

2. Which model for the interaction potential?

2.1. Basic postulate

In the quasi limit $m^3 \ll 1$, at low temperature (say $m^3 > 1$), short range details of $V(r)$ are not important.

Only low k two-body scattering amplitude f_k^{two} matters. \Rightarrow Take a convenient model such that

$f_k^{\text{model}} \approx f_k^{\text{true}}$ at low k (in particular $k \lambda \ll 1$)

2.2. Reminder of scattering theory in 3D

* The Schrödinger equation at $E = \frac{\hbar^2 k^2}{2m} \geq 0$

$$E \phi(r) = \left[-\frac{\hbar^2}{2m} \Delta_r^2 + V(r) \right] \phi(r)$$

with boundary conditions at $r \rightarrow \infty$:

$$\phi(r) \approx e^{i k r} + f_k(m) \frac{e^{i k r}}{r}$$

scattered wave in direction

$$\hat{n} = \hat{r}/r$$

$$= \langle \hat{r} | \hat{n} \rangle$$

incoming wave

$\lambda_e = \text{effective range (any sign)}$

$a = \text{scattering length (any sign)}$

imposed by uncertainty

$$f_k = \frac{\frac{a}{\lambda} + \lambda k - \frac{1}{2} k^2 \lambda_e + \dots}{-1}$$

$f_k(m)$ becomes rationally invariant: δ -wave scattering

* low k limit:

state energy $E_{\text{dim}} = -\frac{\hbar^2 k^2}{2m}$

* pole of f_k for $k = i\eta$, $\eta > 0$ gives bound

Born approximation $G_0(z) = \frac{1}{z - \frac{\hbar^2}{2m}}$

* Born expansion (in powers of V): $T(z) = \underbrace{V}_{\checkmark} + V G_0(z) V + V G_0(z) V G_0(z) V + \dots$

resolvent $G(z) = \frac{1}{z - [\frac{\hbar^2}{2m} + V(z)]}$

T matrix: $T(z) = V + V G(z) V$ $z \in \mathbb{C}$

* formal solution for $f_k(m)$: $f_k(m) = -\frac{m}{\hbar^2} \langle k m | T(E + i0^+) | k \rangle$

in all my lectures, gave wrong limit

$$k^2 |r_e| \ll \left| \frac{a}{1 + ik} \right| \quad \text{for relevant } k$$

in practice, a is the only parameter [Δ only
in weakly interacting regime for $N > 2$ bosons]

2.3 Resonant models

2.3.1 Contact potential (Fermi)

$$V(r) = g \delta(r) = g |r=a| < r=a |$$

has a meaning only in the Born approximation
[or only in 1D]

$$f_0 = -a = -\frac{4\pi\hbar^2}{m} \langle R=a | V | R=a \rangle \Rightarrow g = \frac{4\pi\hbar^2}{m} a$$

[next order in Born expansion gives ∞]

2.3.2 Bethe - Goldstone model

→ potential replaced by contact conditions:

$$\exists A / \phi(r) = A \left[\frac{1}{1} - \frac{a}{r} \right] + O(r)$$

→ large r form of $\phi(r)$ valid $\forall r > 0$:

$$1 + ik \left[\frac{1}{1} + ik \right] = A \left[\frac{1}{1} - \frac{a}{r} \right]$$

$$f_k = \frac{-1}{\frac{1}{a} + ik}$$

→ bound state: $k = i\eta$

$$\frac{1}{a} = \eta$$

$$E_{\text{dim}} = -\frac{\hbar^2}{ma^2}$$

→ which Schrödinger equation?

$$\Delta \frac{1}{r} = -4\pi \delta(r)$$

$$\frac{\hbar^2}{m} \Delta \phi \text{ contains } -\frac{\hbar^2}{m} 4\pi A \delta(r) = g \delta(r) \phi(r)$$

$$E \phi(r) = -\frac{\hbar^2}{m} \Delta \phi(r) + g \delta(r) \phi(r)$$

Ansatz: $\phi(r) = \frac{1}{r} u(r)$

Δ for s-wave, radial equation makes up to 2-body

correlations, not for 3-body correlations.

2.3.3 lattice model

$$R \in \mathbb{Z}^3$$

$$R \in D = [-\frac{a}{2}, \frac{a}{2}]^3$$

$$v = g_0 |R = \vec{a}\rangle \langle R = \vec{a}| \text{ (in COM frame)}$$

* adjust g_0 to have correct scattering length:

$$\frac{g_0}{1} + \int \frac{d^3k}{(2\pi)^3} \frac{1}{m \frac{\hbar^2 k^2}{m} - \frac{1}{g}} = \frac{1}{g}$$

Effort in calculation $\langle R | T(E+i0^+) | R \rangle$, $\langle R=8 | G | R=8 \rangle$

appears, obtained from $G = G_0 + G_0 V G$

* discussion: renormalization of coupling constant

here g_0 depends on momentum cut-off of $\frac{1}{L}$

$$g_0 = \frac{g}{1 - c a^2}$$

$$c = 2, 4, \dots$$

→ Born regime: $|a| \ll 1$, $g_0 \approx g$

$g_0 > 0$ for $a > 0$

→ limit $\ell \rightarrow 0$:

$$g_0 \approx -\frac{c}{4\ell} \frac{\hbar^2 \ell^2}{m} < 0$$

• Bad for $N > 2$ bosons: ground state not a gas but off-gas state

Effort: $\langle N: R=8 | H | N: R=8 \rangle \rightarrow -\infty$

$\ell \rightarrow 0$

• no problem for fermions (fermion, not fermion):

• ground state = gas

• QMC with no sign problem

• equivalence with Bethe Ansatz $\ell \rightarrow 0$