

Particles and strings in six-dimensional $(2, 0)$ theory

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The six-dimensional $(2, 0)$ theories

In 1995, we learned of the existence of

quantum theories in six space-time dimensions with $(2, 0)$ superconformal symmetry.

Here are three reasons to study these theories:

Finding the right conceptual framework to define them is a challenging problem.

This is likely to involve new interesting mathematical structures.

They give rise to certain Yang-Mills theories with maximally extended supersymmetry upon compactification on a two-torus.

This may be a way to find an S -dual formulation of these lower dimensional theories.

They arise within string/ M -theory as decoupled subsectors localized on certain space-time impurities (branes or singularities).

This may provide an opportunity to study aspects of these higher dimensional theories that are not related to quantum gravity.

Tensor multiplet particles

The best known representation of the $(2, 0)$ supersymmetry algebra is the tensor multiplet of massless space-time fields:

<u>field</u>	<u>$SO(5, 1)$</u>	<u>$SO(5)_R$</u>
ϕ	scalar	vector
ψ	anti-chiral spinor	spinor
$h = db$	self-dual three-form	scalar

Second quantization of these fields yields a Fock space of massless particle states.

These particles transform as scalars, spinors, and self-dual tensors under the little group $SO(4) \subset SO(5, 1)$ of transverse rotations.

Dimensional reduction of these fields gives rise to a massless vector multiplet of maximally extended supersymmetry.

But the two-form b does not naturally couple to a particle.

So it appears difficult to construct a 'non-abelian' tensor multiplet theory in six dimensions.

Self-dual strings

Instead, b naturally couples to a self-dual string.

This is also indicated by the $(2, 0)$ supersymmetry algebra extended with 'central charges':

$$\{Q, Q\} = P + Z.$$

<u>generator</u>	<u>$SO(5, 1)$</u>	<u>$SO(5)_R$</u>
Q (supercharge)	chiral spinor	spinor
P (six-momentum)	vector	scalar
Z (string charge)	vector	vector

A straight string element in the spatial direction V coupled to a tensor multiplet ϕ, ψ, h gives rise to the charge $Z = \langle \phi \rangle \otimes V$. Note that $|P| \geq |Z|$.

A multiplet of straight string states transforms as scalars, spinors, and vectors under the little group $SO(4) \subset SO(5, 1)$ of transverse rotations.

Dimensional reduction along a tensile string gives rise to the particles of a massive vector multiplet of maximally extended supersymmetry.

So (spontaneously broken) non-abelian gauge symmetry may appear after compactification!

The Lagrangian formulation

Continuous symmetries spontaneously broken by a string gives rise to Goldstone fields on the world-sheet $\Sigma \subset \mathbb{R}^{1,5}$:

<u>field</u>	<u>$SO(4) \subset SO(5, 1)$</u>
X^\perp (non-chiral)	vector
Θ^+ (left-moving)	chiral spinor
Θ^- (right-moving)	anti-chiral spinor

The dynamics is governed by the action

$$S = \frac{1}{4\pi\lambda^2} \int_{\mathbb{R}^{1,5}} d^6x (\partial\phi\partial\phi + \bar{\psi}\not{\partial}\psi + h^2) \\ + \int_{\Sigma} d^2\sigma \sqrt{-\phi\phi} (DX^\perp DX^\perp + \bar{\Theta}^+ \not{D}\Theta^+ + \bar{\Theta}^- \not{D}\Theta^-) \\ + e \int_{\Sigma} b + \dots,$$

where h fulfills the modified Bianchi identity

$$dh = 2\pi q \delta_\Sigma$$

and

- λ = coupling constant
- e = string electric charge
- q = string magnetic charge.

Decoupling of the anti self-dual part of h (not part of the tensor multiplet) requires $e = \frac{1}{\lambda^2}q$.

The coupling constant and the charges

The quantum theory of a chiral two-form b (with self-dual field strength h) can only be defined for a rational value of the coupling constant λ^2 .

The correct value appears to be $\lambda^2 = 2$.

(Consistent embedding into M -theory, commutation relations of Wilson-'t Hooft surface observables for two linked surfaces,...)

The topological class of the 'wave function' of a system of two strings with electric-magnetic charges (e, q) and (e', q') is determined by the quantity

$$e \cdot q' + q \cdot e' \in \mathbb{Z}.$$

The plus sign means that even two identical strings give a non-zero result...

(Cf. dyonic particles in four dimensions.)

With $e = \frac{1}{2}q$ and $e' = \frac{1}{2}q'$ (decoupling of anti self-dual part of h), we get the condition

$$q \cdot q' \in \mathbb{Z}$$

for all magnetic string charges q and q' .

The normal bundle anomaly

Diffeomorphisms of $\mathbb{R}^{1,5}$ that leave Σ invariant appear as $SO(4)$ gauge transformations on the normal bundle $N = (T\Sigma)^\perp$. The theory must respect this symmetry.

But because of the modified Bianchi identity $dh = 2\pi q \delta_\Sigma$, the electric coupling

$$e \int_\Sigma b = e \int_{\mathbb{R}^{1,5}} b \wedge \delta_\Sigma$$

classically suffers from an anomaly inflow described by descent on the four-form

$$I^{\text{class}} = q \cdot e \delta_\Sigma|_\Sigma = q \cdot e \chi(N).$$

The chiral and anti-chiral fermions Θ^+ and Θ^- give a further one-loop contribution:

$$I^{\text{quant}} = -\chi(N).$$

Anomaly cancellation and the relation $e = \frac{1}{2}q$ thus imply that

$$q \cdot q = 2$$

for all magnetic string charges q .

The ADE-classification

The general solution to the conditions $q \cdot q' \in \mathbb{Z}$ and $q \cdot q = 2$ are labeled by the series $A_{1,2,\dots}$, $D_{4,5,\dots}$, or $E_{6,7,8}$.

This is in agreement with the classification of simple singularities of hyper-Kähler four-folds $K \simeq \mathbb{C}^2/\Gamma$ for $\Gamma =$ discrete subgroup of $SU(2)$. (Type IIB string theory on $\mathbb{R}^{1,5} \times K$ gives a realization of $(2, 0)$ theory.)

More mysteriously, it also indicates that a compactified $(2, 0)$ theory may only give rise to Yang-Mills theories with maximally extended supersymmetry and a simply laced gauge group (*i.e.* $SU(r + 1)$, $SO(2r)$, or $E_{6,7,8}$).

In the Coulomb phase of the Yang-Mills theory:

Massless vector multiplets associated with Cartan generators originate from massless tensor multiplets.

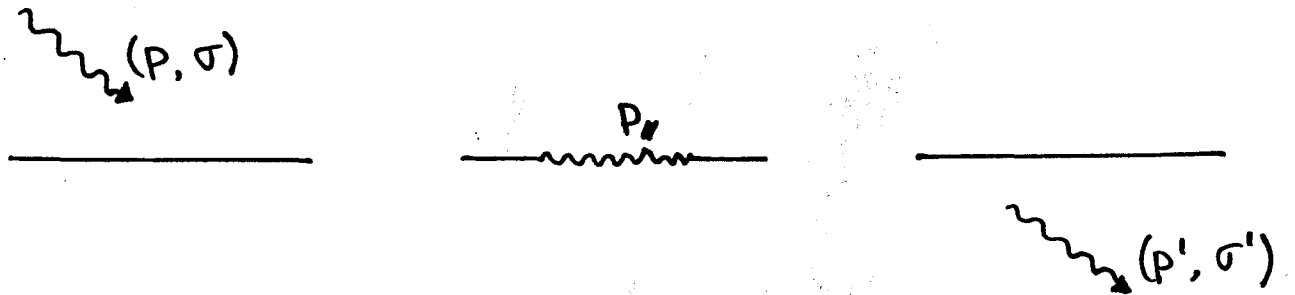
Massive vector multiplets associated with root generators originate from tensile self-dual strings.

Scattering processes

Consider an infinitely extended, approximately straight and static 'bare' string. It is surrounded by a non-vanishing configuration of the space-time tensor multiplet fields ϕ , ψ , and h .

We would like to change variables and work with a 'dressed' string, which includes this self-field.

Because of the $\sqrt{\phi\phi}$ string tension, tensor multiplet quanta of energy E then couple to string wave quanta with strength $E/(\sqrt{\phi\phi})^{1/2}$. Tree-level scattering of tensor multiplets off a string:



For $E^2 \ll \sqrt{\phi\phi}$, the tensor multiplets are described by a free space-time field theory, and the string waves are described by some decoupled world-sheet conformal field theory.

But which two-dimensional CFT?

The decoupled world-sheet theory

The degenerate ground states of the string are a representation of the Clifford algebra of the Θ^+ and Θ^- zero modes.

Above these ground states, there are four bosonic (X^\perp) and four fermionic (Θ^+ and Θ^-) polarizations of string wave excitations.

But because of the electro-magnetic interaction, the bosonic string waves are not free.

Replace X^\perp with 'spherical' variables $r \in \mathbb{R}$ and $g \in S^3 \simeq SU(2)$, modulo $(r, g) \sim (-r, -g)$. The normal bundle structure group acts as $g \mapsto ugv^{-1}$ for $(u, v) \in SU(2) \times SU(2) \simeq Spin(4)$.

The 'radial' coordinate r is a free boson, but g is governed by a 'level one' $SU(2)$ Wess-Zumino-Witten model. This reproduces the classical bosonic normal bundle anomaly.

The WZW kinetic term of g implies that the effective space-time geometry probed by the string wave fluctuations is blown up at the locus of the string.

Conclusion...

The six-dimensional $(2, 0)$ theories exhibit a tightly constrained structure, apparent already in the effective low-energy description in terms of particles and strings:

The classical theory is inconsistent.

Supersymmetry is essential.

There is an ADE-classification.

The string world-sheet theory is non-trivial.

...and outlook

Investigations of scattering processes, particularly in a compactified situation, may give insights into the six-dimensional origins of (spontaneously broken) non-abelian gauge symmetry in certain lower-dimensional Yang-Mills theories.

Particles and strings are useful at a generic point in the moduli space with non-vanishing $\langle \phi \rangle$.

A major challenge is to instead use the full superconformal symmetry of the theory. But we do not yet know the correct framework for such a construction.