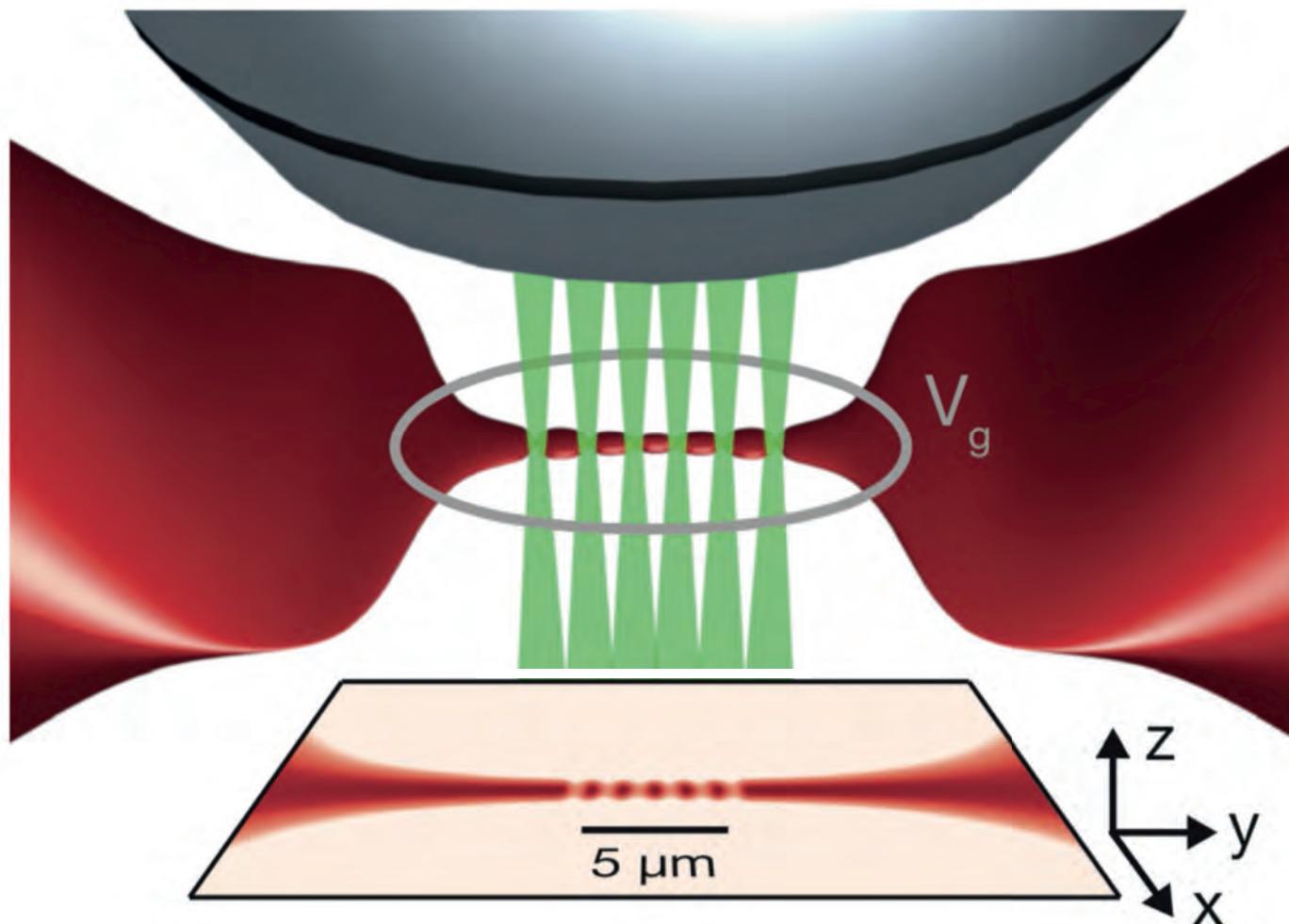
Márton Kanász-Nagy, Leonid Glazman, T. E., Eugene A. Demler, PRL 117, 255302 (2016)

???

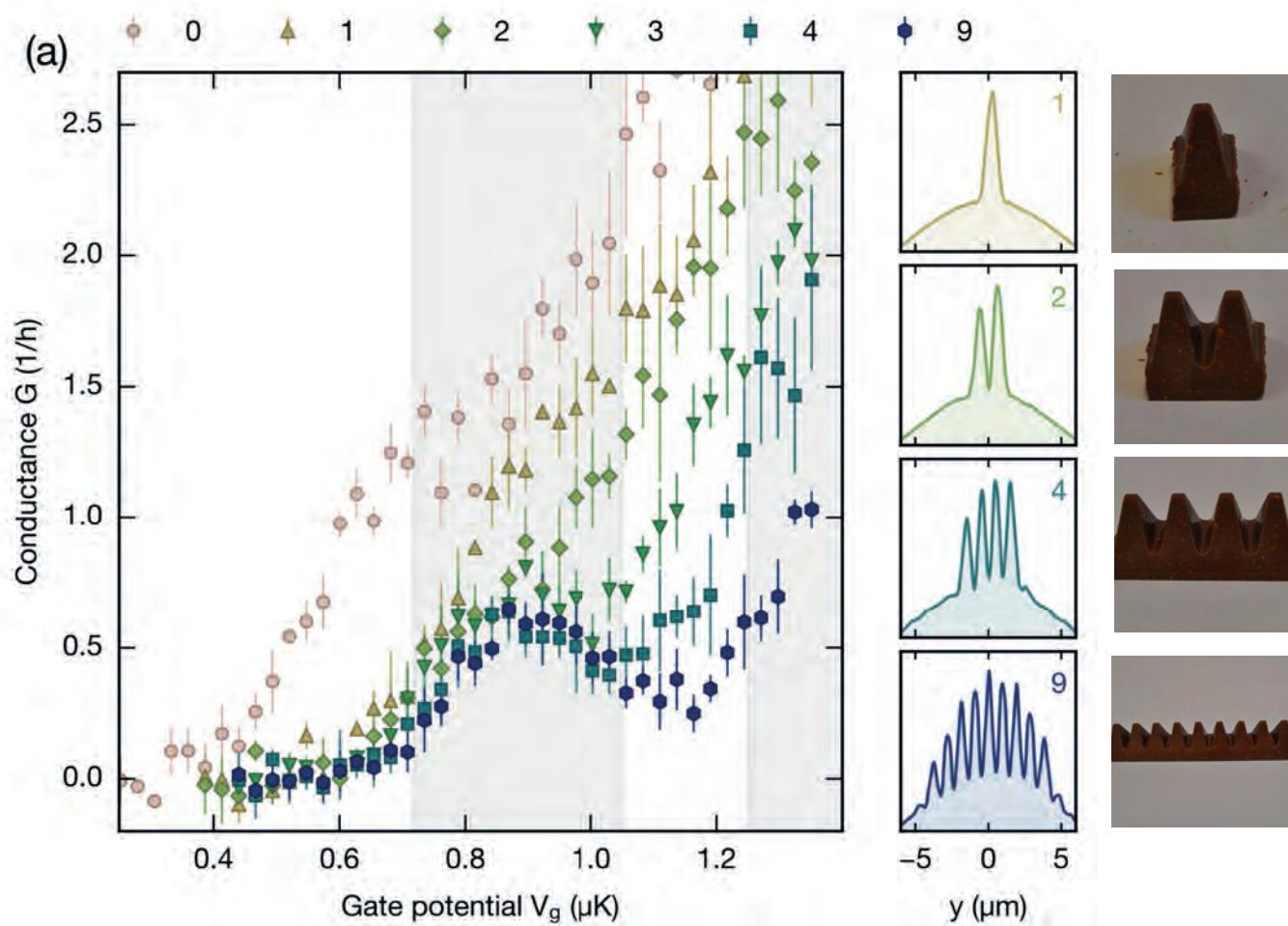


Enhanced pairing in channel

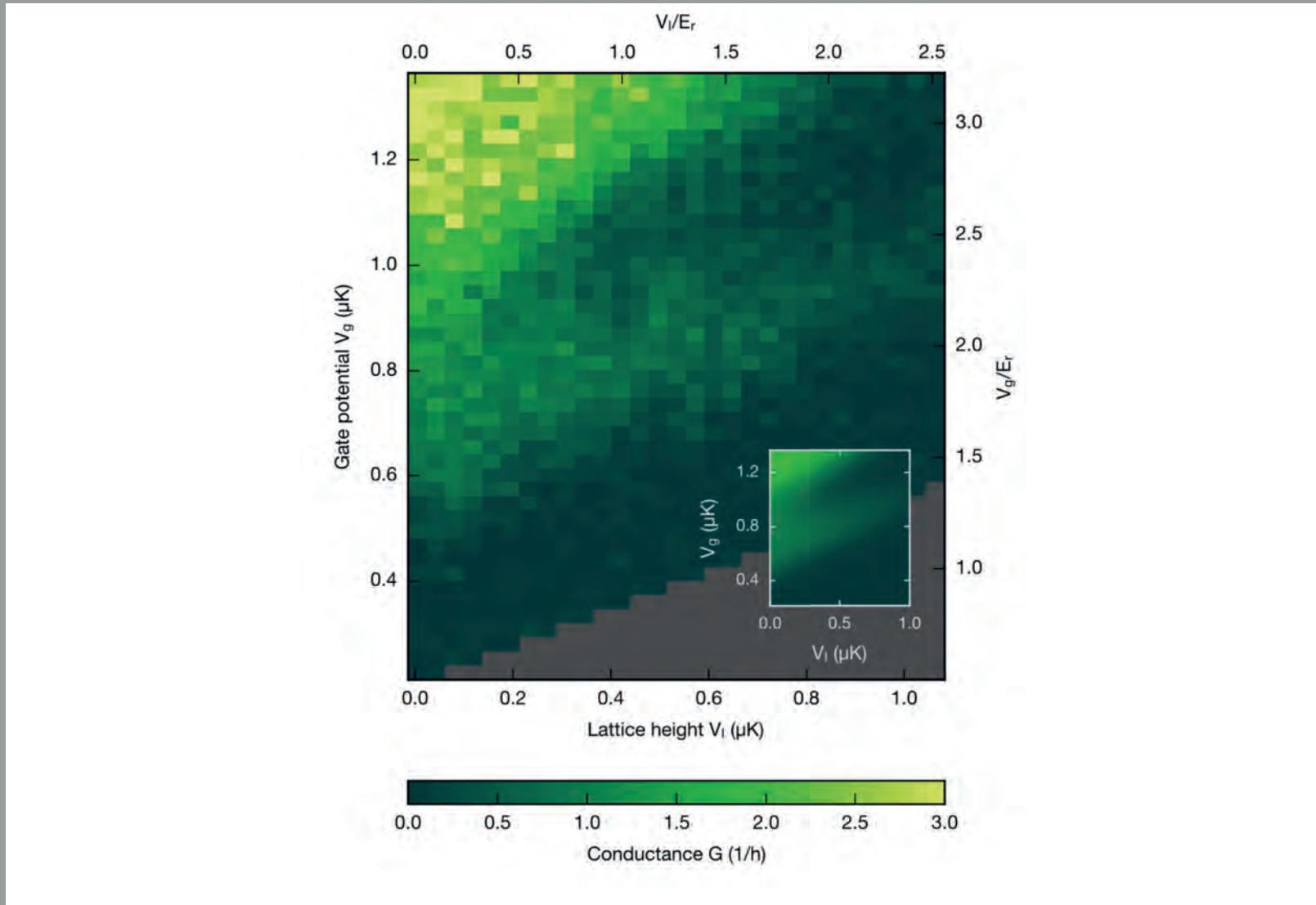
# Transport



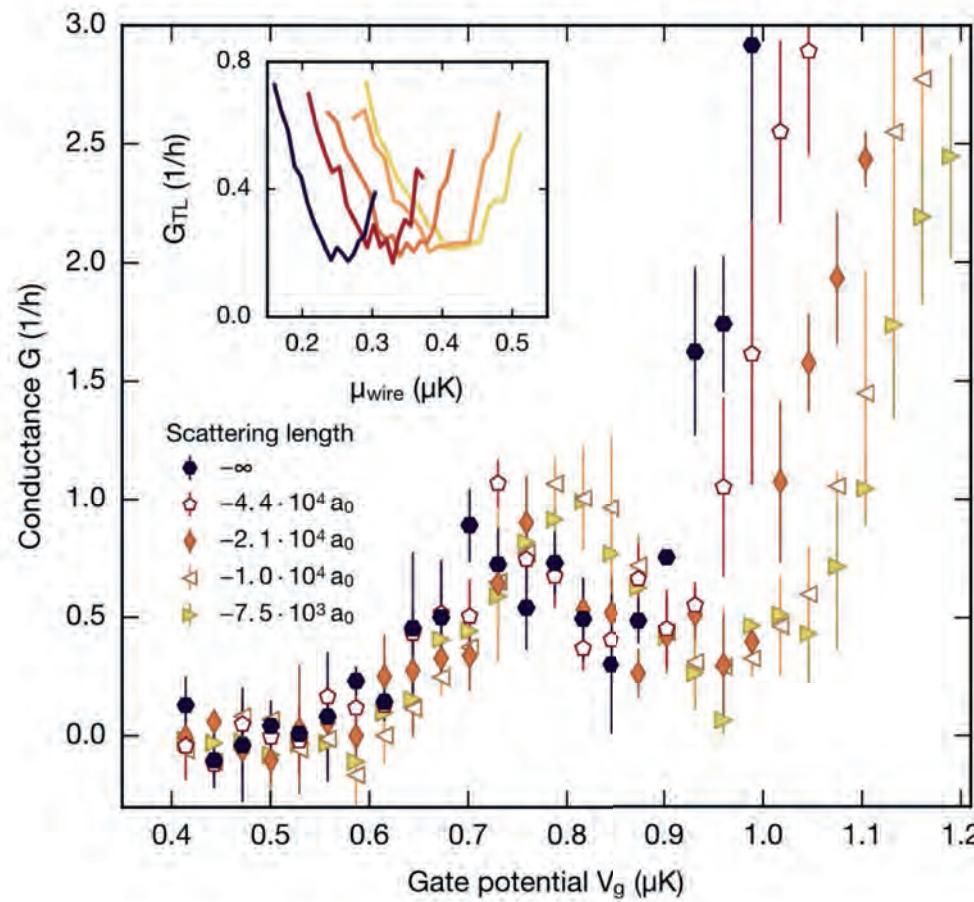
# Site by site



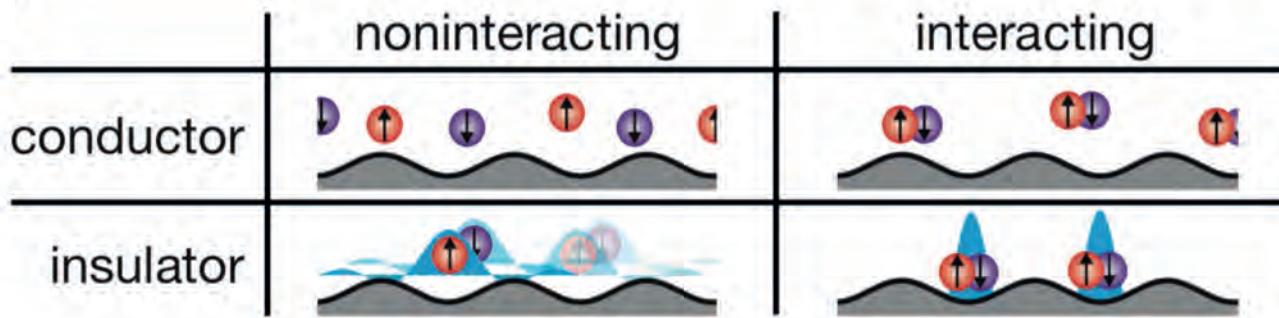
# Increase lattice



# Interactions

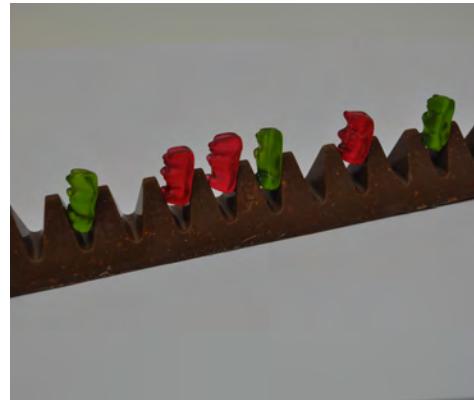
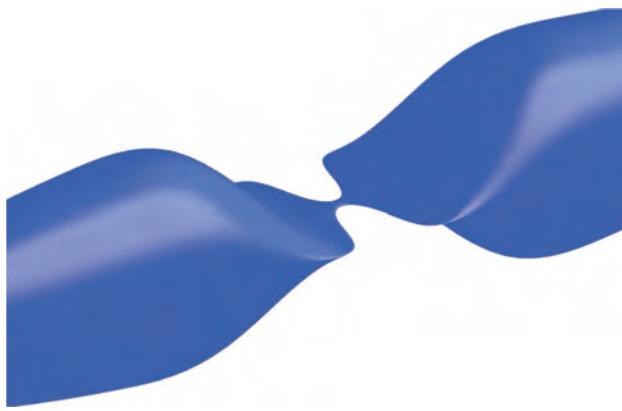


# Luther-Emry liquid



Martin Lebrat, Pjotrs Grišins, Dominik Husmann, Samuel Häusler, Laura Corman,  
Thierry Giamarchi, Jean-Philippe Brantut, T. E., arXiv:1708.01250, PRX 8, 011053 (2018)

# Conclusions



# Thanks!

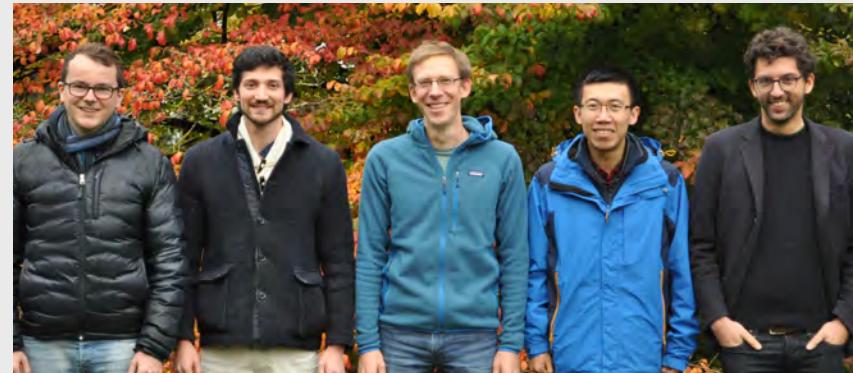
Funding: ETH, SNF, NCCR QSIT, EU SIQS, TherMiQ, QUIC, ERCadv



## Quantum Gases in Optical Lattices

Rémi Desbuquois  
Michael Messer  
**Frederic Görg**  
Kilian Sandholzer,  
Joaquín Minguzzi

Electronics: Alexander Frank  
Administration: Stefanie Ackermann



## Impact experiment

**Tobias Donner**  
Julian Leonard (now Harvard)  
Andrea Morales  
Philip Zupancic  
Xiangliang Li  
Davide Dreon



Wilhelm  
Zwerger



## Lithium Microscope

Laura Corman  
Dominik Husmann  
Martin Lebrat  
Samuel Häusler  
Philipp Fabritius  
Jean-Philippe Brantut (now EPFL)

+ Thierry Giamarchi  
Pjotrs Grisins

## BEC and Cavity

Manuele Landini  
Tobias Donner  
Nishant Dogra  
Katrin Kröger